

# $\pi\pi$ scattering lengths from NA48/2 data on Ke4 and three- pion decays of charged kaons

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on behalf of the **NA48/2** Collaboration:

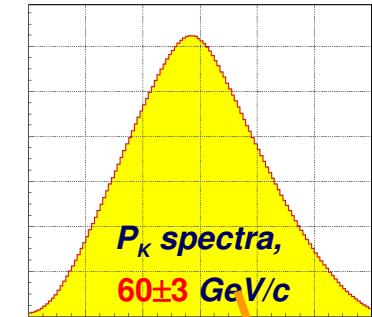
Cambridge, CERN, Chicago, Dubna, Edinburgh, Ferrara,  
Firenze, Mainz, Northwestern, Perugia, Pisa, Saclay,  
Siegen, Torino, Vienna



# Overview

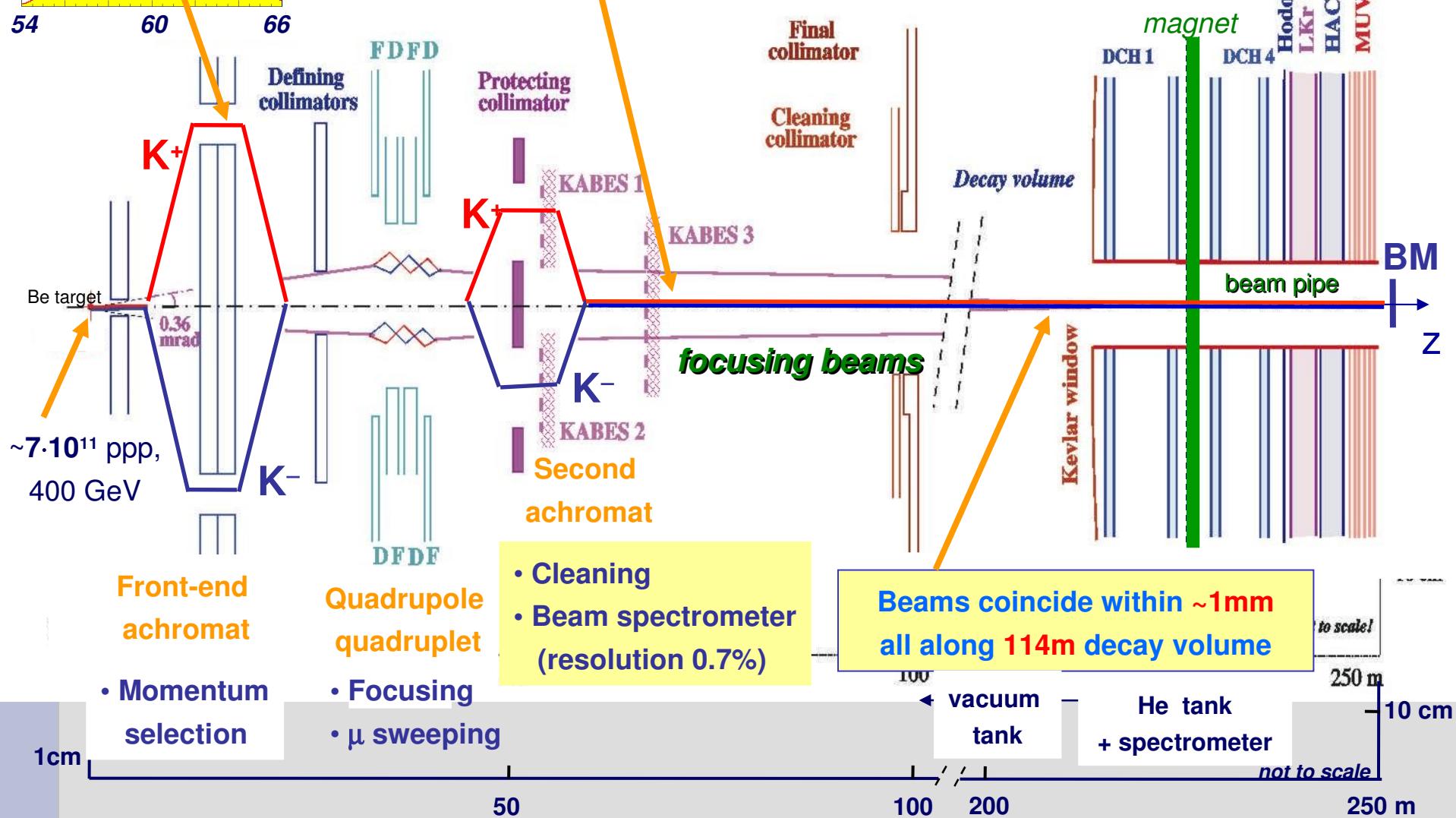
- NA48/2 experiment at CERN: setup, data taking
- Study of  $\mathbf{K}^\pm \rightarrow \pi^\pm \pi^0 \pi^0$  Dalitz plot distribution:  
an updated measurement of  $\pi\pi$  scattering lengths  
( $a_0 - a_2$ ,  $a_2$ ) and the Dalitz plot slopes
- **Ke4** ( $\mathbf{K}^\pm \rightarrow \pi^+ \pi^- e^\pm \nu$ )
  - Form factors and pion scattering lengths
  - Data 2003: Preliminary results

# NA48/2 beam line



2-3M K/spill ( $\pi/K \sim 10$ ),  
 $\pi$  decay products stay in pipe.  
 Flux ratio:  $K^+/K^- \approx 1.8$

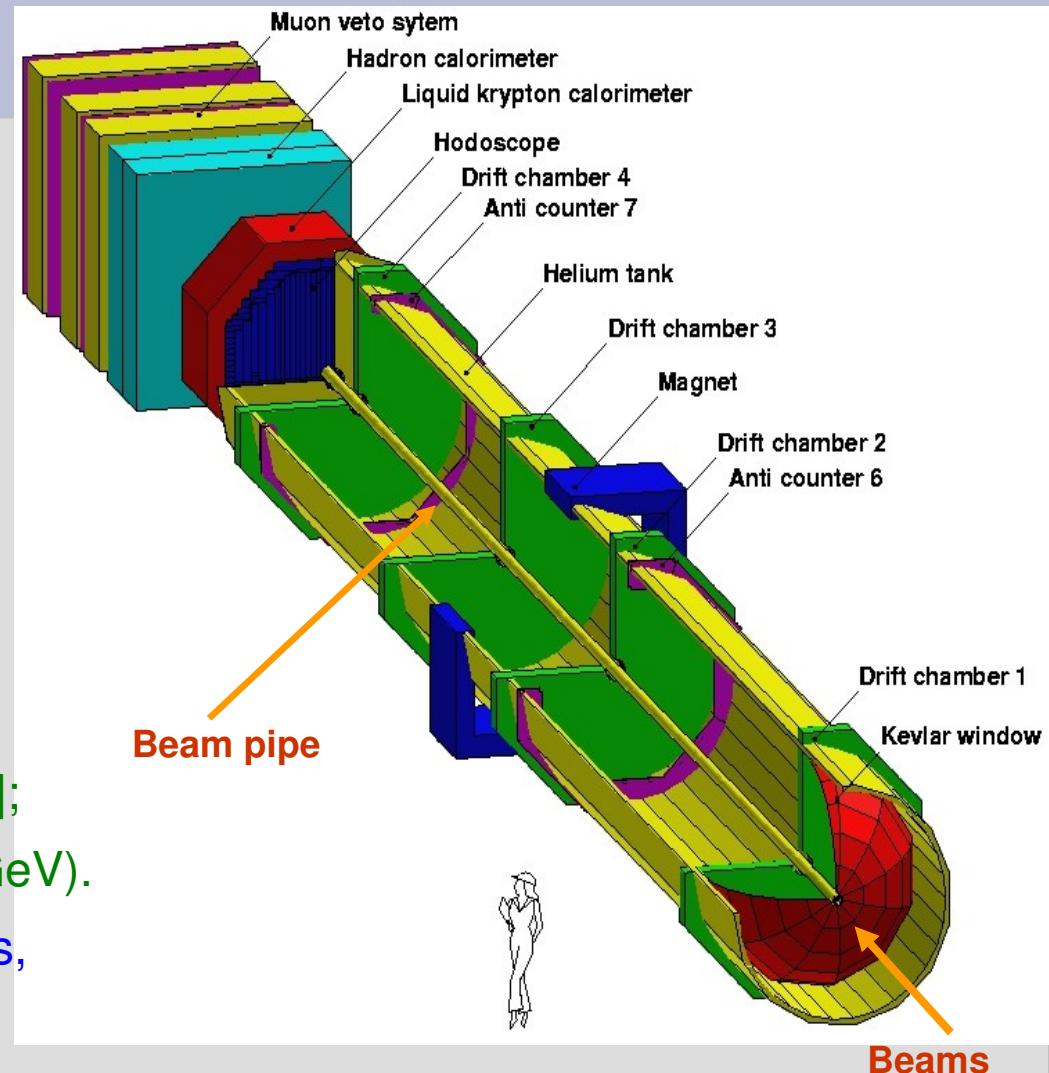
Simultaneous  $K^+$  and  $K^-$  beams:  
 large charge symmetrization of  
 experimental conditions



# The NA48 detector

## Main detector components:

- Magnetic spectrometer (4 DCHs):  
4 views/DCH: redundancy  $\Rightarrow$  efficiency;  
used in trigger logic;  
 $p/p = 1.0\% + 0.044\% * p$  [GeV/c].
- Hodoscope  
fast trigger;  
precise time measurement (150ps).
- Liquid Krypton EM calorimeter (LKr)  
High granularity, quasi-homogenous;  
 $\sigma_E/E = 3.2\%/E^{1/2} + 9\%/E + 0.42\%$  [GeV];  
 $\sigma_x = \sigma_y = 0.42/E^{1/2} + 0.6\text{mm}$  (1.5mm@10GeV).
- Hadron calorimeter, muon veto counters,  
photon vetoes.



# NA48/2 data:

2003 run: ~ 50 days

2004 run: ~ 60 days

A view of the NA48/2 beam line



Total statistics in 2 years:

$K^\pm \rightarrow \pi^-\pi^+\pi^\pm$ : ~ $4 \cdot 10^9$

$K^\pm \rightarrow \pi^0\pi^0\pi^\pm$ : ~ $1 \cdot 10^8$

Rare  $K^\pm$  decays:

BR's down to  $10^{-9}$

can be measured

>200 TB of data recorded

## $\pi\pi$ scattering lengths

The important free parameter of ChPT is the quark condensate  $\langle qq \rangle$ , it determines the relative size of mass and momentum terms in the power expansion.

