

# Electron and Photon Interactions at Very High Energies

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# Overview

- Overall introduction
- Higgs sector
- Supersymmetry
- Heavy Majorana neutrinos
- Inverse neutrinoless double beta decay
- Photon collider — general attractions
- Compton polarimetry at SLC
- Deeply virtual Compton scattering
- Radiative Corrections and polarization effects
- Compatibility of  $e^+e^-$  and  $e^-e^-$  interaction regions
- Specific machine aspects (CLIC)

# Overall Introduction

## Definitions

- Moeller scattering:  $e^-e^- \Pi \dots$
- Bhabha scattering:  $e^+e^- \Pi \dots$
- Compton scattering:  $\gamma C \Pi \dots$  ( $C =$  charged particle)

# Additive Quantum Numbers of Different $e^-e^-$ initial states

Table 1. Additive quantum numbers of different-helicity  $e^-e^-$  initial states: judicious choices of polarization parameters can tune the final-state representations, fix chiral couplings of the incoming electrons.

	$Q_{el}$	$S_Z$	$L$	$L_e$	$I_3^W$	$Y_W$
$e_L^-e_L^-$	-2	1	2	2	-1	-2
$e_L^-e_R^-$	-2	0	2	2	-1/2	-3
$e_R^-e_R^-$	-2	-1	2	2	0	-4

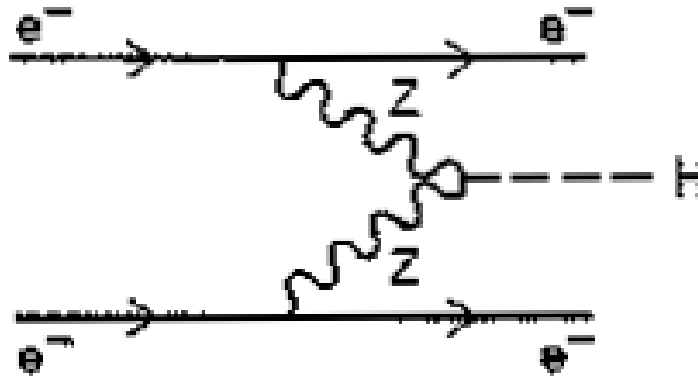
# Advantages of using $e^-e^-$ interactions

- $e^-$  can be highly polarized ( $e^+$  cannot!)
- Easily available
- Clearly described (no form factor)

Clean description, clean experimentation

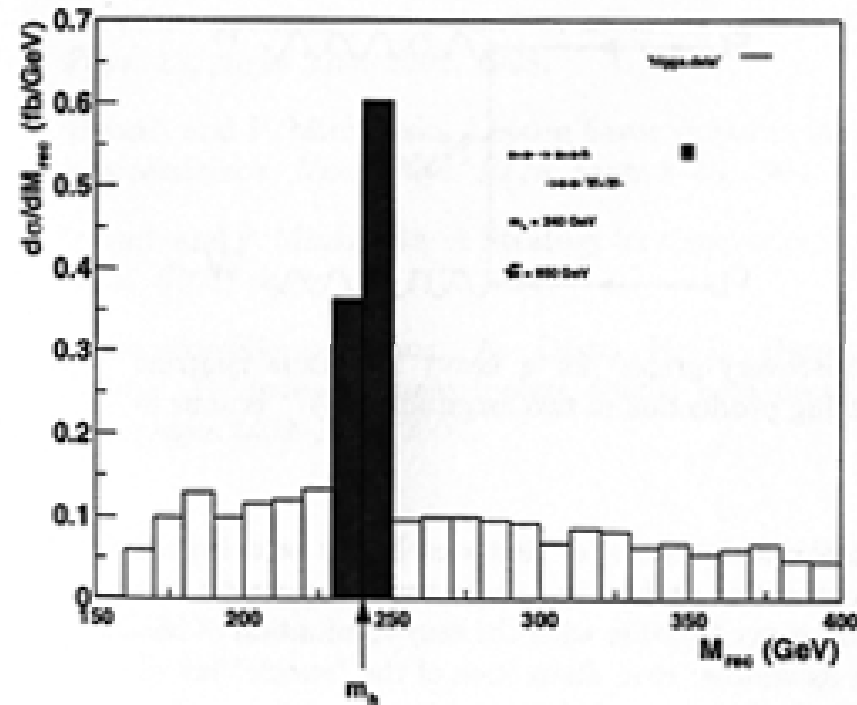
# Higgs Sector

# The Clear Signal for H Production in $e^-e^-$ Interactions



Clean, clear signal for H production!

# Differential cross section of standard $H^0$ scalar production



Differential cross section of standard  $H^0$  scalar production for  $m_h = 240$  GeV — at c.m. energy of 850 GeV, and the background reaction  $e^-e^- \rightarrow e^-e^- W^+W^-$ , with respect to the recoil mass.



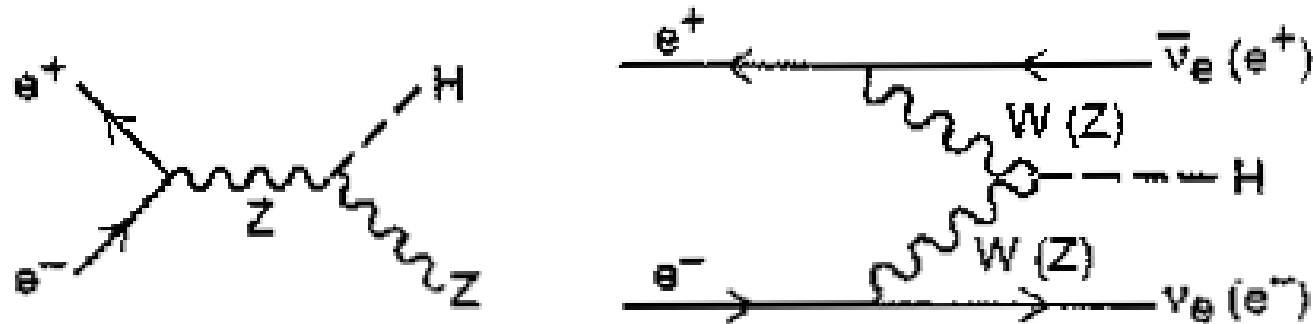
# Examples for Outstanding Possibilities of $e^-e^-$ Collisions

- In the context of this Lomonosov conference, it will be useful to consider a few specific examples of the capabilities of using highly polarized electron beams as the principal input to experimentation. While there is no need to be exhaustive, I will give a few outstanding interactions that delineate the strengths of this particular  $e^-e^-$  state.
- Altogether, we cannot find a better defined initial state than two monoenergetic pointlike electrons of chosen spin orientation.
- The ILC will be the first machine where we find the option of colliding  $e^+e^+$ ,  $e^+e^-$ , or  $e^-e^-$  at full energy as well as, at slightly reduced energy,  $\gamma\gamma$  and  $e^-\gamma$ , but still at full luminosity.

Let us look at a few illustrative examples.

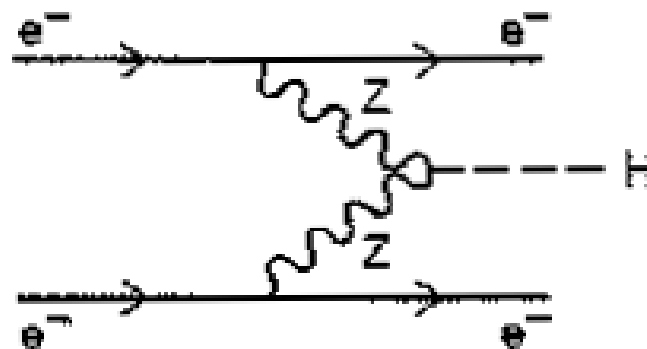
## A Study of the Higgs Sector

The standard production process of the neutral Higgs boson is by  $e^+e^-$  annihilation into a virtual  $Z$  boson, followed by "Higgsstrahlung".



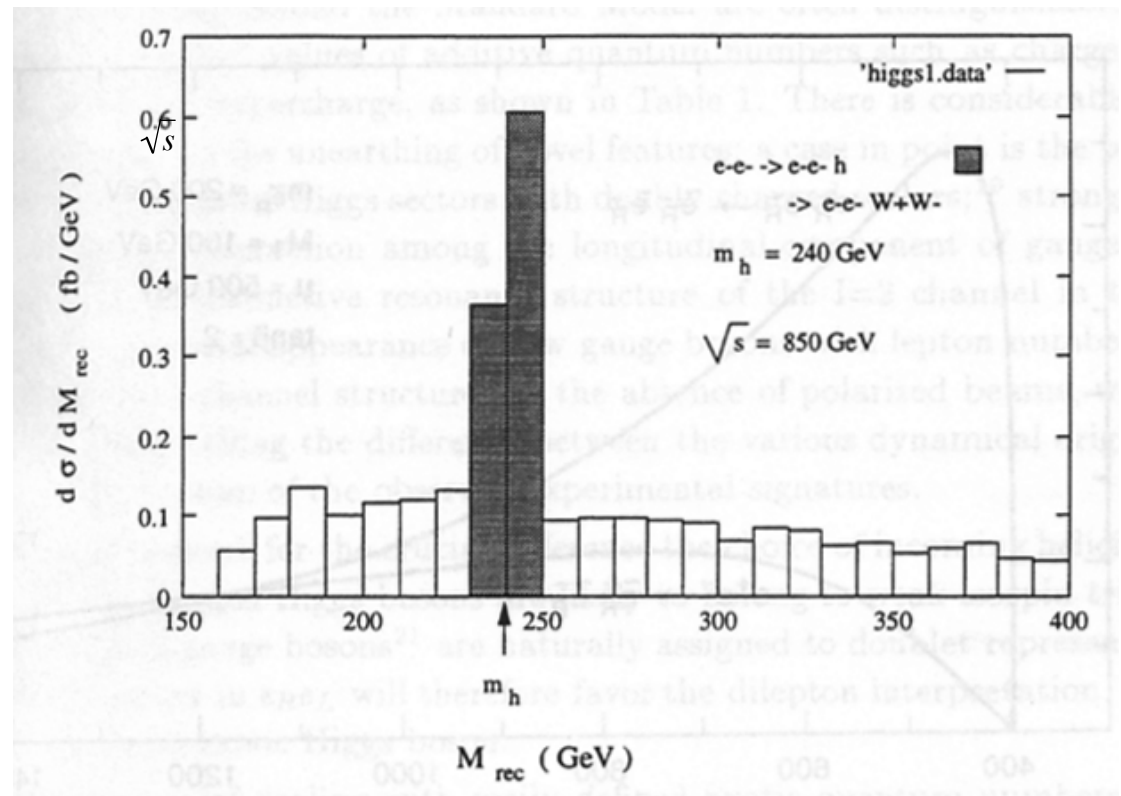
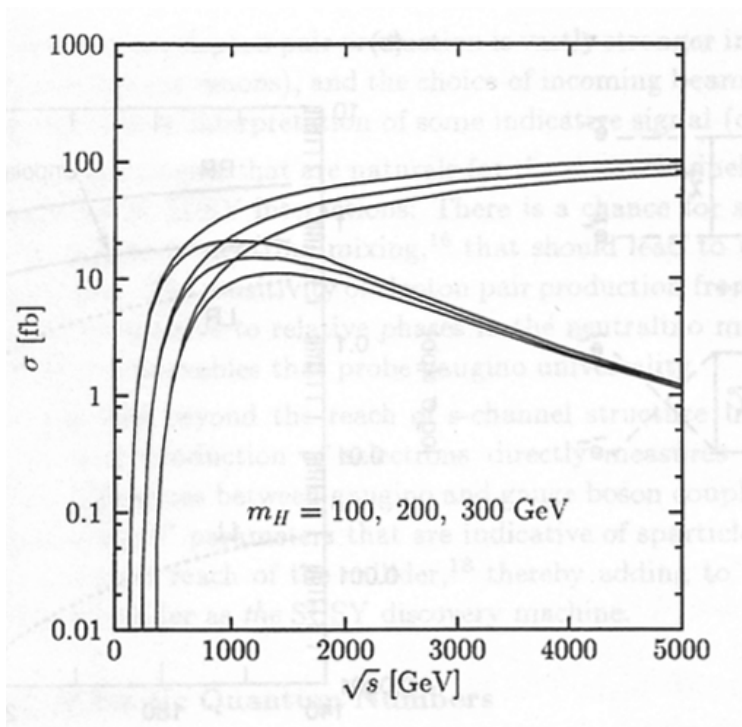
Alternatively  $e^+$  and  $e^-$  emit virtual  $W$ 's or  $Z$ 's that fuse into an  $H^0$ , where the kinematics are not fully measurable because of the emitted neutrinos.

