

Semiclassical soliton type solutions of the nonlocal Gross-Pitaevsky equation

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Nonstationary multidimensional Gross-Pitaevsky equation

$$-ih\frac{\partial}{\partial t}\Psi + \hat{H}_\varkappa\Psi = 0, \quad \|\Psi\|_{\mathbb{L}_2(\mathbb{R}_x^n)} = 1, \quad (1)$$

$$\hat{H}_\varkappa = \hat{H}_\varkappa[\Psi] = \hat{H}_0 + \varkappa \int_{\mathbb{R}^n} |\Psi(y, t)|^2 V(x - y) dy, \quad (2)$$

$$\hat{H}_0 = \frac{\hat{p}^2}{2m} + U(x), \quad \hat{p} = -ih\nabla.$$

$\varkappa = \text{Const}$, h -small parametr, $h \rightarrow 0$.

$U(x)$ is the smooth potential of the external field,

$V(x - y)$ is the smooth translate-invariant potential of the self-action.

Statment of problem: semiclassically soliton-type solutions , $h \rightarrow 0$

$$\begin{cases} -ih\frac{\partial}{\partial t}\Psi + \hat{H}_\varkappa[\Psi]\Psi = O(h^{3/2}), & (O(h^{3/2}) \text{ in the norm of } \mathbb{L}_2(\mathbb{R}_x^n)) \\ \Psi|_{t=0} = \Psi_0(x, h) = \exp\left[\frac{i}{h}((p_0, x - x_0) + (x - x_0, B_0(x - x_0)))\right] \end{cases} \quad (3)$$

$$B_0 = B_0^t, \quad \text{Im}B_0 > 0$$

In the linear case ($\varkappa = 0$) in (2), $\hat{H}_\varkappa = \hat{H}_0 = \frac{\hat{p}^2}{2} + U(x)$
 asymptotic solutions ($h \rightarrow 0$) have the following form:

$$\Psi(x, t, h) = N \exp\left[\frac{i}{h}(S_{\text{cl}}(t) + (p_{\text{cl}}(t), x - x_{\text{cl}}(t)))\right] \cdot \exp\left[\frac{i}{h}(x - x_{\text{cl}}(t), BC^{-1}(t)(x - x_{\text{cl}}(t)))\right] (\det C(t))^{-1/2}. \quad (1)$$

Where $S_{\text{cl}}(t)$ is the classical action

$$S_{\text{cl}} = \int_0^t \left(\frac{\dot{x}_{\text{cl}}^2(\tau)}{2} - U(x_{\text{cl}}(\tau)) \right) d\tau, \quad (2)$$

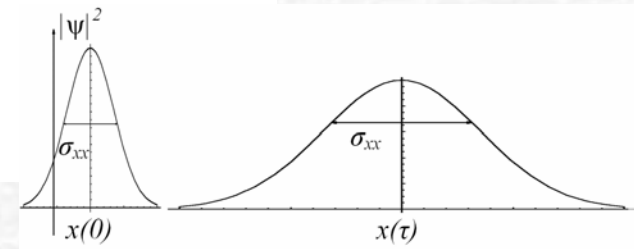
system in variations

$$\begin{pmatrix} \dot{B} \\ \dot{C} \end{pmatrix} = \begin{pmatrix} 0 & -U''_{xx}(x_{\text{cl}}(t)) \\ \mathbb{E}_n & 0 \end{pmatrix} \begin{pmatrix} B \\ C \end{pmatrix}, \quad (3)$$

$$B(0) = B_0, C(0) = \mathbb{E}_n = ((\delta_{ij}))_{n \times n}.$$

Classical system $\left\{ \begin{array}{l} \dot{p}_{\text{cl}} = \frac{\partial U(x_{\text{cl}})}{\partial x_{\text{cl}}} \\ \dot{x}_{\text{cl}} = p_{\text{cl}} \\ p_{\text{cl}}|_{t=0} = p_0, x_{\text{cl}}|_{t=0} = x_0 \end{array} \right.$ - apriory Hamiltonian system

$$\sigma_{xx}(t) = \langle \Psi, (x - x_{\text{cl}})^2 \Psi \rangle_{\mathbb{L}_2} \rightarrow \infty, t \rightarrow \infty$$



The first main questions arises: What is classical dynamical system corresponding to nonlinear \hat{H}_\varkappa in the limit $\hbar \rightarrow 0$?

Classical trajectory of quantum particle in the state $\Psi(x, t, \hbar) : i\hbar\Psi_t = \hat{H}_\varkappa\Psi$
 $Z_\Psi(t, \hbar) = (P(t, \hbar), X(t, \hbar)) \in \mathbb{R}^{2n}$, $P(t, \hbar) = \langle \Psi, -i\hbar\nabla\Psi \rangle_{\mathbb{L}_2}$, $X(t, \hbar) = \langle \Psi, x\Psi \rangle_{\mathbb{L}_2}$.