$a_0$  and  $a_2$  are S-wave  $\pi\pi$  scattering lengths in isospin states  $I=0$  and  $I=2$ , correspondingly. They enter into all  $\pi\pi$  scattering amplitudes.

The relation between  $\langle qq \rangle$  and the scattering lengths  $a_0$  and  $a_2$  is known from this theory with a high precision, so the experimental measurement of  $a_0$  and  $a_2$  provides an important constraints for ChPT Lagrangian parameters.

# **Study of the cusp-effect in $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ decays**

[2003 result: PLB 633 (2006) 173,  
Cabibbo-Isidori theoretical framework]

[Now: a preliminary 2003+2004 result,  
updated selection and analysis of systematic effects,  
the same theoretical framework]

# $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ selection

For each photon pair  $(i, k)$  a decay vertex reconstructed along beam axis under the assumption of  $\pi^0 \rightarrow \gamma\gamma$  decay

$$m_0^2 = 2E_i E_k (1 - \cos\alpha) \approx E_i E_k \alpha^2 = E_i E_k \frac{(D_{ik})^2}{(z_{ik})^2}$$

$m_0^2$  – mass of  $\pi^0$

$E_i, E_k$  – energy of  $\gamma_i, \gamma_k$

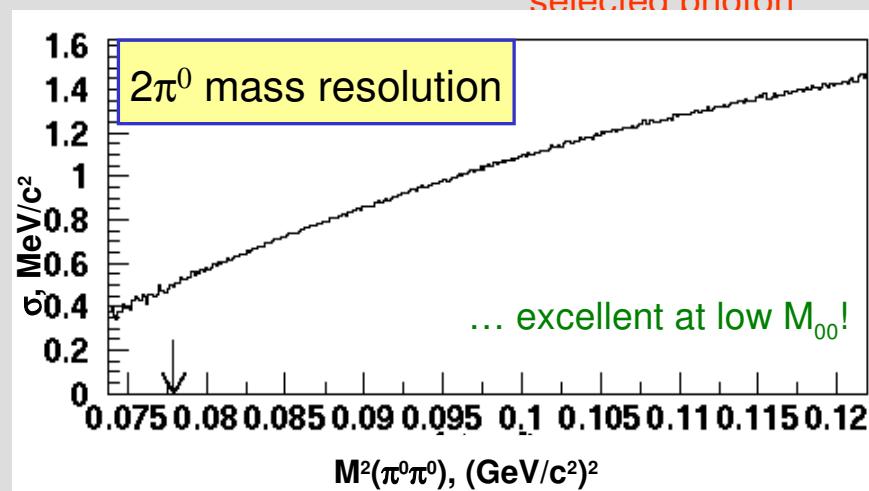
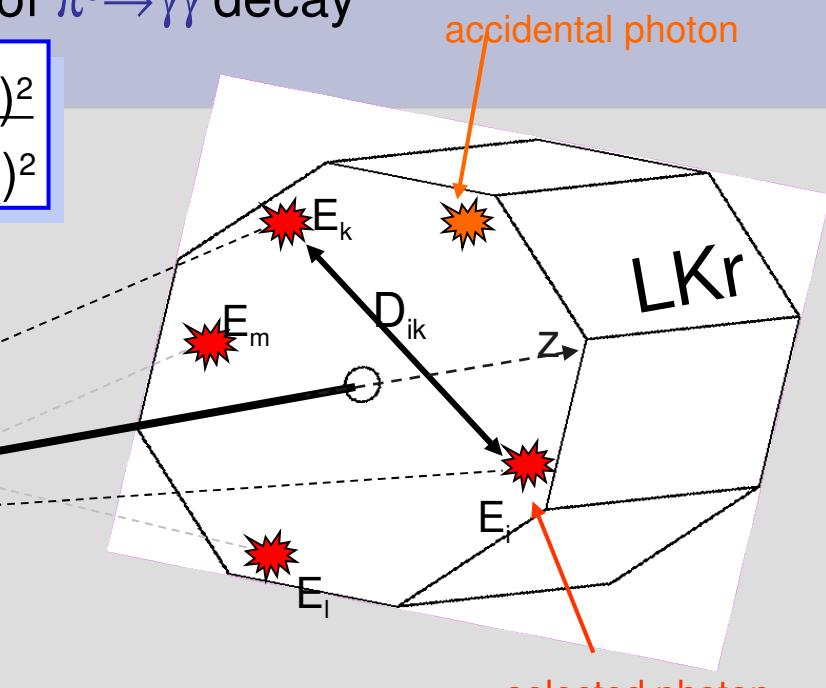
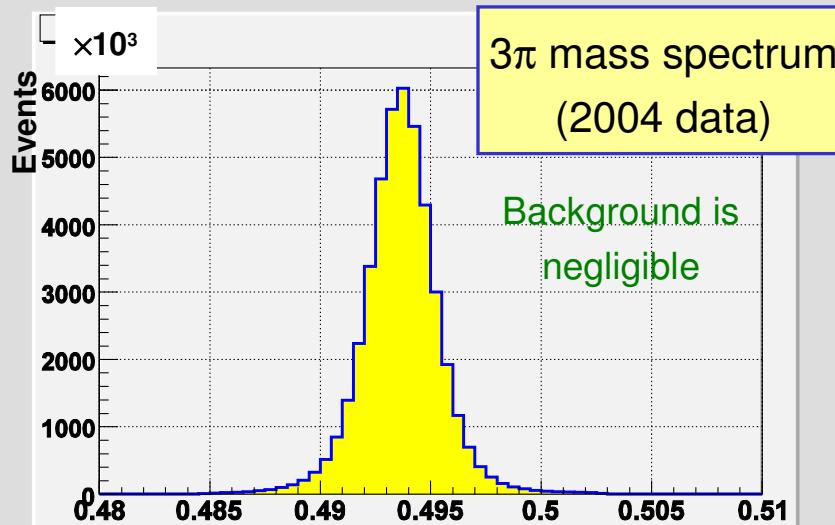
$D_{ik}$  – distance between  $\gamma_i$  and  $\gamma_k$  on LKr

$z_{ik}$  – distance from  $\pi^0 \rightarrow \gamma\gamma$  decay vertex to LKr

**K-decay vertex**

$z_{ik}$

$\Delta z$



# Observation of the cusp

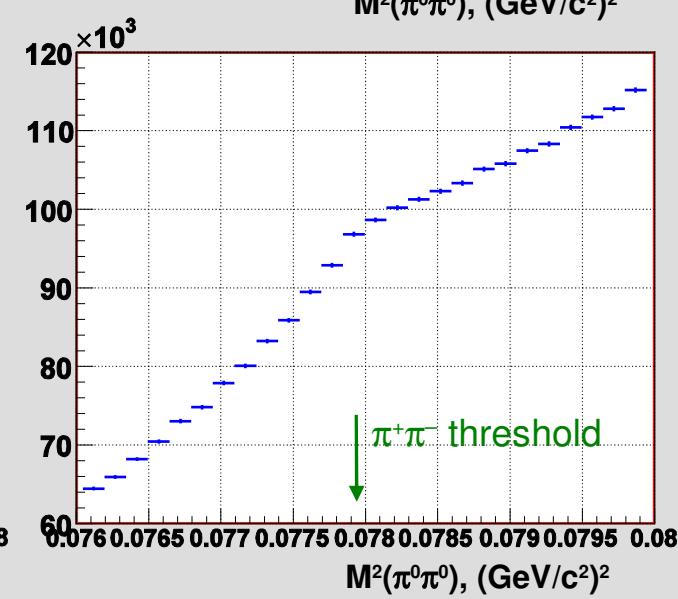
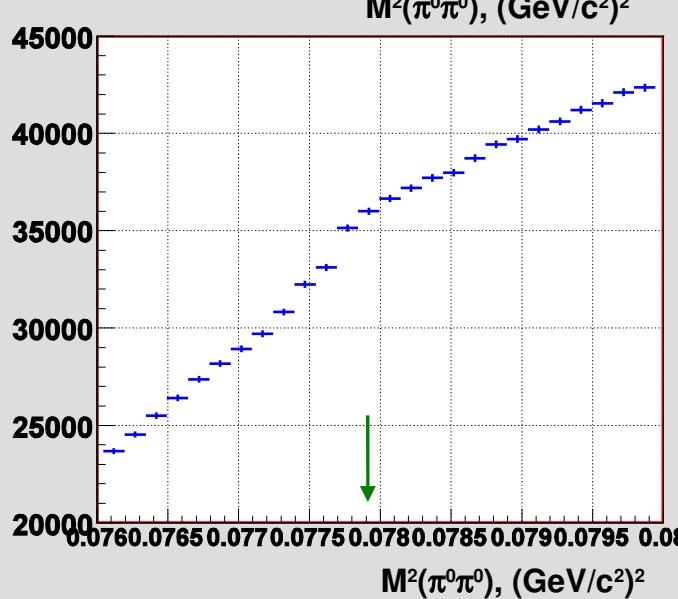
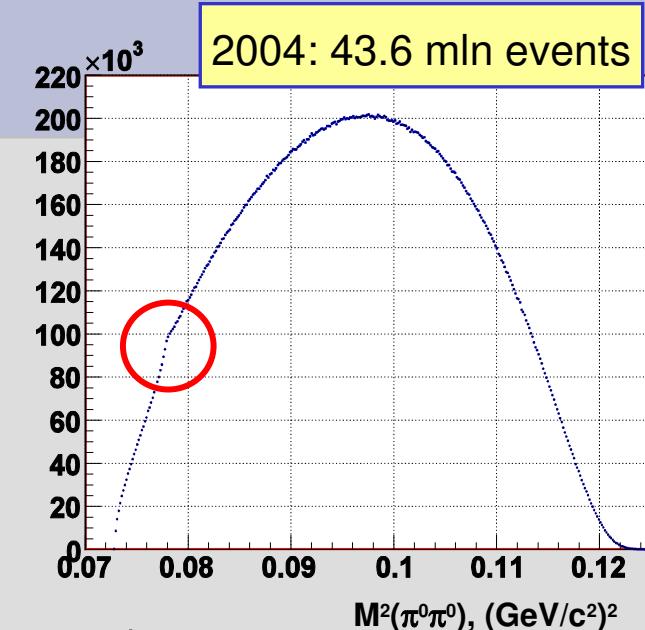
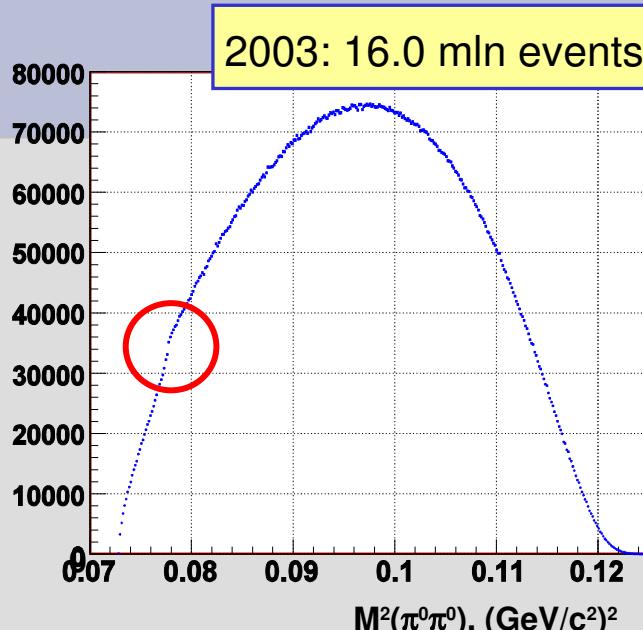
First observation of the cusp was made with the 2003 data

Inclusion of 2004 data: statistics increased by a factor of 3.7!