The preferable mode of Higgs boson production in  $e^+e^-$  interactions, where the final-state electrons define the Higgs mass fully, is



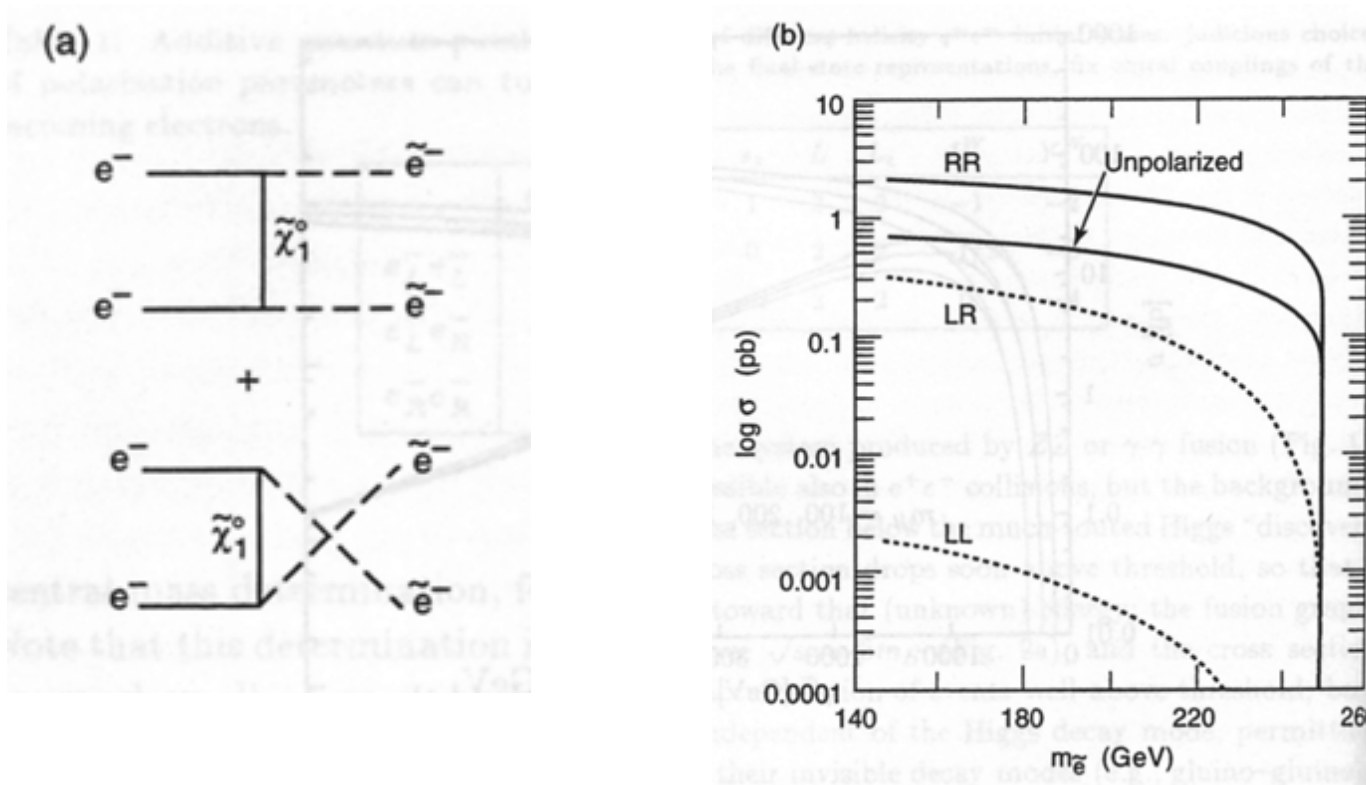
# Central Production Cross Section for Higgs Bosons

Rises steeply and saturates above  $3m_H$



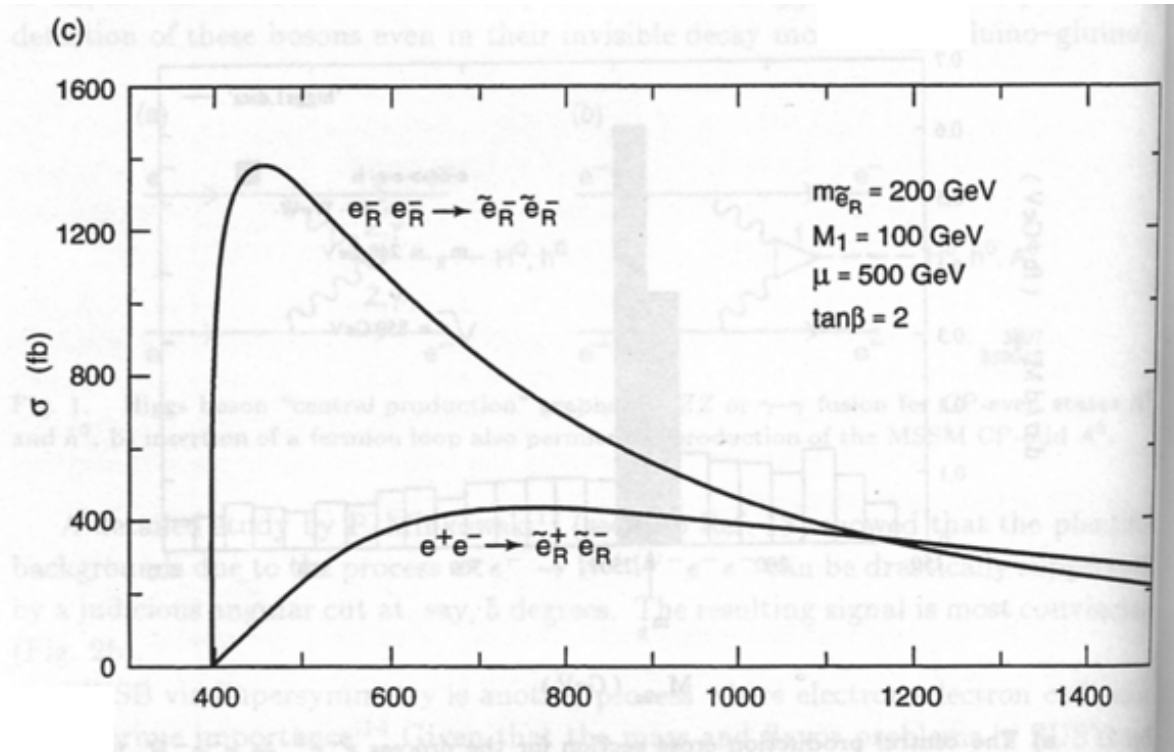
# Supersymmetry

# Supersymmetry Applications



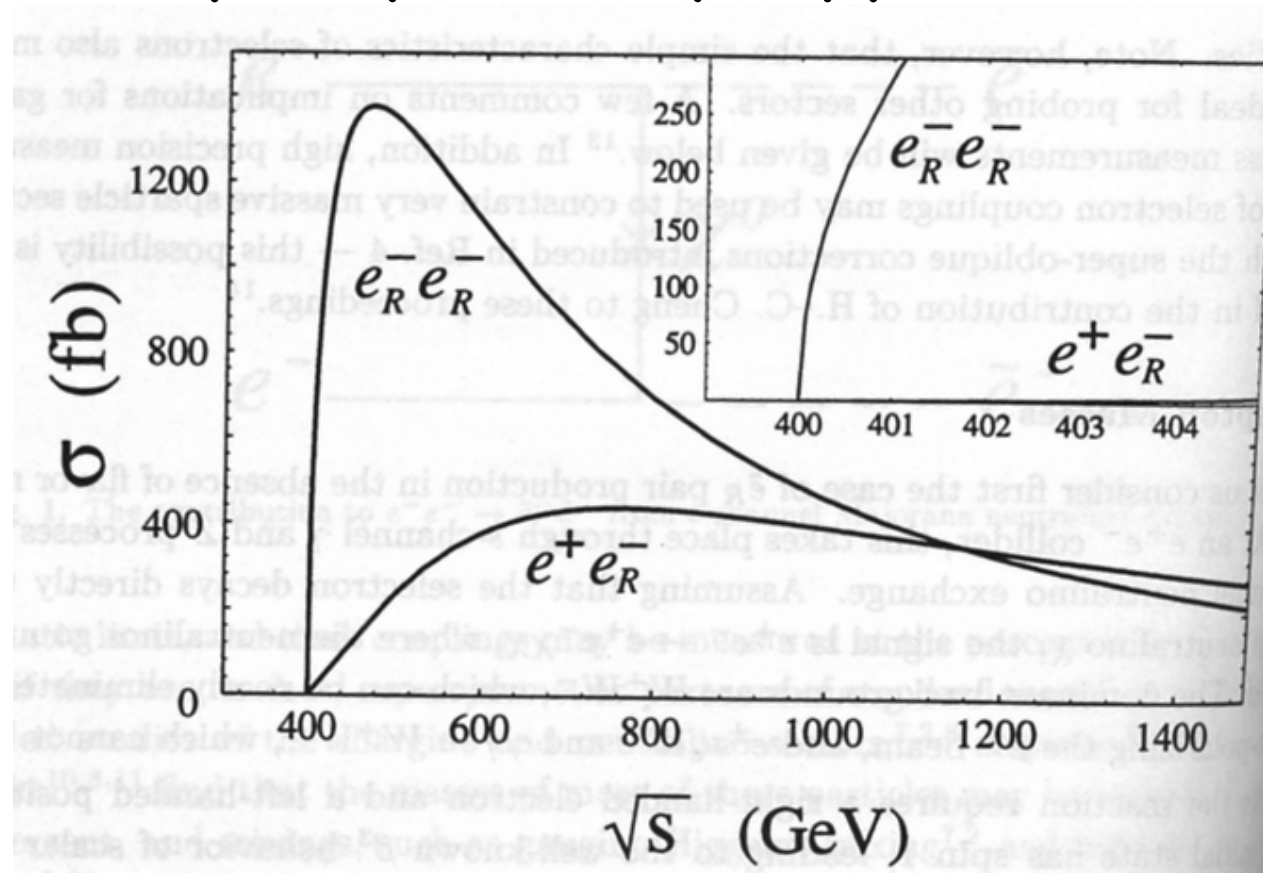
Selectron pair production in  $e^-e^-$  collisions. a) Basic graphs with neutralino exchange. b) The cross-section as a function of the selectron mass shows a strong polarization dependence. The  $e_R e_R$  mode is vastly preferred, backgrounds can be identified by  $e_L e_L$  operation.