2003 and 2004 data: slightly different selection conditions and different data/MC sample ratios.



Analysis is made independently for the two data sets, and then the results are averaged.



# Theory: final state rescattering

N. Cabibbo, PRL 93 (2004) 121801

$$M(K^\pm \rightarrow \pi^\pm \pi^0 \pi^0) = M_0 + M_1$$

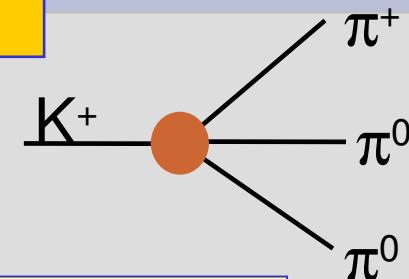
Direct emission:

$$M_0 = A_0(1 + g_0 u/2 + h' u^2/2 + k' v^2/2)$$

Kaon rest frame:

$$u = 2m_K(m_K/3 - E_{\text{odd}})/m_\pi^2$$

$$v = 2m_K(-E_1 - E_2)/m_\pi^2$$



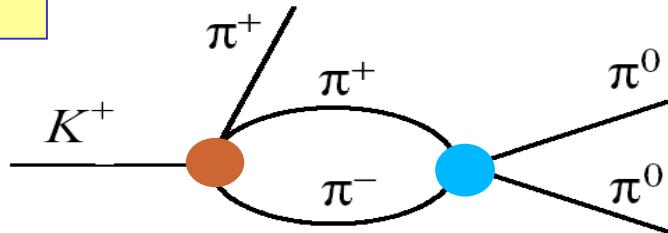
Rescattering amplitude:

$$M_1 = -2/3(a_0 - a_2)m_+ M_+ \sqrt{1 - (\frac{M_{00}}{2m_+})^2}$$

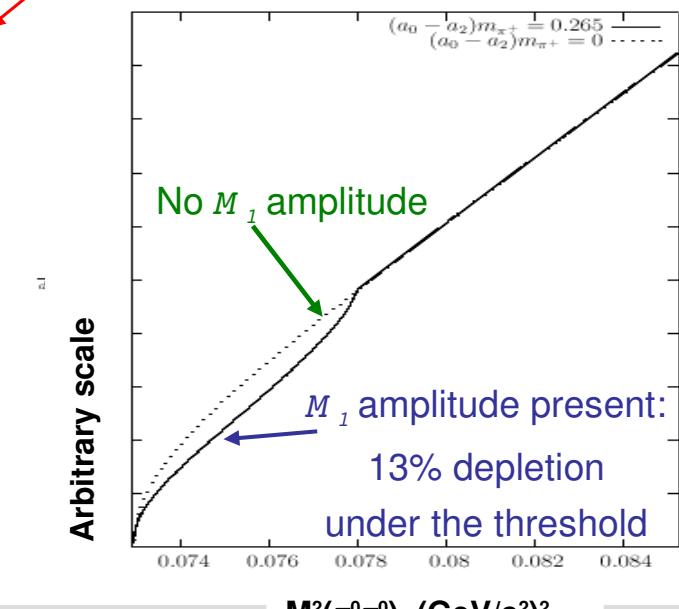
Combination  
of S-wave  $\pi\pi$  scattering  
lengths

(isospin symmetry  
assumed here)

$K^\pm \rightarrow 3\pi^\pm$  amplitude  
at threshold



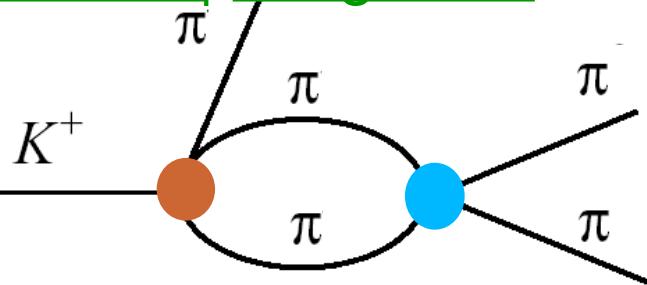
Negative interference  
under threshold



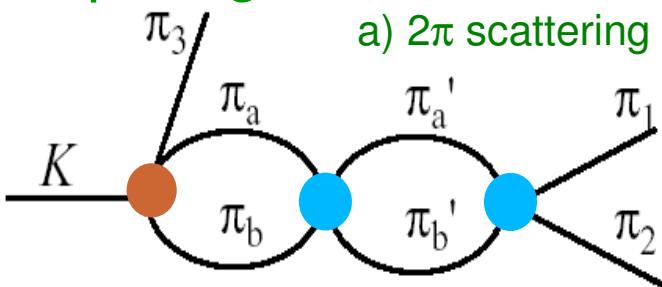
# Theory: two-loop diagrams

N. Cabibbo and G. Isidori,  
JHEP 503 (2005) 21

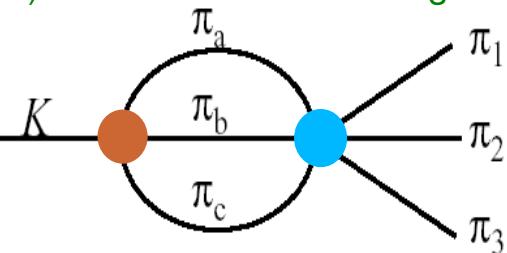
## One-loop diagrams:



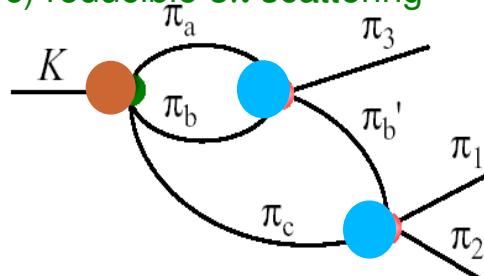
## Two-loop diagrams:



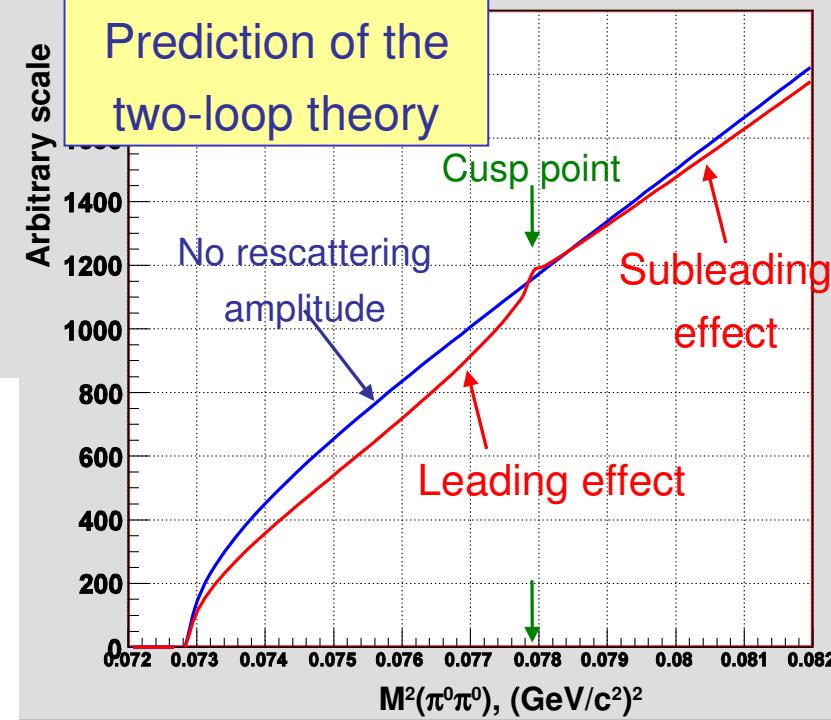
b) irreducible 3 $\pi$  scattering



c) reducible 3 $\pi$  scattering



- Five S-wave scattering lengths ( $a_x, a_{++}, a_{+-}, a_{+0}, a_{00}$ ) expressed as linear combinations of  $a_0$  and  $a_2$
- Isospin symmetry breaking accounted for following J. Gasser
  - For example,  $a_x = (1+\varepsilon/3)(a_0 - a_2)/3$ , where  $\varepsilon = (m_+^2 - m_0^2)/m_+^2 = 0.065$  is isospin breaking parameter
  - All rescattering processes at one- & two-loop level
  - Radiative corrections missing:  $(a_0 - a_2)$  precision  $\sim 5\%$

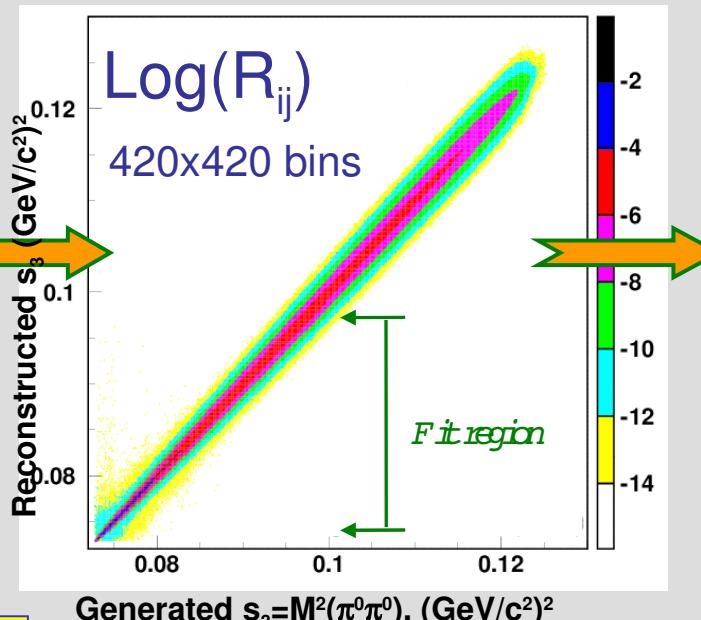


# Fitting procedure

## 1-dimensional fit of the $M_{00}$ projection

Detector response matrix  $R_{ij}$  obtained with a GEANT-based Monte-Carlo simulation

Generated distribution  
 $G(M_{00}) = G(g_0, h', a_0, a_2, M_{00})$



Reconstructed distribution:

$$F_j^{MC} = \sum R_{ij} G_i$$

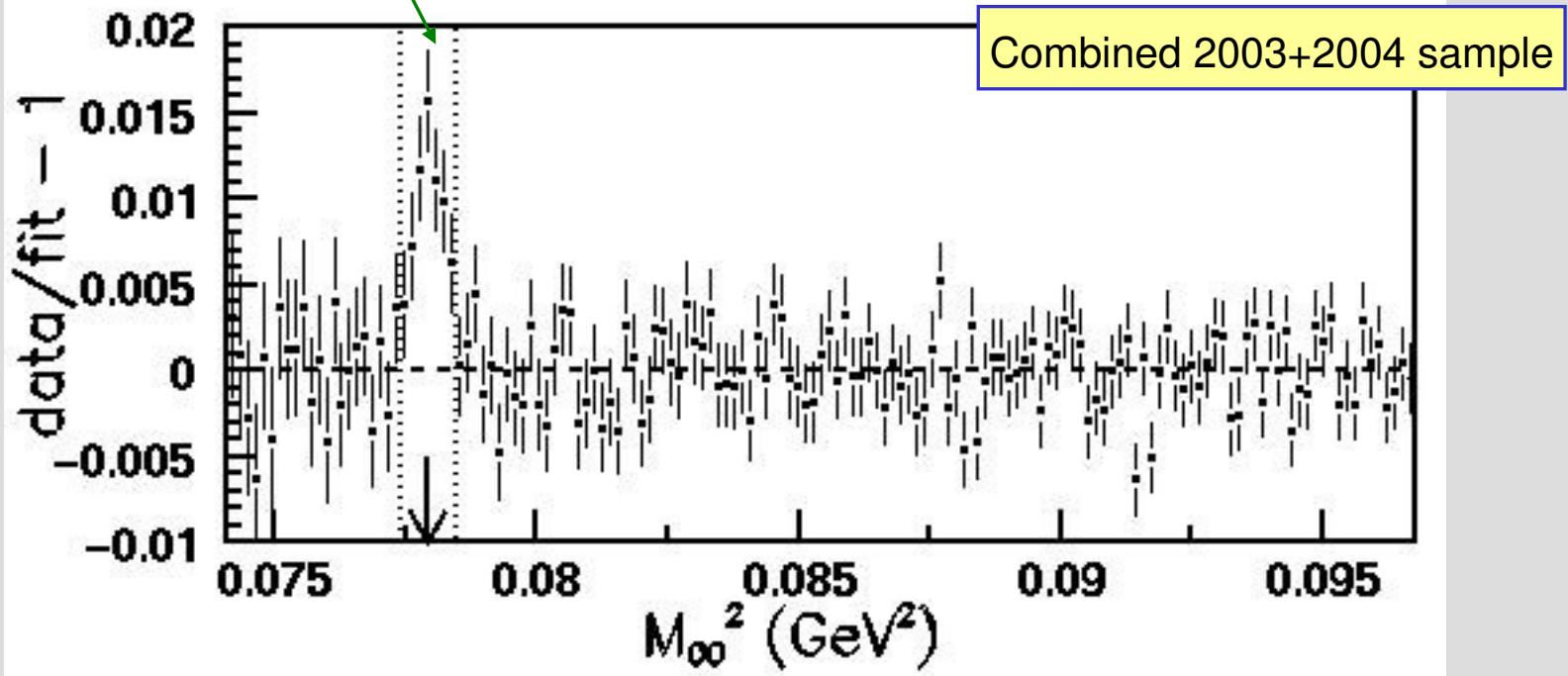
MINUIT minimization of  $\chi^2$   
of data/MC spectra shapes

$$\chi^2(g, h', m_+ (a_0 - a_2), m_+ a_2, N) = \sum_{s_3 \text{ bins}} \frac{(F_{\text{DATA}} - N F_{\text{MC}})^2}{\delta F_{\text{DATA}}^2 + N^2 \delta F_{\text{MC}}^2}$$

# Fit quality & pionium signature

Points excluded from the fit  
due to absence of EM corrections  
in the used model

7 data bins skipped around  
the  $M(\pi^+\pi^-)$  threshold



Excess of events in the excluded interval,

if interpreted as due to pionium decaying as  $A_{2\pi} \rightarrow \pi^0\pi^0$ ,

gives  $R = \Gamma(K^\pm \rightarrow \pi^\pm A_{2\pi}) / \Gamma(K^\pm \rightarrow \pi^\pm \pi^+ \pi^-) = (1.82 \pm 0.21) \times 10^{-5}$ .

➡ Prediction [Z.K. Silagadze, JETP Lett. 60 (1994) 689]:  $R = 0.8 \times 10^{-5}$ .

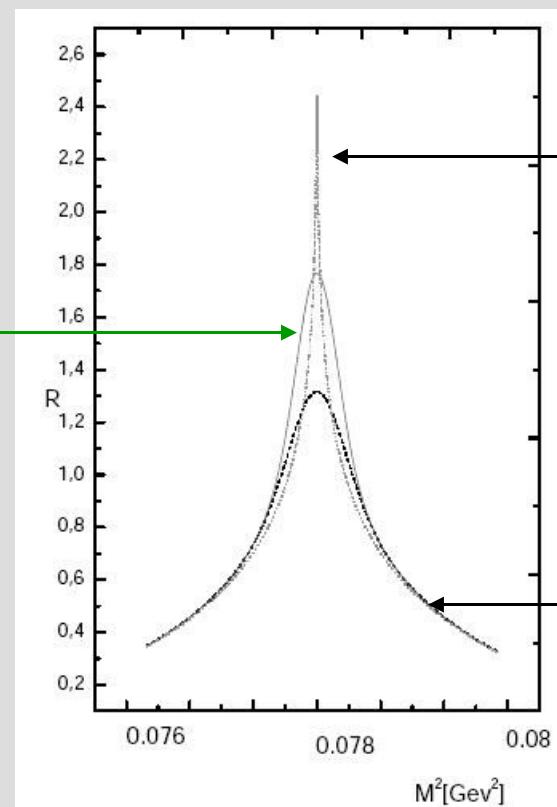
# Electromagnetic corrections to final state interactions in $K \rightarrow 3\pi$ decays

(Gevorkian, Tarasov, Voskresenskaya, hep-ph / 0612129)

**Two contributions from  $K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$  decay to the  $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$  cusp region:**

- Pionium formation :  $\pi^+ \pi^-$  atom  $\rightarrow \pi^0 \pi^0$  (negligible width)
- Additional  $\pi^+ \pi^-$  unbound states with resonance structure  $\rightarrow \pi^0 \pi^0$

$\pi^+ \pi^-$  atoms and  
 $\pi^+ \pi^-$  resonant structure  
with experimental resolution



$\pi^+ \pi^-$  resonant structure  
(no experimental resolution)

$\pi^+ \pi^-$  resonant structure  
with experimental resolution

# Uncertainties & results (1)

(preliminary)

NA48/2 result with 2003+2004 statistics

$$(a_0 - a_2)m_+ = 0.261 \pm 0.006_{\text{stat.}} \pm 0.003_{\text{syst.}} \pm 0.0013_{\text{ext.}}$$

$$a_2 m_+ = -0.037 \pm 0.013_{\text{stat.}} \pm 0.009_{\text{syst.}} \pm 0.0018_{\text{ext.}}$$

- External uncertainty: due to  $R = (A_{++}/A_{+00})|_{\text{threshold}} = 1.975 \pm 0.015$ ;
- Theory precision (rad.corr. & higher order terms neglected):  $\delta(a_0 - a_2)m_+ = 0.013$ .

Systematic effect	$(a_0 - a_2) \times 10^2$	$a_2 \times 10^2$
Analysis technique	$\pm 0.10$	$\pm 0.20$
Trigger inefficiency	negl.	$\pm 0.50$
Description of resolution	$\pm 0.06$	$\pm 0.11$
LKr non-linearity	$\pm 0.06$	$\pm 0.26$
Geometric acceptance	$\pm 0.02$	$\pm 0.01$
MC sample	$\pm 0.03$	$\pm 0.21$
Simulation of LKr showers	$\pm 0.17$	$\pm 0.50$
V-dependence of amplitude	$\pm 0.17$	$\pm 0.38$
Total	$\pm 0.28$	$\pm 0.90$

# Uncertainties & results (2)

Using a chiral symmetry constraint [Colangelo et al., PRL 86 (2001) 5008]:

$$a_2 = -0.0444 + 0.236(a_0 - 0.22) - 0.61(a_0 - 0.22)^2 - 9.9(a_0 - 0.22)^3$$

(preliminary)

$$(a_0 - a_2)m_+ = 0.263 \pm 0.003_{\text{stat.}} \pm 0.0014_{\text{syst.}} \pm 0.0013_{\text{ext.}}$$

→ Theory precision uncertainty is also applicable:  $\delta(a_0 - a_2)m_+ = 0.013$ .