# Supersymmetry Applications



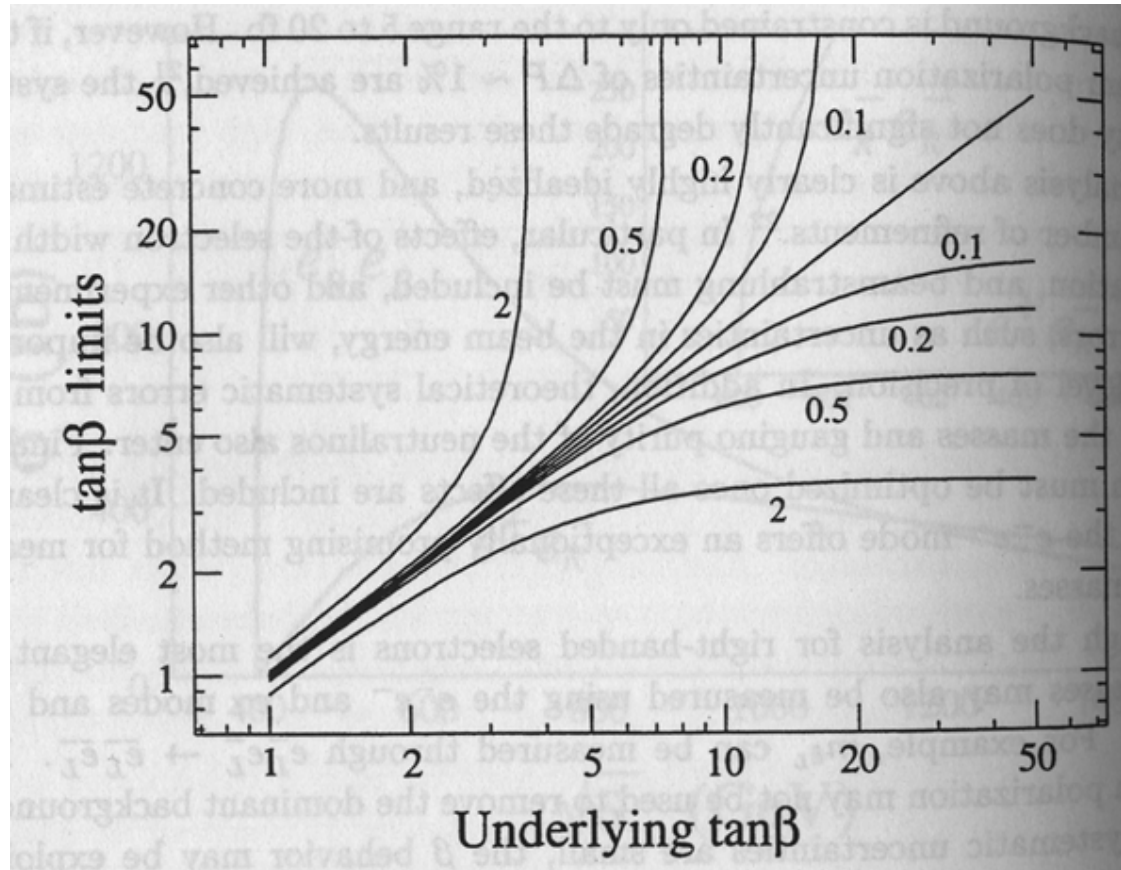
Selectron pair production in  $e^-e^-$  collisions. c) The sharp turn-on at threshold show why  $e^-e^-$  collisions are the natural choice for slepton pair production; mass determination is at least an order of magnitude better than in  $e^+e^-$  annihilation.

# Supersymmetry Applications



Cross sections for selectron pair production. Note the huge advantage of  $e^-e^-$  vs.  $e^+e^-$ .

# Supersymmetry Applications

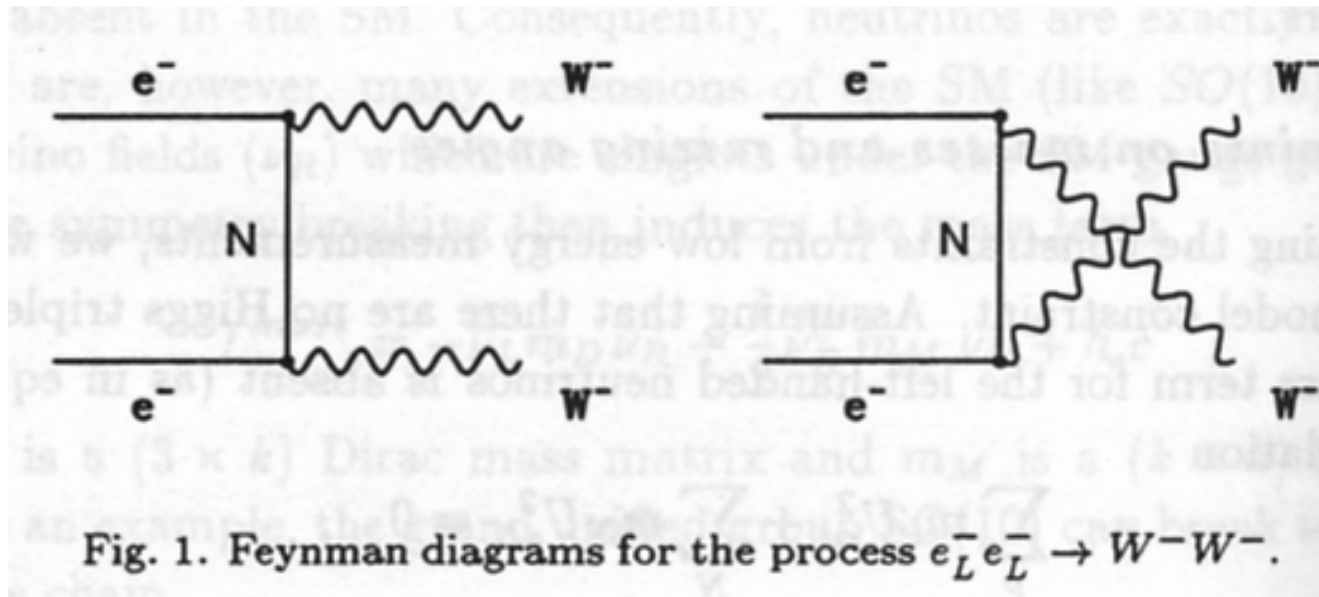


Measurement of  $\tan\beta$ , which has important implications for Yukawa couplings and other SUSY measurements.

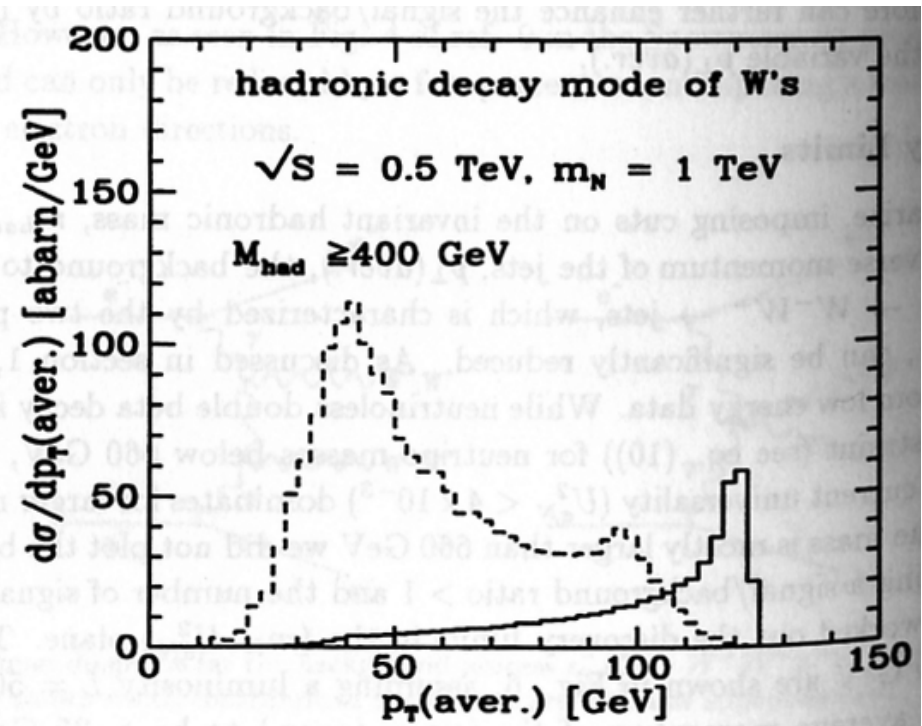
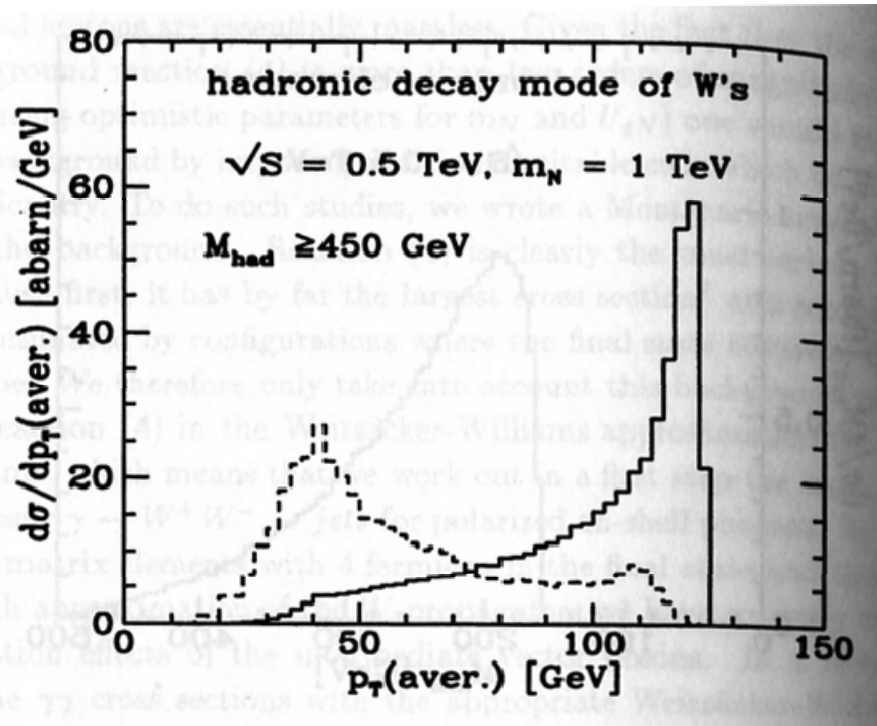