Systematic effect	$(a_0 - a_2) \times 10^2$
Analysis technique	$\pm 0.08$
Trigger inefficiency	negl.
Description of resolution	$\pm 0.06$
LKr non-linearity	$\pm 0.05$
Geometric acceptance	$\pm 0.02$
MC sample	$\pm 0.06$
Simulation of LKr showers	$\pm 0.04$
V-dependence of amplitude	$\pm 0.02$
Total	$\pm 0.14$

# Results: Dalitz plot slopes

$$M(K^\pm \rightarrow \pi^\pm \pi^0 \pi^0) = M_0 + M_1$$

$$M_0 \sim (1 + g_0 u/2 + h' u^2/2 + k' v^2/2)$$

NB: not the same parameterization as the PDG one:

$$|M_0|^2_{(\text{PDG})} \sim (1 + gu + hu^2 + kv^2) \quad [g_0 \approx g, h' \approx h - g^2/4, k' \approx k]$$

Technique: consecutive 1D-fit for rescattering and 2D-fit for the rest.

1.  $(a_0, a_2, g_0, h')$  measurement (fixed  $k'=0$ ) [2003 data] →
2.  $k'$  measurement (fixed  $a_0, a_2, g_0, h'$ ) [2003 data] →
3.  $(a_0, a_2, g_0, h')$  second iteration (fixed  $k'$ ) [2004 data]

$$g = (64.9 \pm 0.3_{\text{stat.}} \pm 0.4_{\text{syst.}})\%$$

$$h' = (-4.7 \pm 0.7_{\text{stat.}} \pm 0.5_{\text{syst.}})\%$$

$$k' = (-0.97 \pm 0.03_{\text{stat.}} \pm 0.08_{\text{syst.}})\%$$

NA48/2 result:

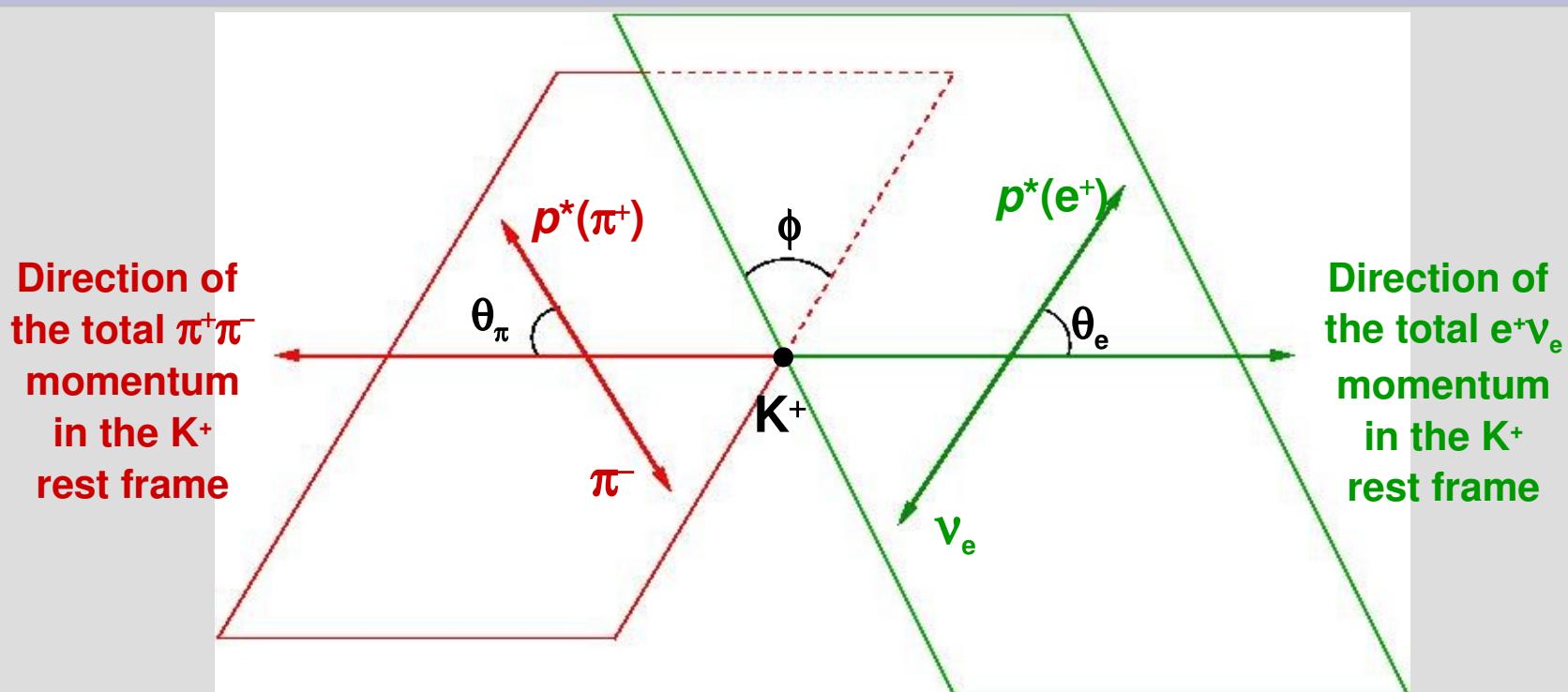
*preliminary*

$(a_0, a_2)$  are not significantly affected by neglecting the  $k'v^2$  term,  
but  $g_0, h'$  are biased by  $\delta g_0 = -1.5\%$ ,  $\delta h = -1.2\%$ )

# $\pi\pi$ scattering in $K^\pm \rightarrow \pi^+\pi^-e^\pm\nu$ decay

# $K_{e4}$ : formalism

A rare decay [ B.R. =  $(4.09 \pm 0.09) \times 10^{-5}$  ] described by five independent variables



Cabibbo – Maksymowicz variables :  $s_\pi \equiv M_{\pi\pi}^{-2}$

$$s_e \equiv M_{e\nu}^{-2}$$

$$\begin{matrix} \theta_e \\ \theta_\pi \\ \phi \end{matrix}$$

For  $K^+ \Rightarrow K^-$

$$\phi \Rightarrow \pi + \phi$$

$$\theta_e \Rightarrow \pi - \theta_e$$

# $K_{e4}$ : formalism

In the partial wave expansion (only S and P waves) the amplitude can be written using form factors:

$$F = F_s e^{i\delta_s} + F_p e^{i\delta_p} \cos\theta_\pi$$

$$G = G_p e^{i\delta_p}$$

$$H = H_p e^{i\delta_p}$$

$F$  ( $F_p, F_s$ ),  $G$ ,  $H$  and  $\delta = \delta_p - \delta_s$  will be used as fit parameters

The form factors can be expanded as a function of  $M_{\pi\pi}^2$  and  $M_{ev}^2$ :

$$F_s = f_s + f_s' q^2 + f_s'' q^4 + f_e' (M_{ev}^2 / 4m_\pi^2) + \dots$$

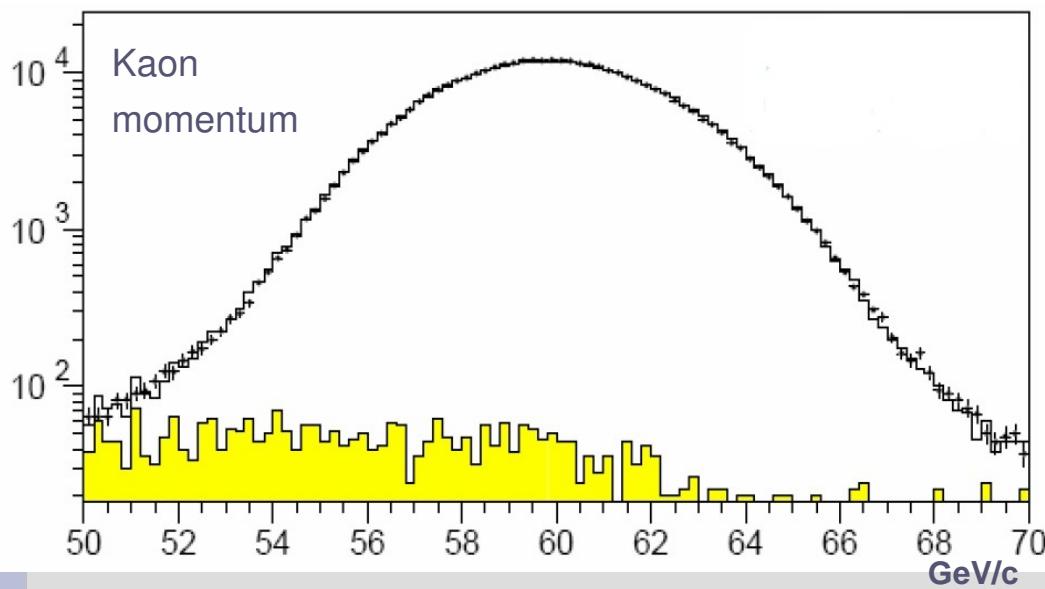
$$F_p = f_p + f_p' q^2 + \dots$$

$$q^2 = (M_{\pi\pi}^2 / 4m_\pi^2) - 1$$

$$G_p = g_p + g_p' q^2 + \dots$$

$$H_p = h_p + h_p' q^2 + \dots$$

# $K_{e4}$ : selection & background



The background is studied using the electron “wrong” sign events (we assume  $\Delta Q = \Delta S$  and total charge  $\pm 1$ ) and cross check with MC. The total bkg is at level of 0.5%.