Heavy Majorana neutrinos

# Feynman diagrams for heavy Majorana neutrino exchange

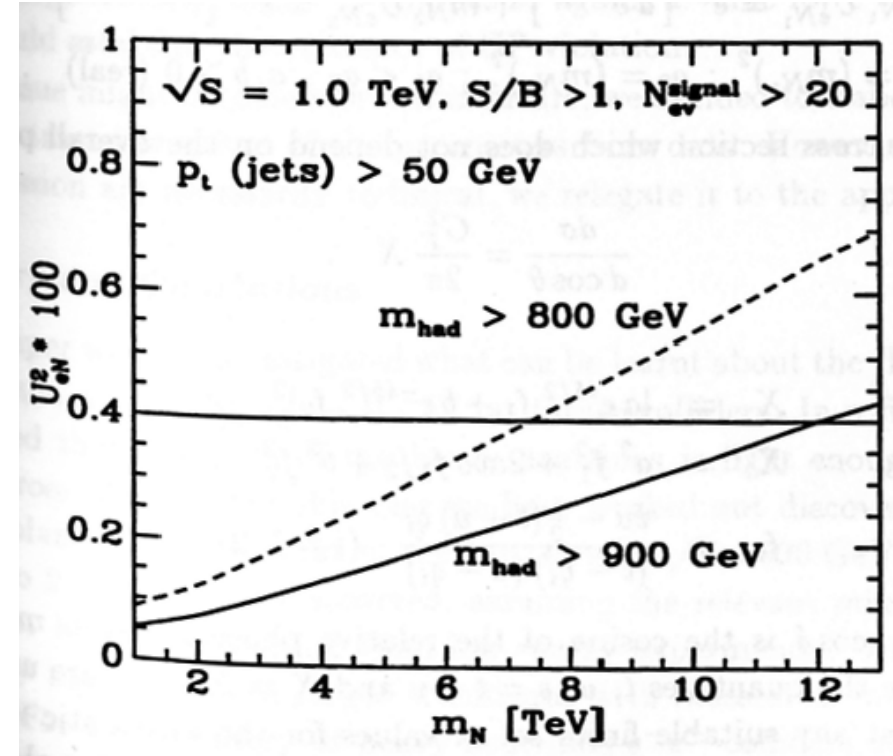
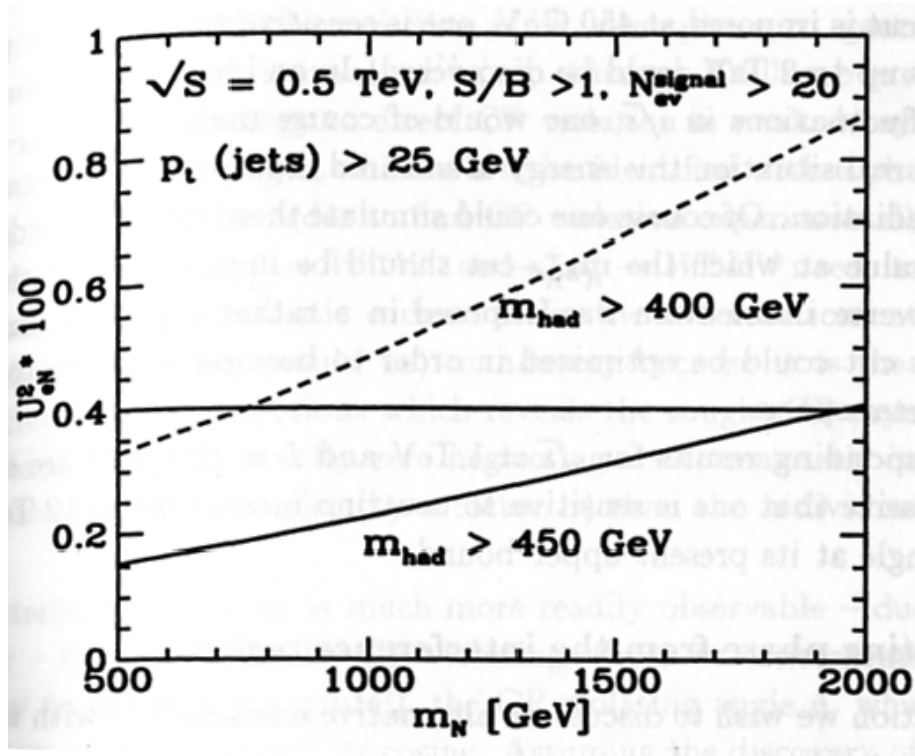


# Heavy Majorana neutrinos



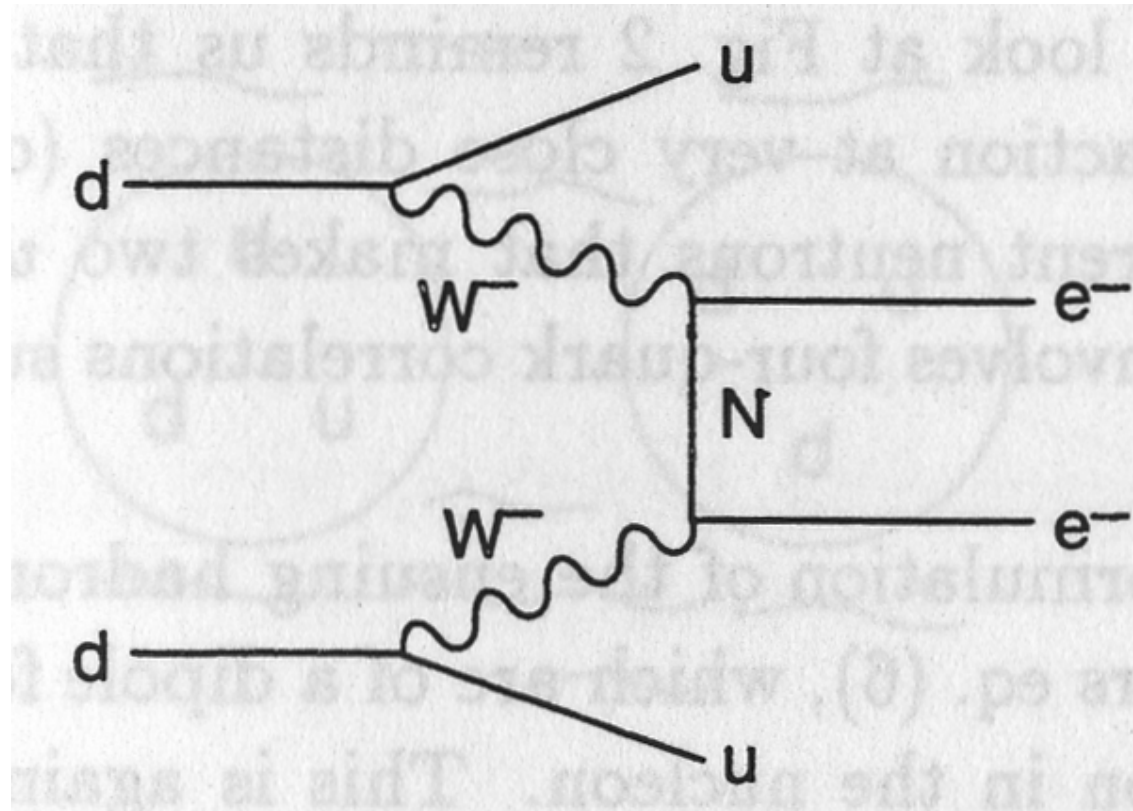
Discovery limits for heavy Majorana neutrinos

# Heavy Majorana neutrinos



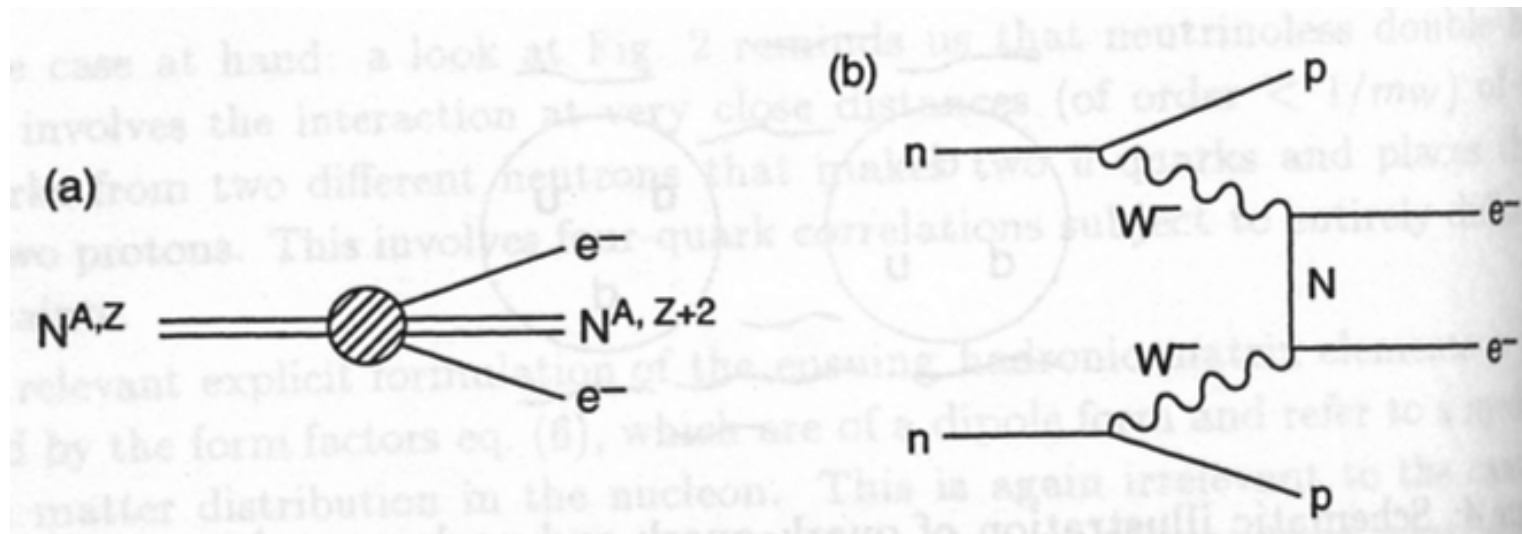
Inverse neutrinoless double beta decay

## Neutrinoless Double Beta Decay



The quark-level subprocess in neutrinoless double beta decay mediated by a Majorana neutrino  $N$ .

## Neutrinoless Double Beta Decay

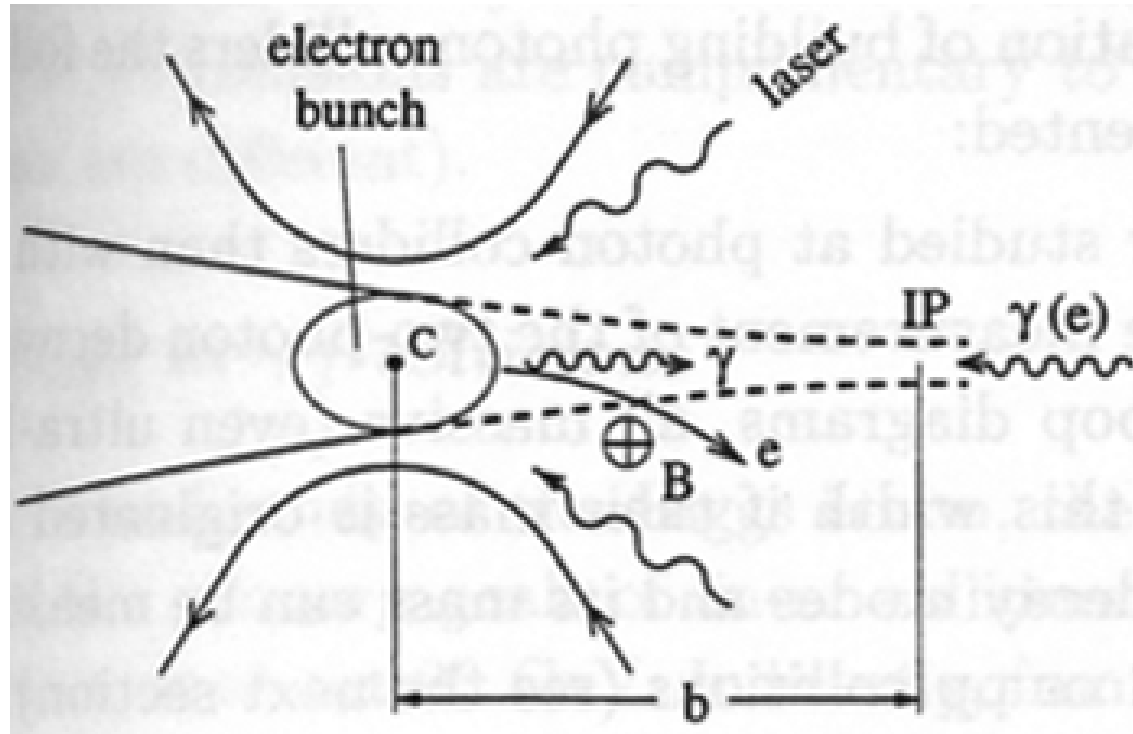


The constraints on the quark-level diagram for neutrinoless double beta decay process shown in previous slide are imposed on the nuclear (a) and nucleon (b) level.