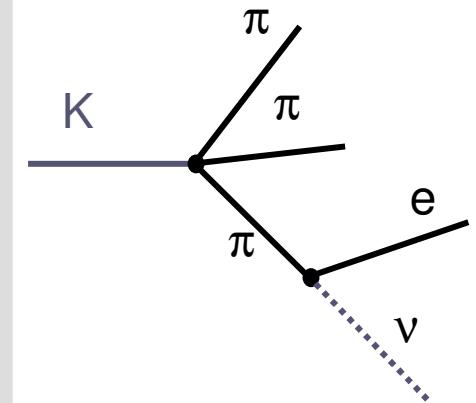
## Selection:

- 3 tracks
- Missing energy and missing Pt
- LKr/DCH energy to electron PID

677500 decays

## Main background sources:

- $\pi\pi\pi + \pi \rightarrow e\nu$
- $\pi\pi\pi$  with  $\pi$  misidentified
- $\pi\pi^0\pi^0$  or  $\pi^0 + \pi^0$  (Dalitz) +  $e\nu$  misidentified and  $\gamma s$  outside the LKr

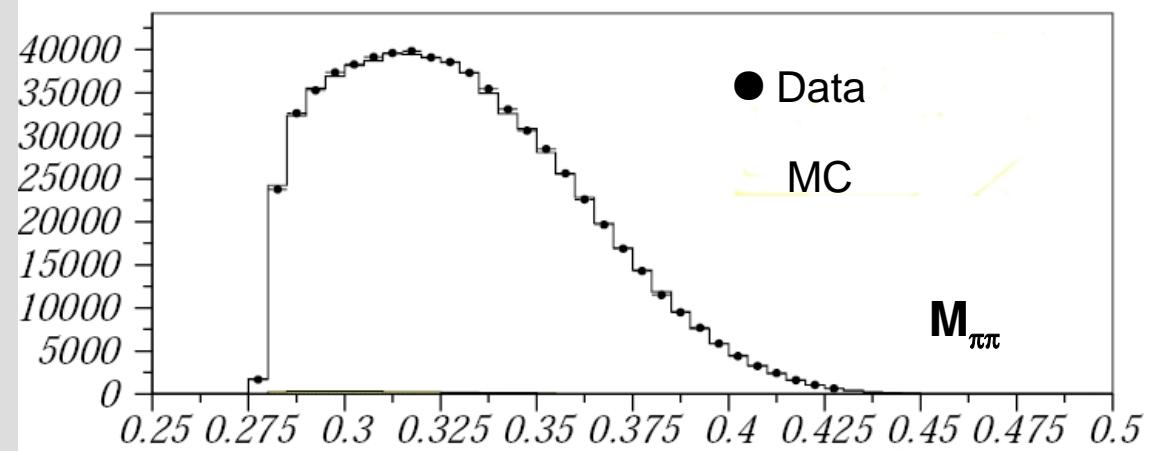


# $K_{e4}$ : Fitting procedure

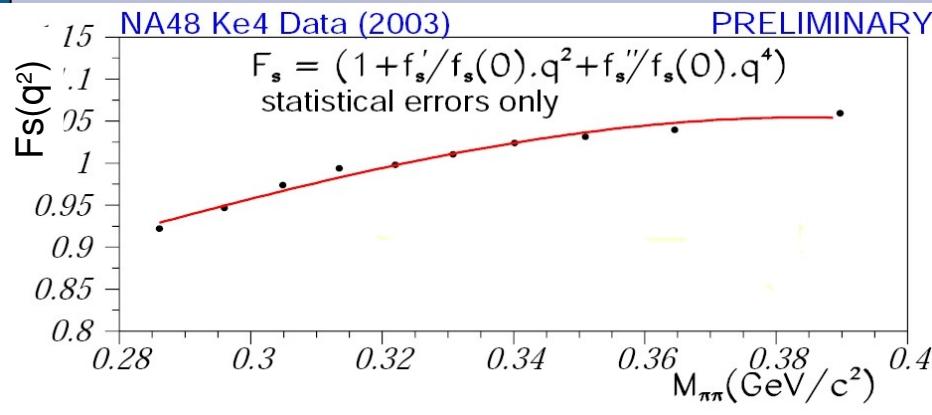
- The form factors ( $F, G, H$  and  $\delta$ ) are extracted minimizing a log-likelihood estimator in each of  $10(M\pi\pi) \times 5(\text{Mev}) \times 5(\cos\theta_e) \times 5(\cos\theta_\pi) \times 12(\phi) = 15000$  equi-populated bins. In each bin the correlation between the 4+1 parameters is taken into account.

- The form factors structure is studied in 10 bins of  $M\pi\pi$ , assuming constant form factors in each bins
- A 2D fit ( $M\pi\pi$ , Mev) is used to study the Fs expansion
- All the results are given wrt to  $F_s(q=0)$  constant term, due to the unspecified overall normalization (BR is not measured)

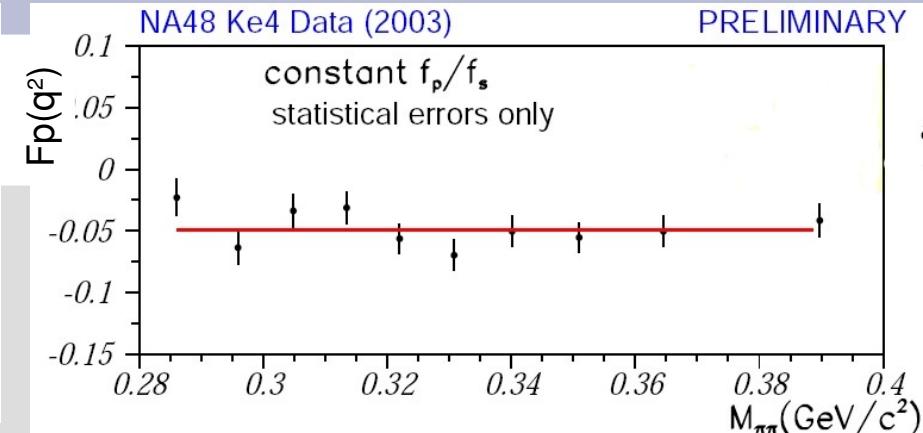
		Data	MC
K+	evts	435654	10.0 M
	Evts/bin	29	667
K-	evts	241856	5.6 M
	Evts/bin	16	373



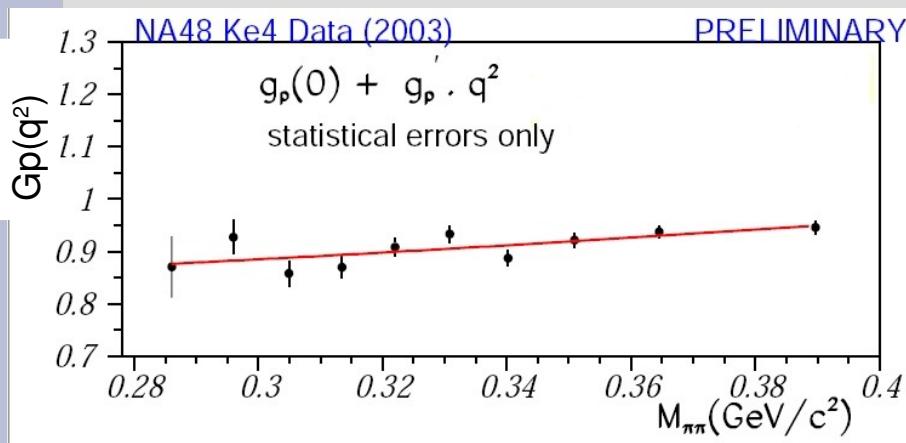
# K<sub>e4</sub>: Fitting results



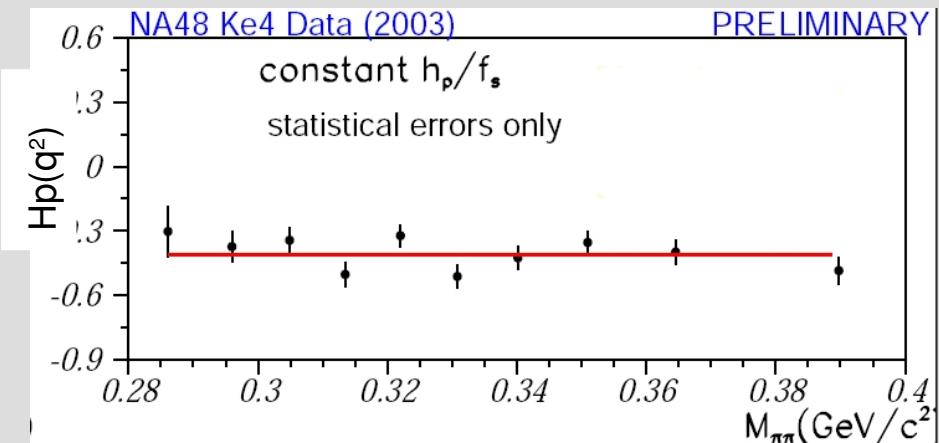
- $F_s$  is quadratic in  $q^2$



- First measurement of  $F_p \neq 0$



- Linear in  $q^2$



- No linear term ( $h_p'$ )

# $K_{e4}$ : form factors result

$$f_s/f_s = 0.165 \pm 0.011 \pm 0.006$$

$$f''_s/f_s = -0.092 \pm 0.011 \pm 0.007$$

$$f'_e/f_s = 0.081 \pm 0.011 \pm 0.008$$

$$f_p/f_s = -0.048 \pm 0.004 \pm 0.004$$

$$g_p/f_s = 0.873 \pm 0.013 \pm 0.012$$

$$g'_p/f_s = 0.081 \pm 0.022 \pm 0.014$$

$$h_p/f_s = -0.411 \pm 0.019 \pm 0.007$$

Preliminary  
(2003 data)

- Separately measured on  $K^+$  and  $K^-$  and then combined (different statistical error)

## Systematics checks:

- Acceptance
- Background
- PID
- Radiative corrections
- Evaluation of the sensitivity of the form factors on the Mev dependence of the normalization
- All the Form factors are measured relatively to  $f_s$
- first evidence of  $f_p \neq 0$  and  $f_e' \neq 0$
- The f.f. are measured at level of <5% of precision while the slopes at ~15% (factor 2 or 3 improvement wrt previous measurements)

The Ke4 decay amplitude depends on two complex phases identified with

- $\delta_0$  : the  $\pi\pi$  scattering phase shift in the  $I = 0, l = 0$  state ( $s$  – wave)
- $\delta_1$  : the  $\pi\pi$  scattering phase shift in the  $I = 1, l = 1$  state ( $p$  – wave)

The Ke4 decay rate depends on the phase shift difference  $\delta = \delta_0 - \delta_1$

$\delta$  is an increasing function of  $M_{\pi\pi}$  ;  $\delta \rightarrow 0$  for  $M_{\pi\pi} \rightarrow 2m_+$

( from scattering theory :  $\delta(k) \approx ak$  at very low centre-of-mass momentum  $k$  ;  
 $a$  is the scattering length ;  $a \neq 0$  for  $s$  – waves only )