Photon collider — general attractions



# Photon Colliders

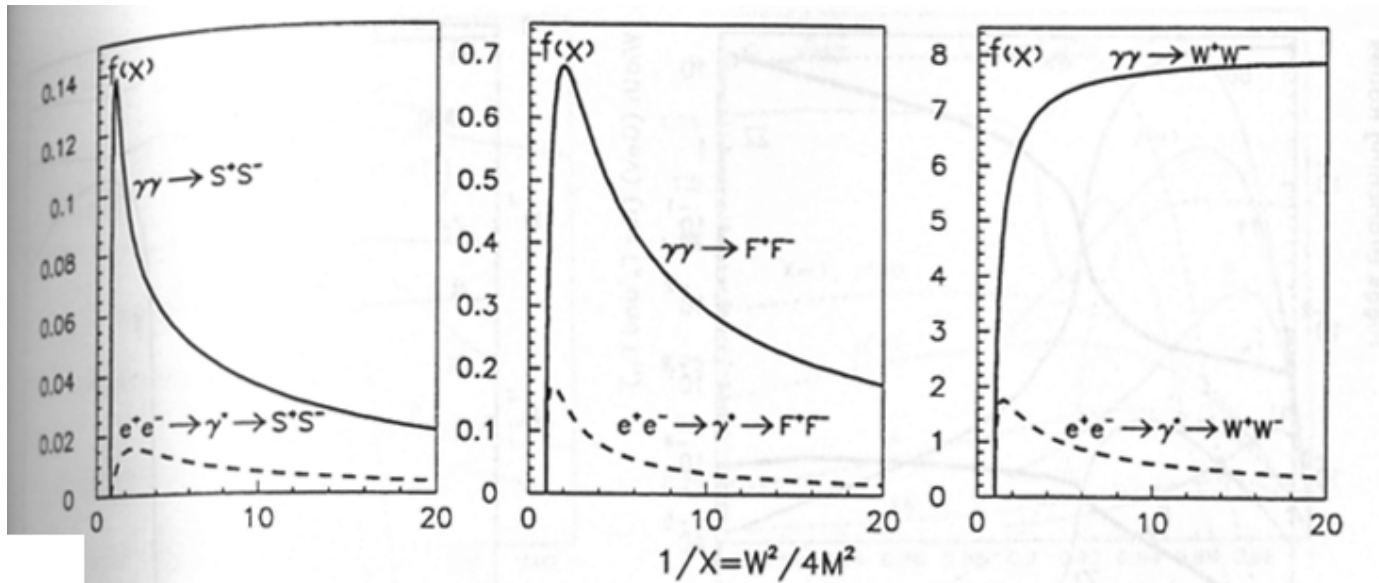


Scheme for  $\gamma\gamma, \gamma e$  collisions

# Photon Colliders

## $\gamma\gamma$ and $\gamma e$ Collisions

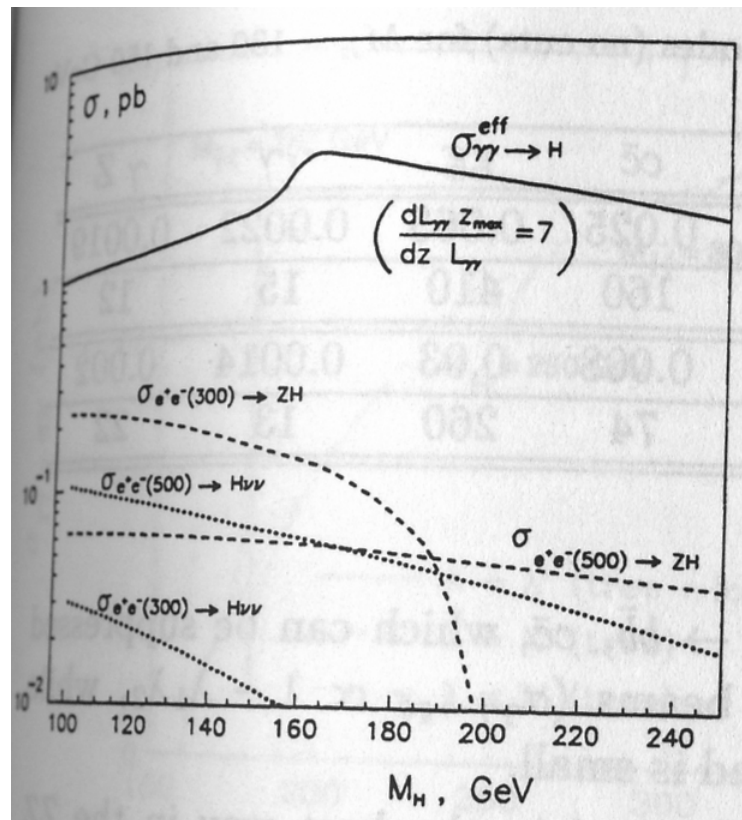
The special usefulness of  $\gamma\gamma$  collisions — no overall charge, no polarization



Comparison of cross sections for charged pair production in  $e^+e^-$  and  $\gamma\gamma$  collisions. The cross section  $\sigma = (\pi\alpha^2/M^2)f(x)$ ,  $P=S$  (scalars),  $F$  (fermions),  $W$  (W-bosons);  $M$  is particle mass,  $x = W_{pp}^2/4M^2$ . The functions  $f(x)$  are shown.

# Photon Colliders

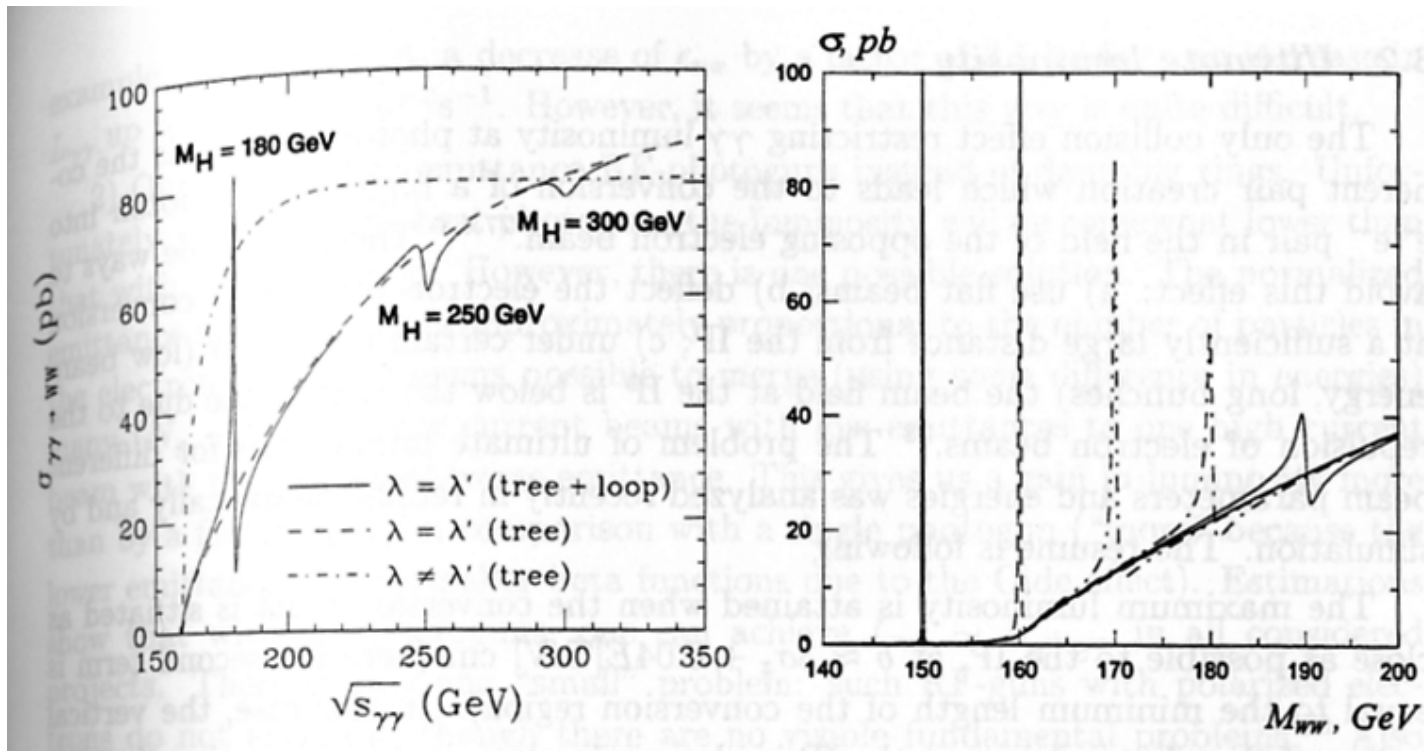
## Higgs production in $\gamma\gamma$ collisions



Comparison of Standard Model Higgs production in  $\gamma\gamma$  and  $e^+e^-$  collisions.

# Photon Colliders

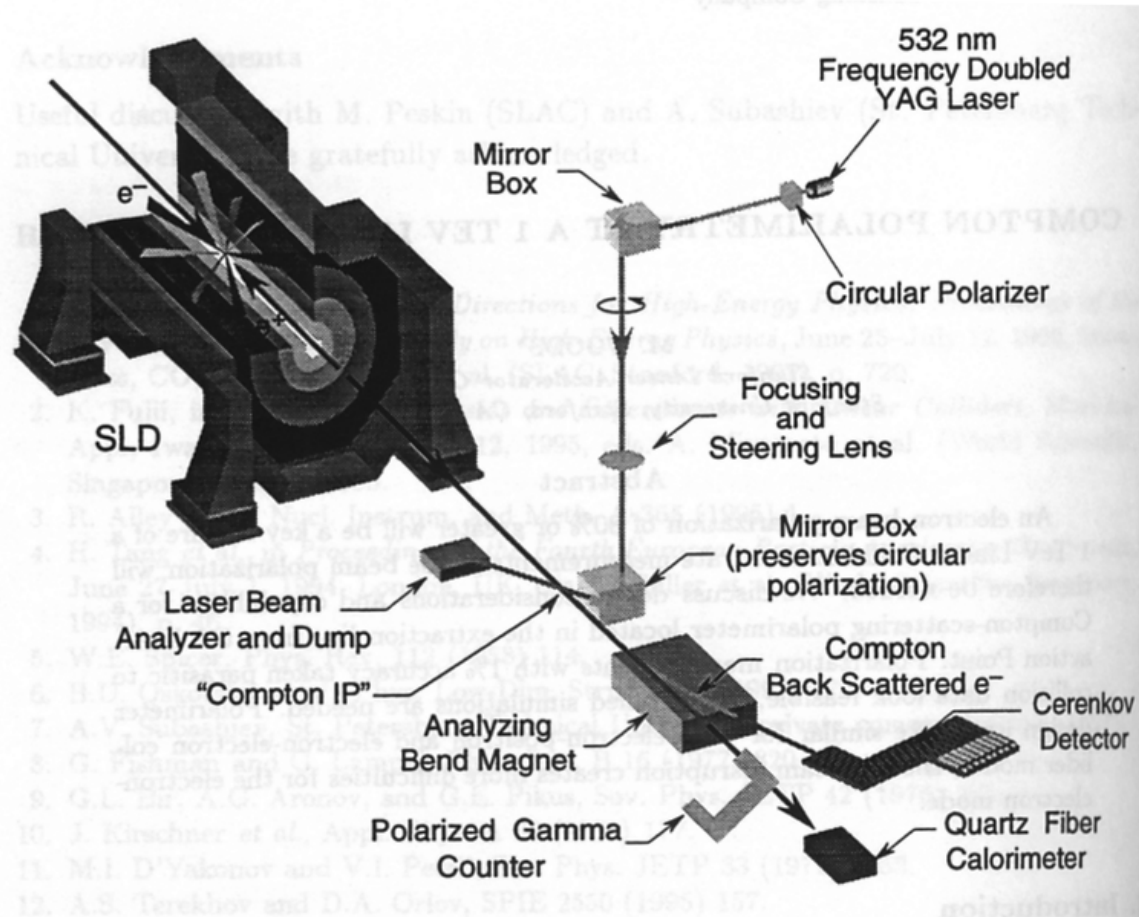
## Producing WW in $\gamma\gamma$ collisions



The  $\gamma\gamma \rightarrow WW$  cross section for various Higgs masses. The influence of Higgs masses on WW production.