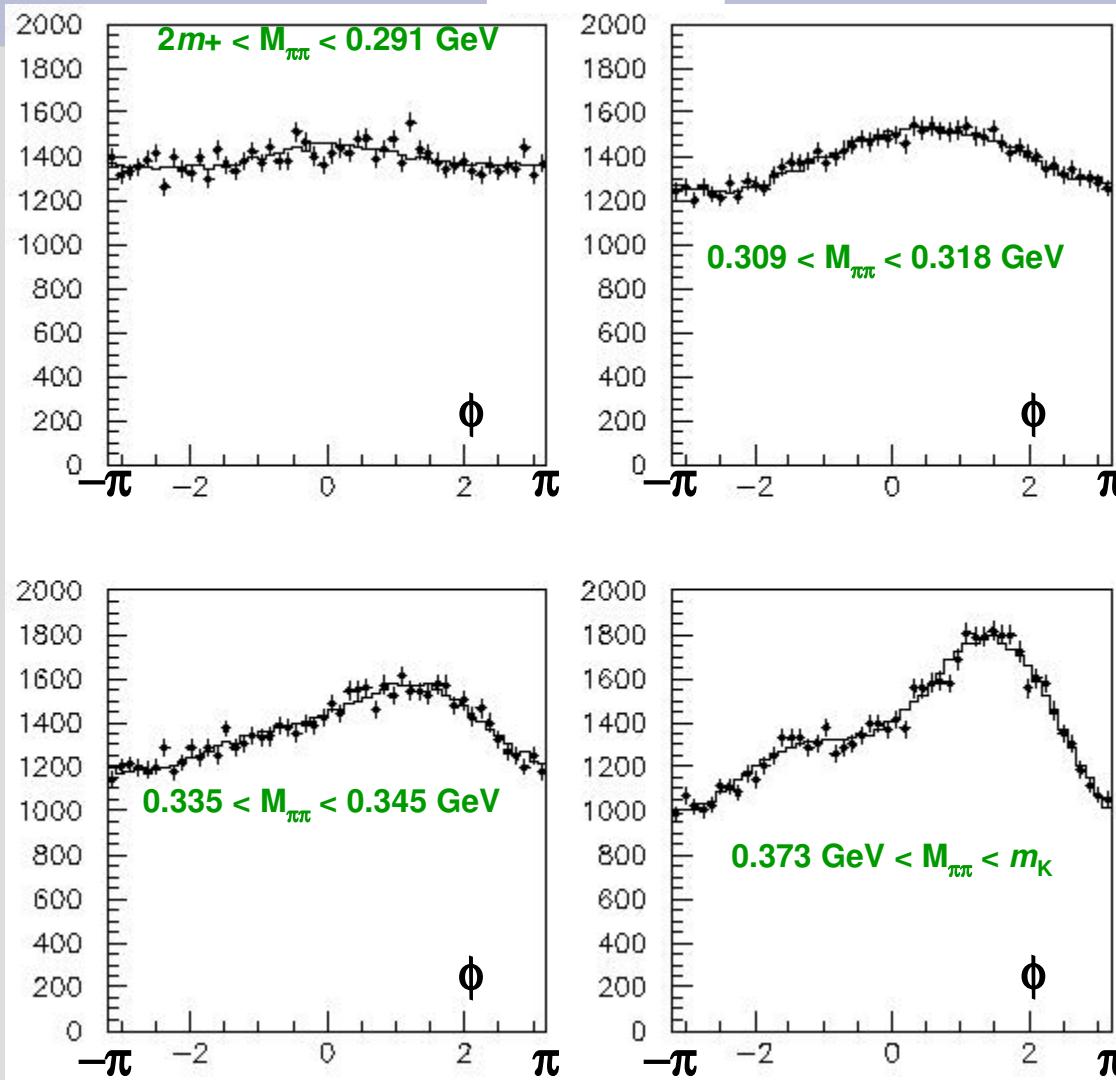
$\delta \neq 0$  in the Ke4 decay amplitude  $\rightarrow$  asymmetric distribution of  
the charged lepton direction with respect to the plane defined  
by the two pions ( $\phi$  distribution)

(Shabalin 1963)

**The asymmetry of the  $\phi$  distribution increases with  $M_{\pi\pi}$**

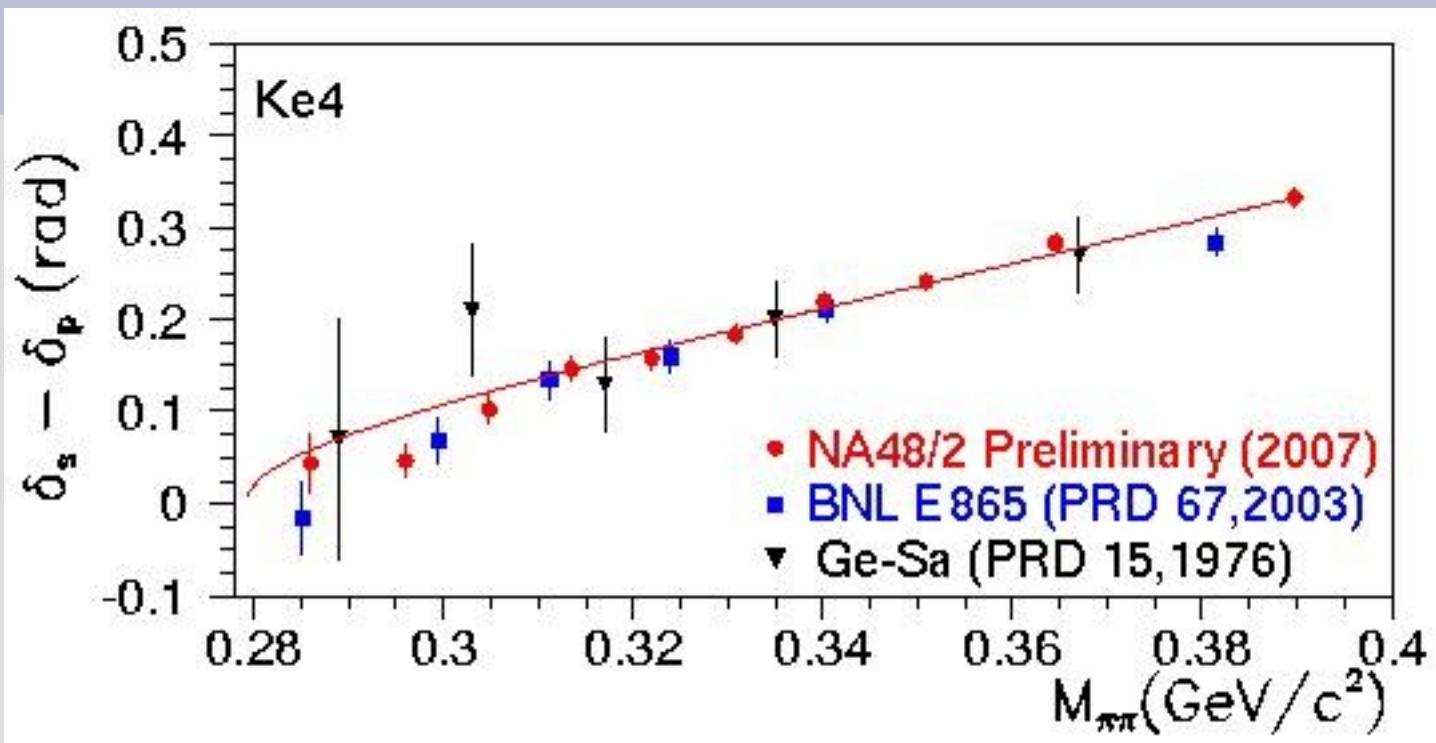
# NA48 / 2 (Preliminary)

## $\phi$ distributions for four $M_{\pi\pi}$ bins



# $\delta$ versus $M_{\pi\pi}$

(no isospin – breaking corrections)



Geneva – Saclay :  $\sim 30,000$  events ,  $p_{K^\pm} = 2.8 \text{ GeV}/c$

BNL E865 : 406,103 events (with  $\sim 4.4\%$  background),  $p_{K^\pm} = 6 \text{ GeV}/c$

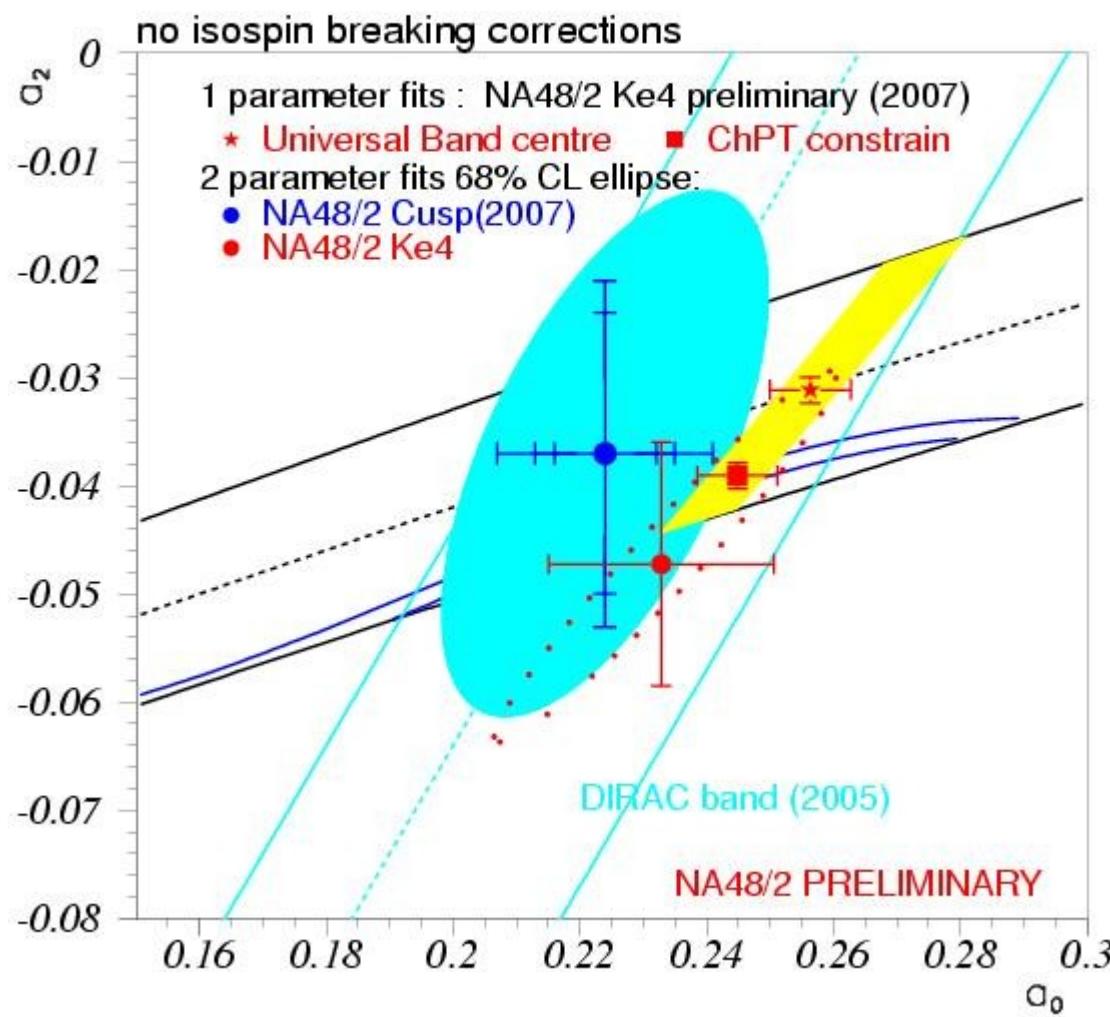
NA48/2 : 677,510 events (with  $\sim 0.5\%$  background),  $p_{K^\pm} = 60 \text{ GeV}/c$