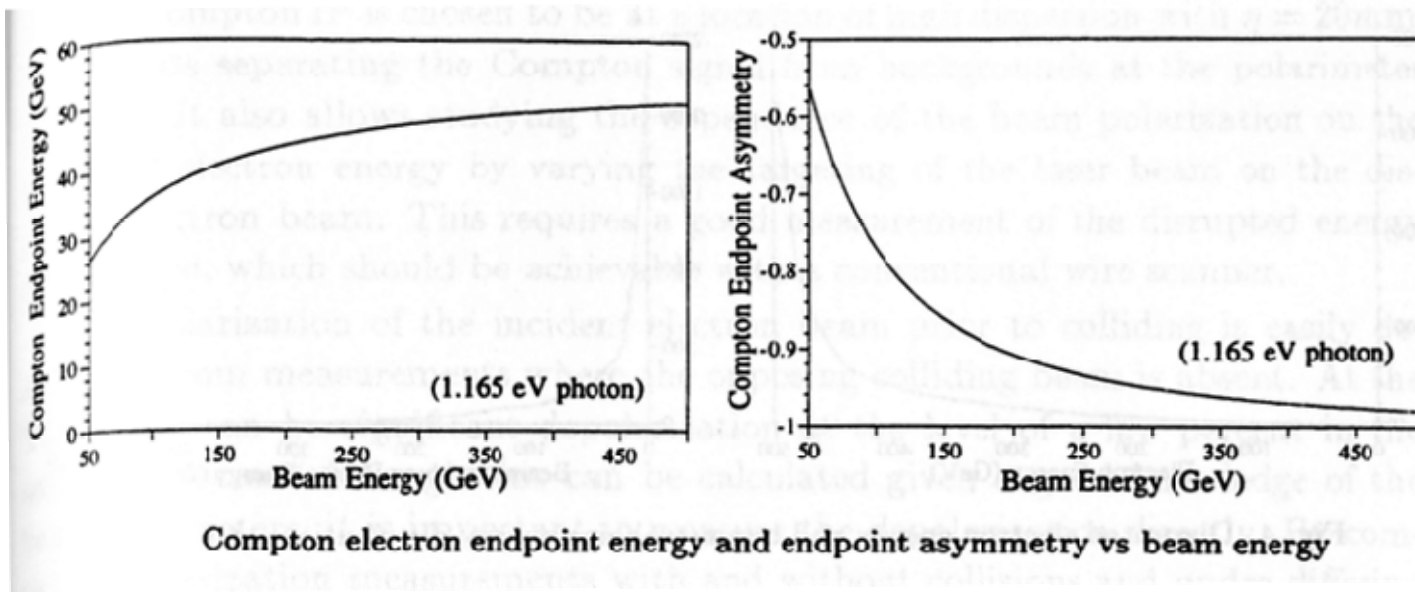
# Compton polarimetry at SLC

# Compton Polarimetry

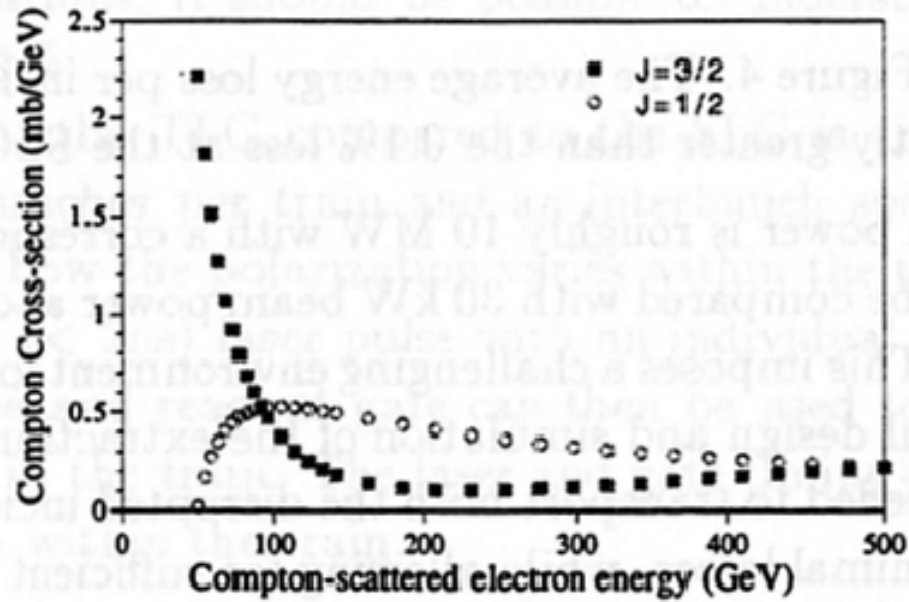


Compton polarimeter for the 1 TeV Linear Collider being tried out at the SLD experiment.

# Compton Polarimetry



# Compton Polarimetry

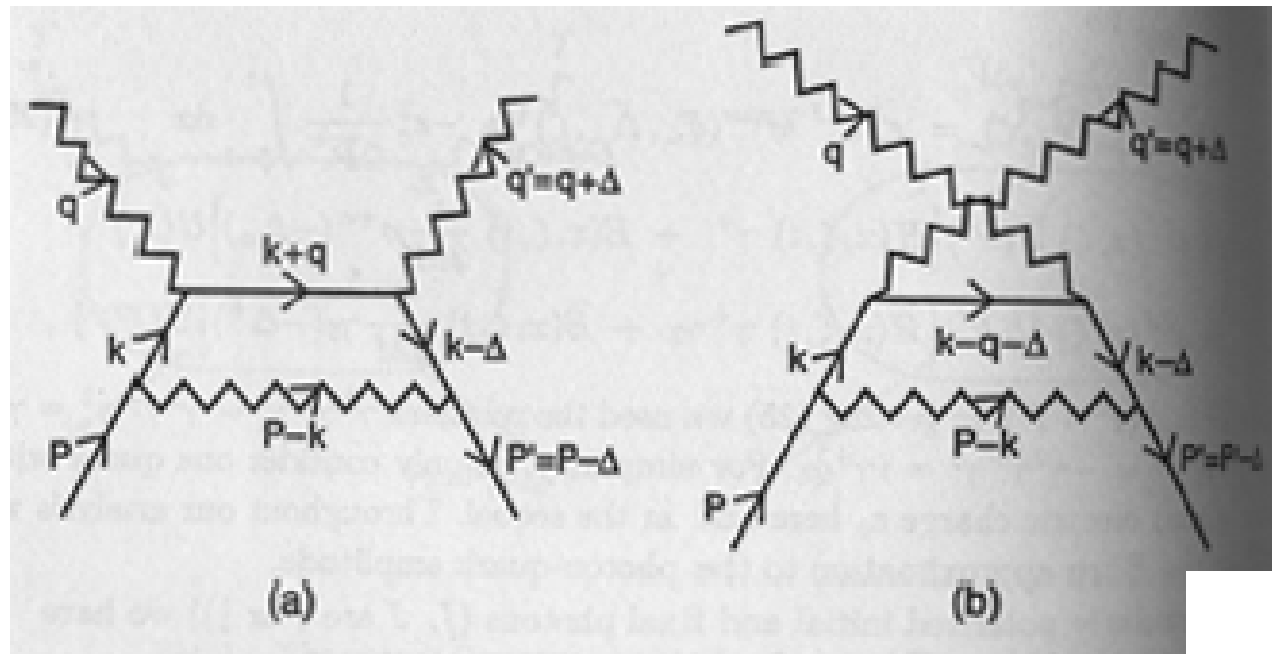


Compton scattering cross-sections for the  $J=3/2$  and  $J=1/2$  polarization states



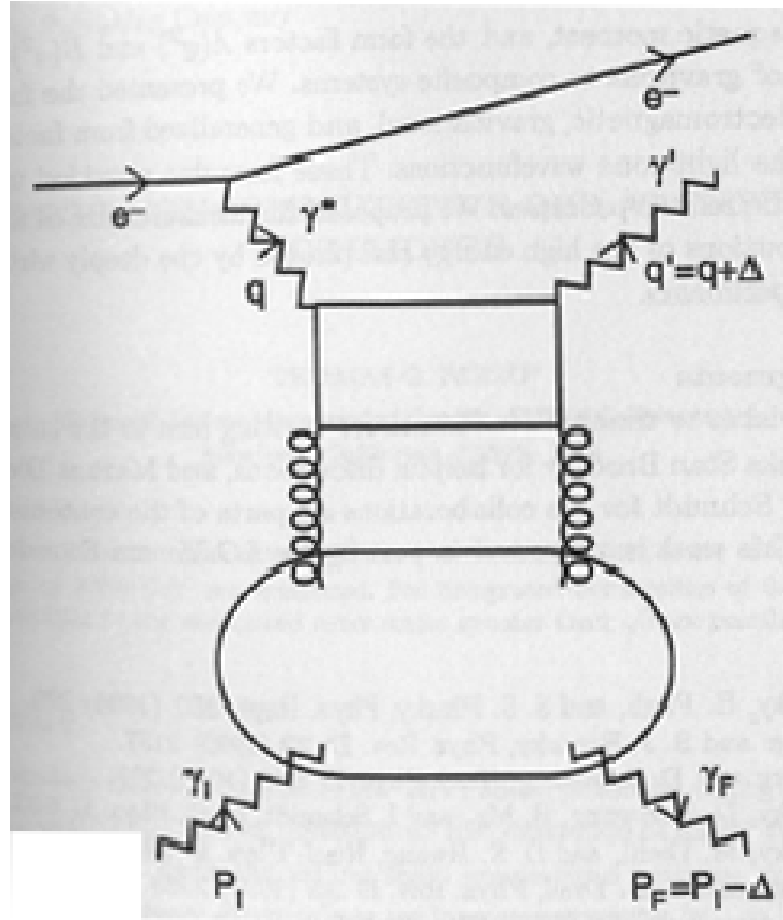
# Deeply virtual Compton scattering

# Deeply virtual Compton scattering



One-loop covariant Feynman diagrams for virtual Compton scattering.

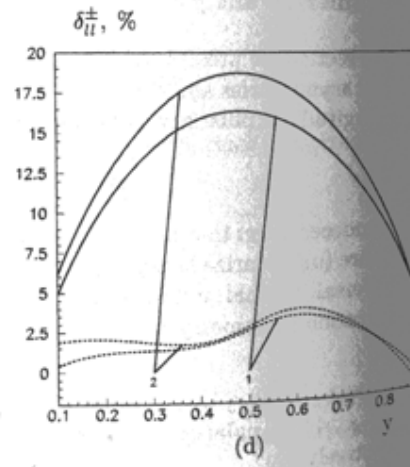
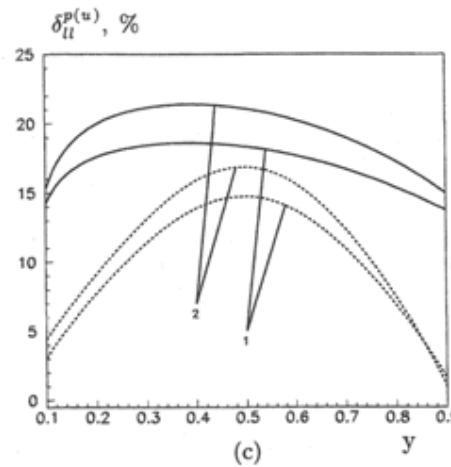
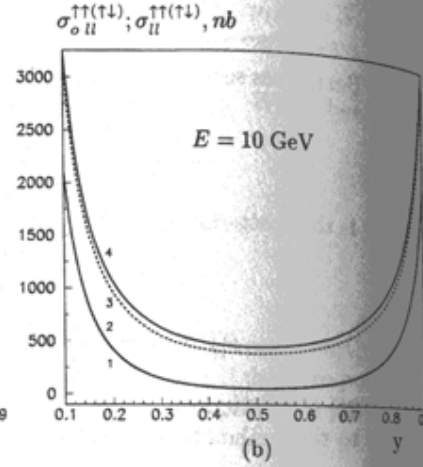
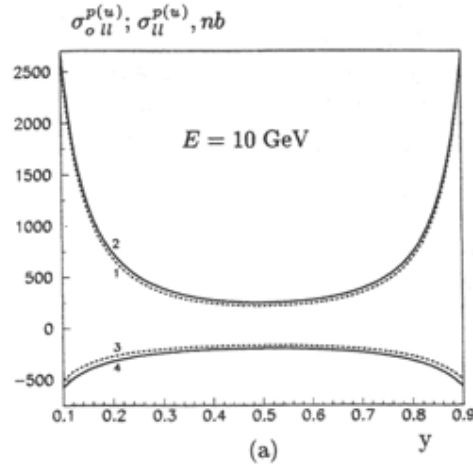
# Deeply virtual Compton scattering