**NOTE:** the isospin – breaking corrections reduce  $\delta$  by 0.01 – 0.012  
(J. Gasser)

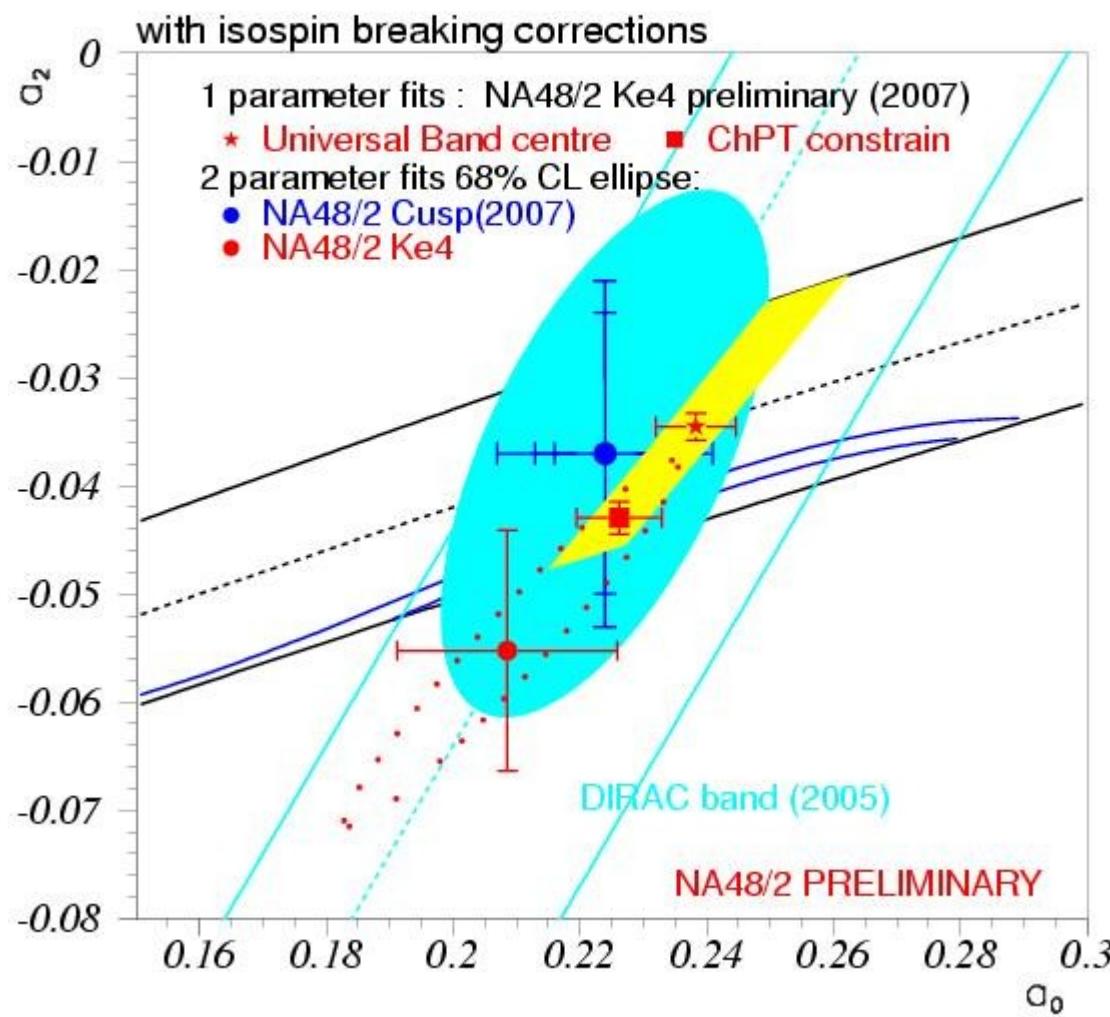
To extract scattering lengths from  $\delta = \delta_0 - \delta_1$  some external data (l=2  $\pi\pi$  scatteringg at higher energies) and theoretical input are needed, for example:

- Solution of Roy equations - ACGL, Phys.Rep.353(2001), DFGS, EPJ C24 (2002) that relates  $\delta$  and  $(a_0, a_2)$  .
- The Universal Band centre line parametrisation corresponds to 1-parameter fit with  $a_2 = f(a_0)$
- ChPT constrain gives another relation between  $(a_0, a_2)$  inside Universal Band.

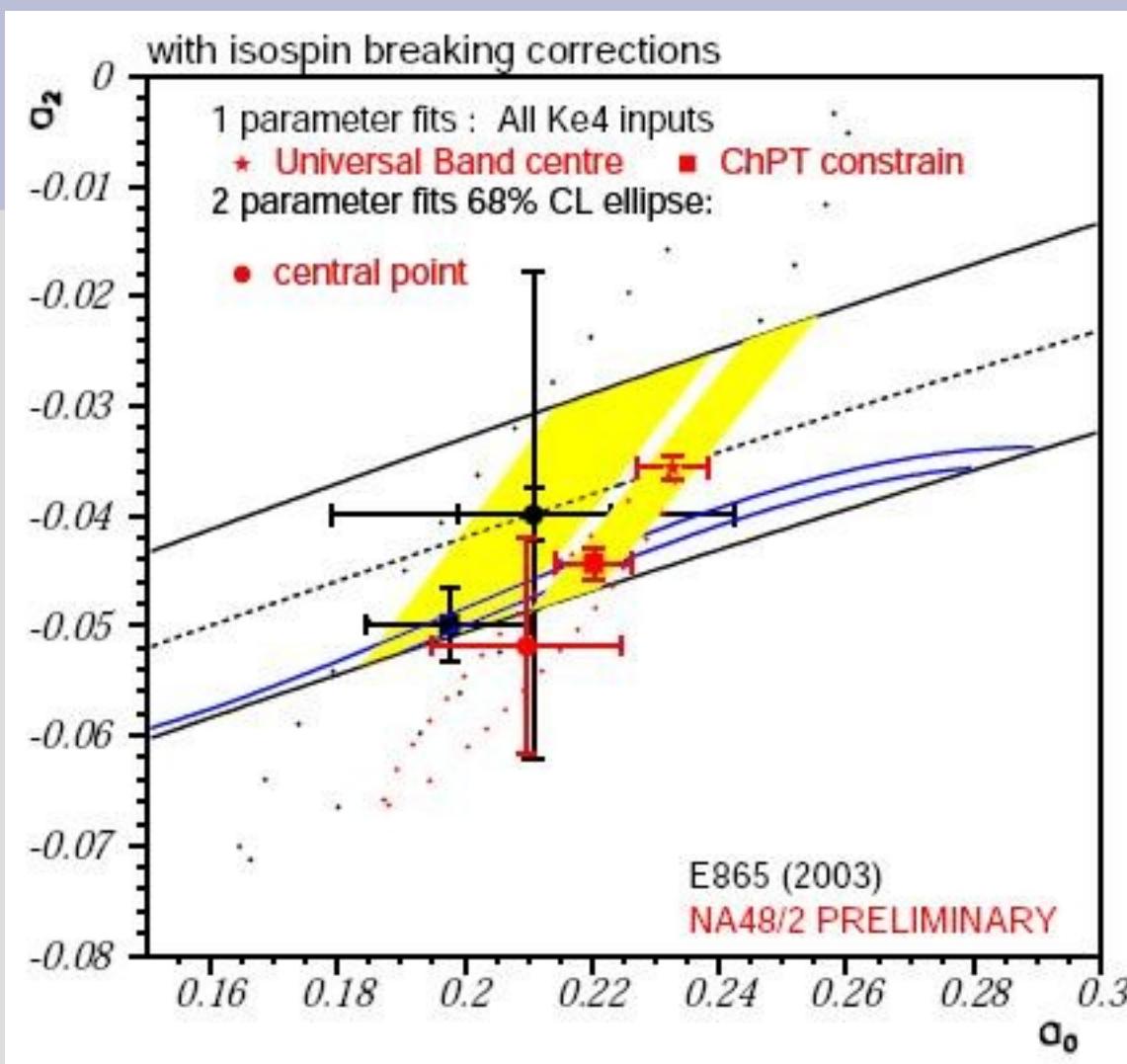
## NA48/2 Cusp – Ke4 comparison



## NA48/2 Cusp – Ke4 comparison



## Ke4: NA48/2 – BNL E865 comparison



Marginal agreement for constrained fits

# Conclusions

- A new step made in  $(a_0 - a_2)$  measurement by analysis of  $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$  decay with the Cabibbo-Isidori framework:
  - Data sample size increased by almost a factor of 4;
  - A significant improvement in experimental statistical and systematic precision on scattering lengths: in particular,  
 $(a_0 - a_2) = (0.268 \pm 0.011_{\text{exper.}})$  [2003] →  $(0.261 \pm 0.007_{\text{exper.}})$  [2003+2004]
  - However the largest uncertainty comes from the theory precision:  $\delta(a_0 - a_2) = 0.013$  → a challenge to the theory!
- Ke4
  - Expect further improvement from NA48/2 by the addition of the 2004 data ( $\Rightarrow$  total number of events  $> 10^6$ ).
  - A new development: isospin – breaking corrections. Better agreement between NA48/2 Ke4 and cusp.
  - Agreement between NA48/2 and E865 is marginal and should be understood