The virtual Compton amplitude  $\gamma^*(q) + \gamma_I(P) \Pi \gamma(q') + \gamma_F(P')$

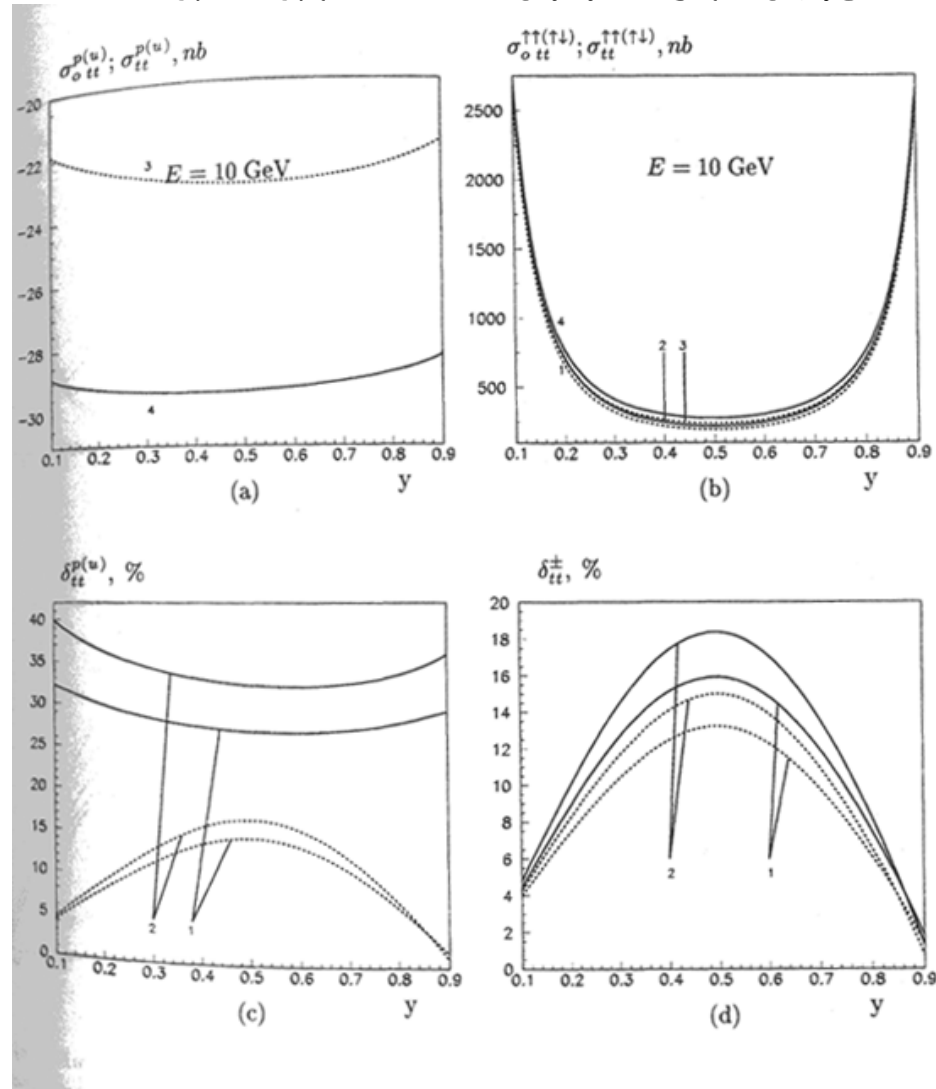
# Radiative Corrections and Polarization Effects

# Radiative Corrections



$y$ -dependence of  $e^-e^-$ -scattering cross section and corrections to it.  $S=2mE$ . (a):  $\sigma_o^u(1)$ ,  $\sigma^u(2)$ ,  $\sigma_{o\ell\ell}^p(3)$ ,  $\sigma_{\ell\ell}^p(4)$ ; (b):  $\sigma_{o\ell\ell}^{\uparrow\uparrow}(1) \simeq \sigma_{\ell\ell}^{\uparrow\uparrow}(2)$ ,  $\sigma_{o\ell\ell}^{\uparrow\downarrow}(3)$ ,  $\sigma_{\ell\ell}^{\uparrow\downarrow}(4)$ ; (c):  $\delta_{\ell\ell}^{p(u)}$  -solid(dashed) line; (d):  $\delta_{\ell\ell}^+$  -dashed,  $\delta_{\ell\ell}^-$  -solid line.  $E=10(1), 100(2) \text{ GeV}$ .

# Radiative Corrections

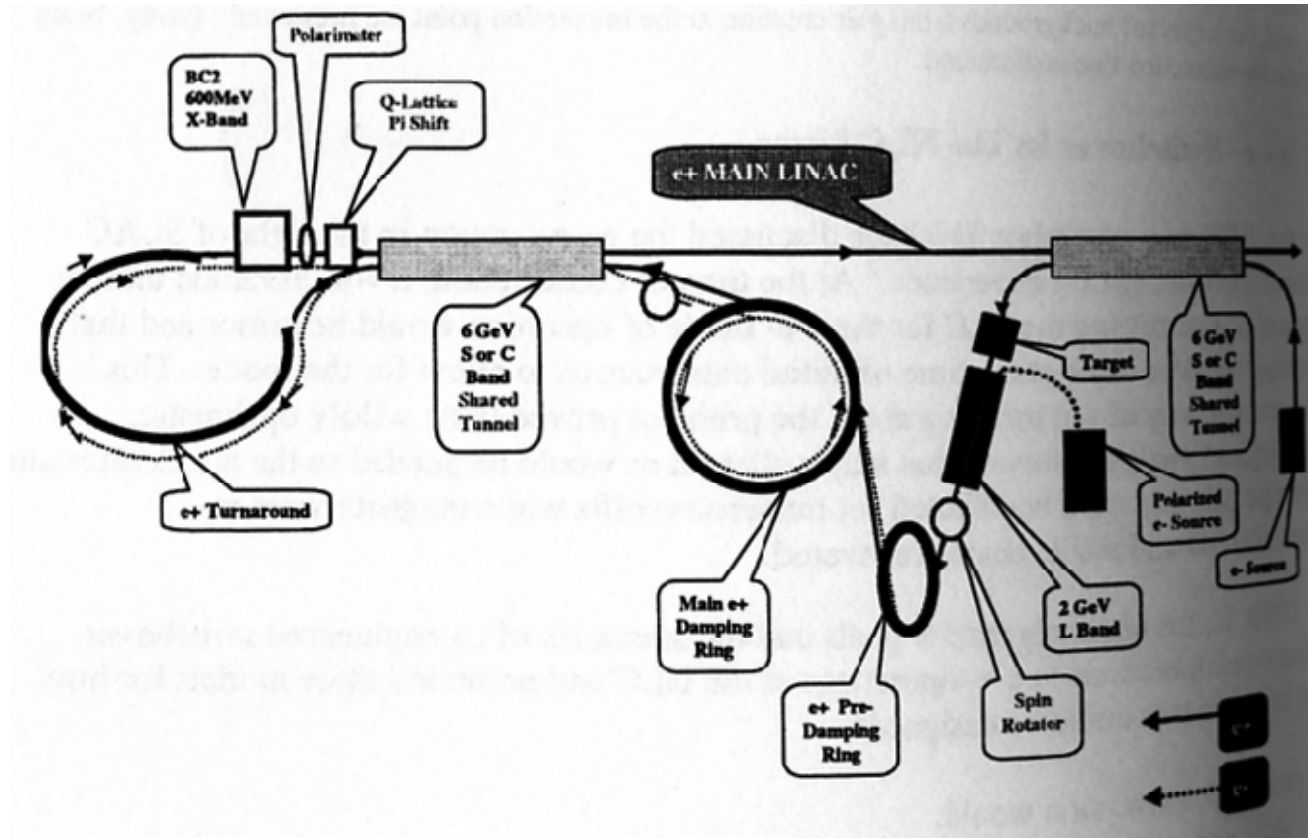


$y$ -dependence of  $e^+e^-$ -scattering cross section and corrections to it.  $S=2mE$ . (a):  $\sigma_{0\,tt}^p(3), \sigma_{tt}^p(4)$ ; (b):  $\sigma_{0\,tt}^{\uparrow\uparrow}(1), \sigma_{tt}^{\uparrow\uparrow}(2) \simeq \sigma_{0\,tt}^{\uparrow\downarrow}(3), \sigma_{tt}^{\uparrow\downarrow}(4)$ ; (c):  $\delta_{tt}^{p(u)}$ -solid(dashed) line; (d):  $\delta_{tt}^+$ -dashed,  $\delta_{tt}^-$ -solid line.  $E=10(1), 100(2)$  GeV.

# Compatibility of $e^+e^-$ and $e^-e^-$ Interaction Regions

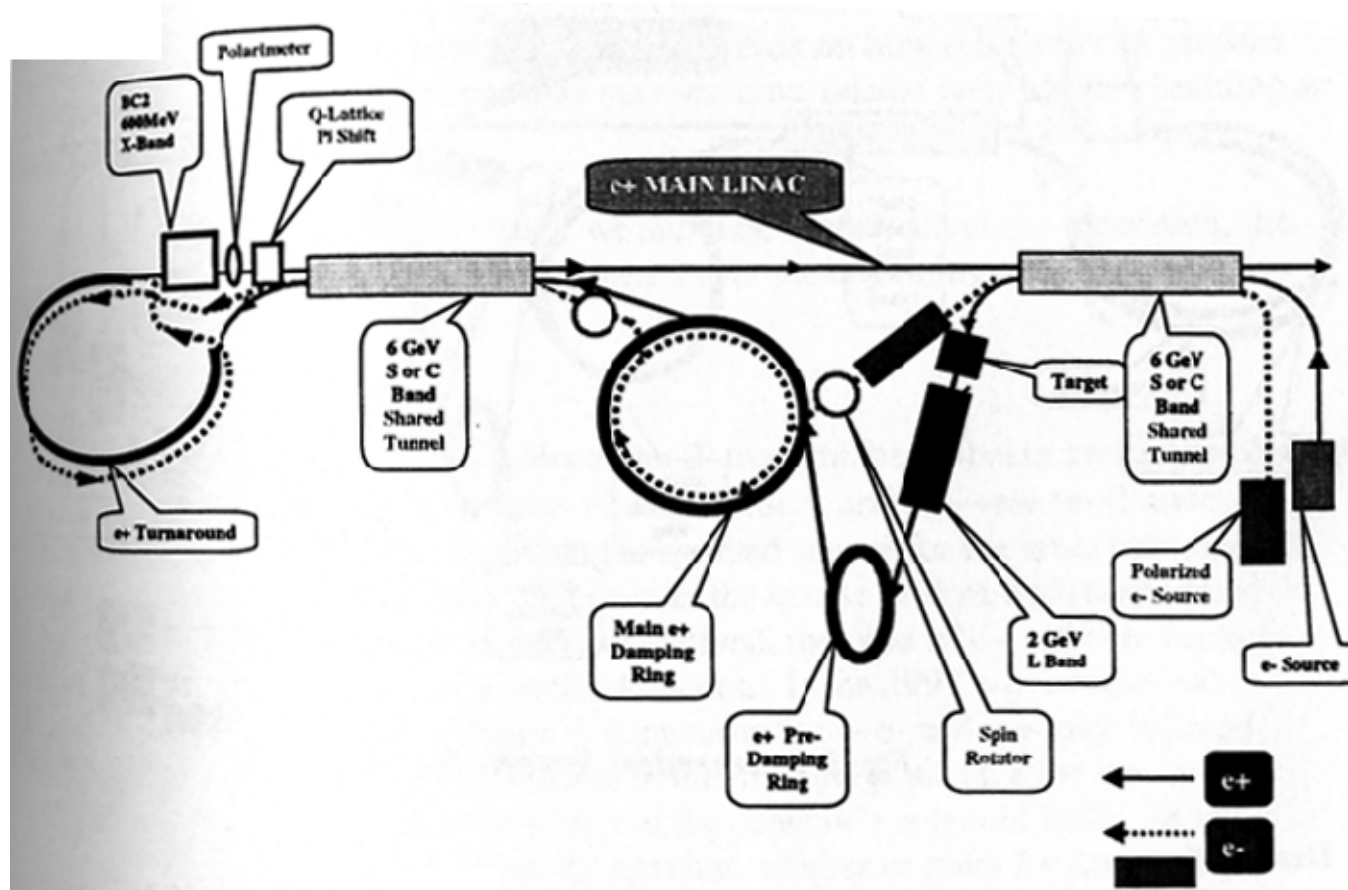
These compatibility issues are being evaluated at the  
Stanford Linear Accelerator site.

# Polarity Reversal Model

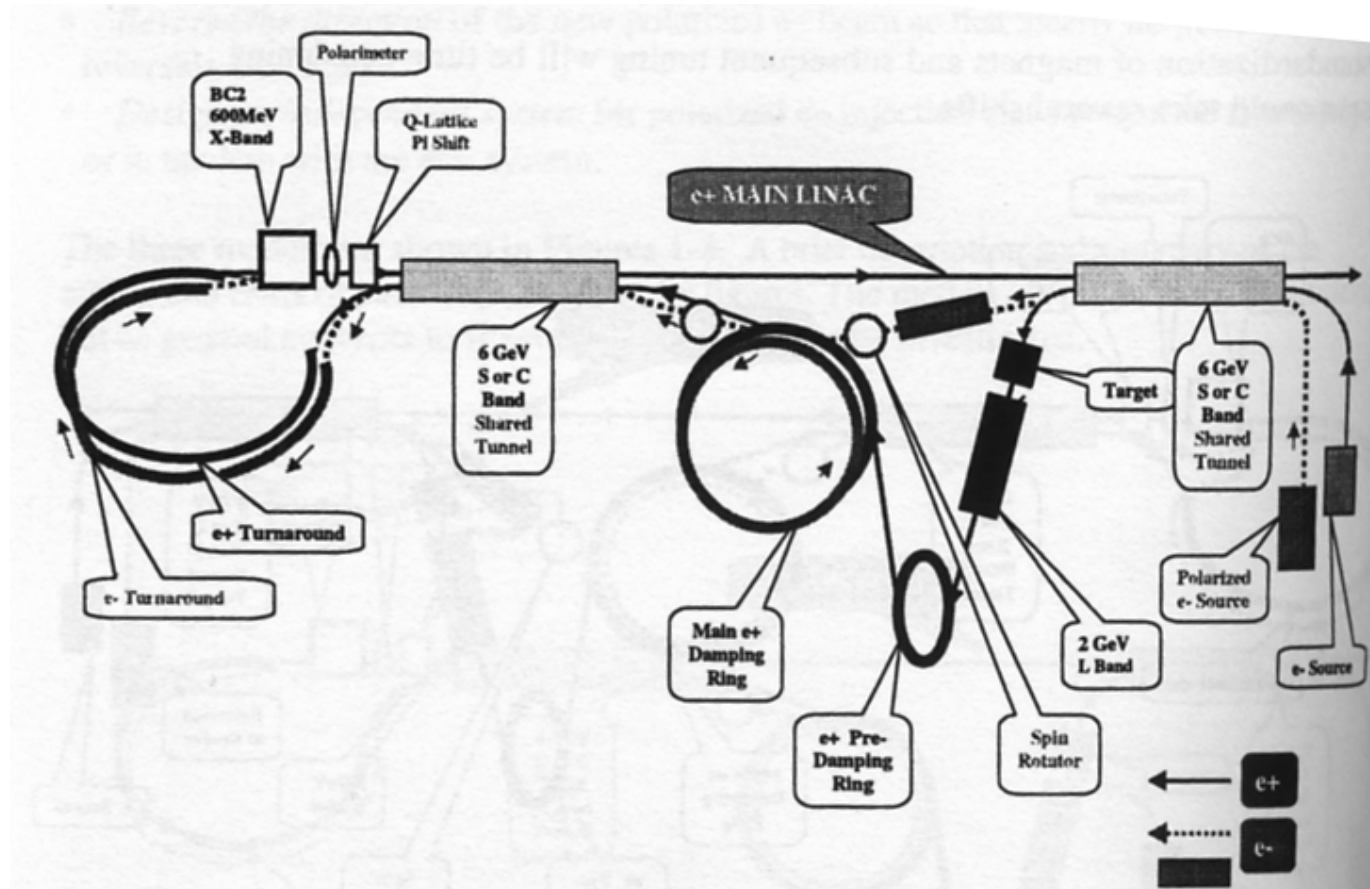




# Direction Reversal Model



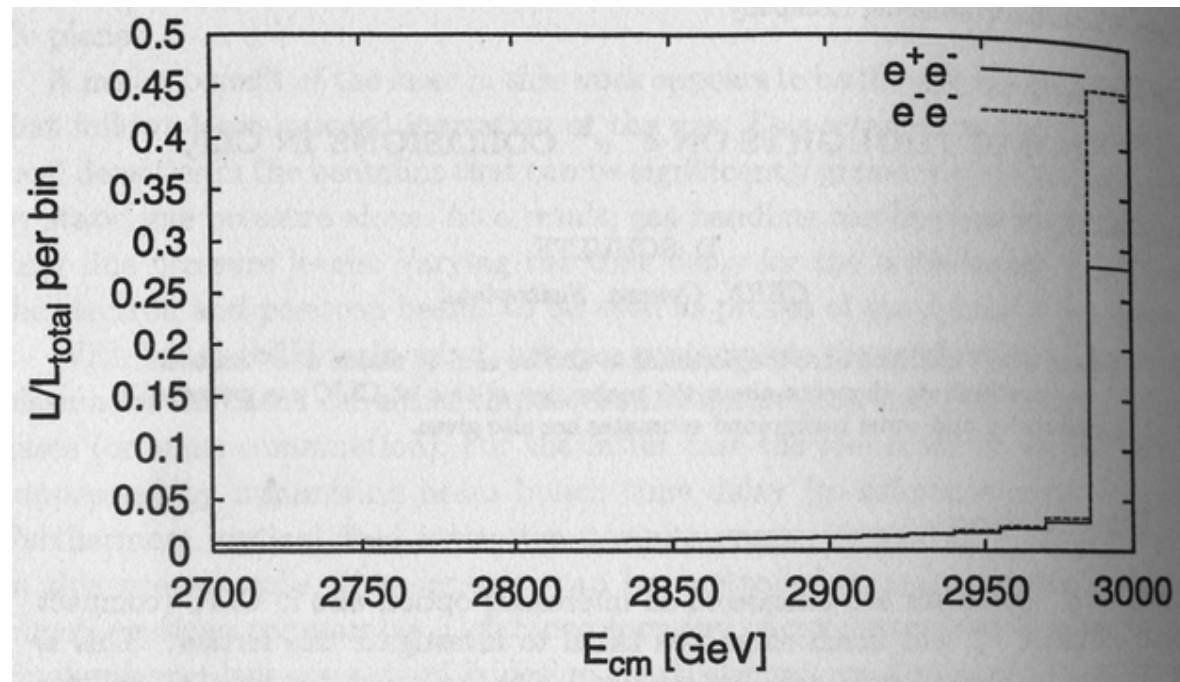
# Independent Systems Model



# Specific Machine Aspects (CLIC)

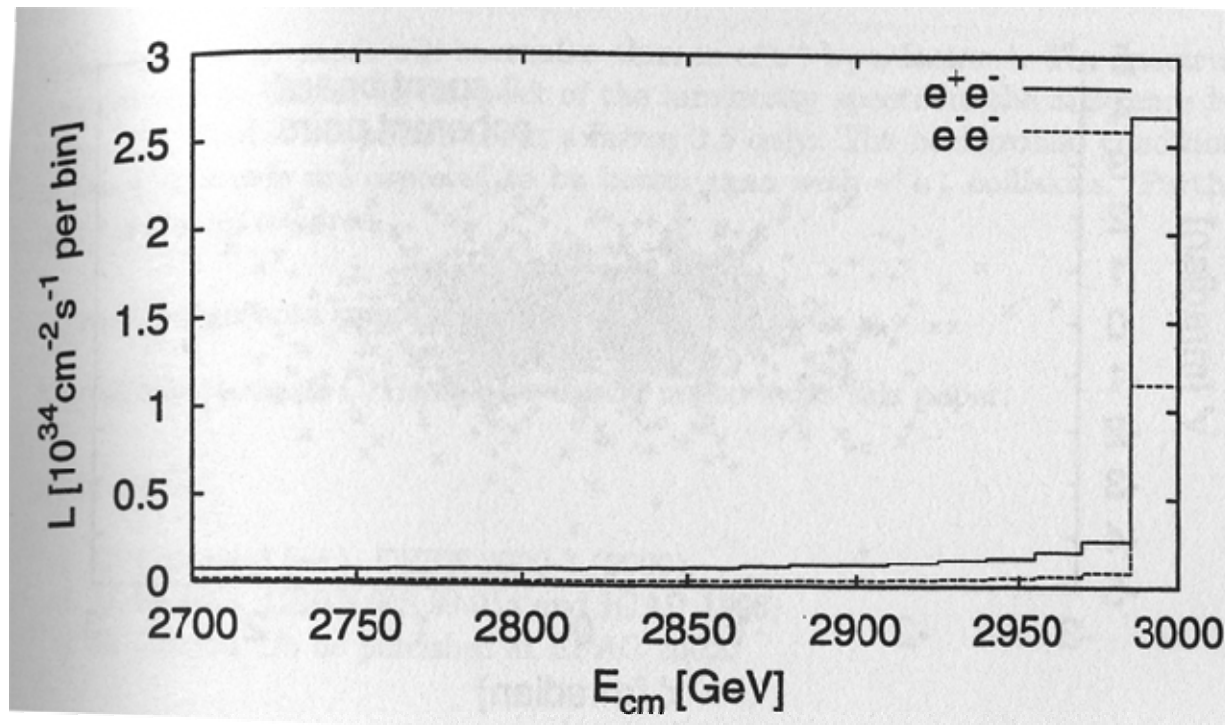
- The CERN laboratory plans to have a high-energy linear collider that can go to relatively high energies and good luminosities. It is worthwhile considering the advisability of planning for electron-electron as well as electron-positron interactions in that machine project.
- It is also significant that, irrespective of the interaction and detection equipment, the relative luminosities of  $e^+e^-$  and  $e-e^-$  colliders are considerably in favor of the  $e-e^-$  version, as shown in the two figures added below.

# Specific Machine Aspects (CLIC)



The relative luminosity spectrum for the  $e^+e^-$  and the  $e^-e^-$  case.  
The bins have a width of  $0.5\%E_{\text{cm}}$ .

# Specific Machine Aspects (CLIC)



The absolute luminosity spectrum for the  $e^+e^-$  and the  $e^-e^-$  case.