$P(t, \hbar), X(t, \hbar)$ are smooth functions, $\hbar \rightarrow 0$.

Statment 1. Classical system corresponding to the nonlinear operator \hat{H}_\varkappa (**Hamilton -Erenfest system**):

$$\begin{cases} \dot{p} = -\nabla_x U(x) - \varkappa \nabla_\tau V(0) - \\ \quad - \nabla_x \left[\frac{1}{2} \text{Sp}(U''_{xx}(x) + 2\varkappa V''_{xx}(x-y)) \sigma_{xx} \right] \Big|_{y=x} \\ \dot{x} = p \\ \dot{\Delta} = JM_\varkappa(x)\Delta - \Delta M_\varkappa(x)J, \quad (p, x) \in \mathbb{R}^{2n}. \end{cases} \quad (1)$$

$$\Delta = \begin{pmatrix} \sigma_{pp} & \sigma_{px} \\ \sigma_{xp} & \sigma_{xx} \end{pmatrix} - \text{are of order } O(\hbar), \quad \sigma_{xp}^t = \sigma_{px}, \quad \sigma_{xx}^t = \sigma_{xx}, \quad \sigma_{pp}^t = \sigma_{pp},$$

$$M_\varkappa(x) = \begin{pmatrix} \mathbb{E}_n & 0 \\ 0 & U''_{xx}(x) + \varkappa V''_{\tau\tau}(0) \end{pmatrix}, \quad U''_{xx}(x) = \left(\frac{\partial^2 U(x)}{\partial x_i \partial x_j} \right)_{n \times n},$$

Initial conditions

$$\begin{cases} p|_{t=0} = \langle \Psi_0, \hat{p}\Psi_0 \rangle = p_0, \quad x|_{t=0} = \langle \Psi_0, x\Psi_0 \rangle = x_0 \\ (\sigma_{pp}|_{t=0})_{km} = \langle \Psi_0, \hat{\Delta}p_k \hat{\Delta}p_m \Psi_0 \rangle = \frac{\hbar}{4}(B_0 B_0^+ + B_0^* B_0^t) \\ (\sigma_{xx}|_{t=0})_{km} = \langle \Psi_0, \hat{\Delta}x_k \hat{\Delta}x_m \Psi_0 \rangle = \frac{\hbar}{4} \\ (\sigma_{px}|_{t=0})_{km} = \frac{1}{2} \langle \Psi_0, (\hat{\Delta}p_k \hat{\Delta}x_m + \hat{\Delta}x_m \hat{\Delta}p_k) \Psi_0 \rangle = \frac{\hbar}{4}(B_0 + B_0^t), \quad k, m = \overline{1, n} \end{cases}$$

Hamiltonian system in variations with self- action:

$$\begin{aligned}
 \begin{pmatrix} \dot{B} \\ \dot{C} \end{pmatrix} &= JM_{\varkappa}(X_0(t)) \begin{pmatrix} B \\ C \end{pmatrix}, & \Delta(t) &= A(t)\Delta|_{t=0}A^+(t) \\
 B|_{t=0} &= B_0, C|_{t=0} = \mathbb{E}_n & & A(t) \text{ -the Cauchy matrix}
 \end{aligned} \tag{1}$$

where X_0 - is the main part of the $X(t, h) = X_0(t) + hX_1(t) + O(h^2)$

Remark: $A(t) = \begin{pmatrix} B(t) & B^*(t) \\ C(t) & C^*(t) \end{pmatrix} \begin{pmatrix} B(0) & B^*(0) \\ C(0) & C^*(0) \end{pmatrix}^{-1}$

Statement

Let the following conditions to be satisfy: $U(x), V(\tau) \in C^3(\mathbb{R}^n)$. Then localized asymptotic solution (mod $h^{3/2}$), $h \rightarrow 0, t \in [0, T]$ has the form

$$\begin{aligned}
 \Psi_{\nu}(x, t, h) &= N_{\nu} \exp\left[\frac{i}{h}(S_{\varkappa}(t, h) + (P(t, h), x - X(t, h)))\right] \\
 &\exp\left[\frac{i}{2h}(x - X(t, h), BC^{-1}(t)(x - X(t, h)))\right] \frac{1}{\sqrt{\det C(t)}}.
 \end{aligned}$$

where $S_{\varkappa}(t, h)$ is the classical phase

$$\begin{aligned}
 S_{\varkappa}(t, h) &= \int_0^t \left(\frac{\dot{X}^2}{2}(\tau, h) - U(X(\tau, h)) \right) d\tau - \varkappa V(0)t - \\
 &\quad - \frac{\varkappa}{2} \int_0^t \text{Sp}(\underline{V''_{\tau\tau}(0)\sigma_{xx}(\tau, h)})d\tau.
 \end{aligned}$$

Semiclassical solyton type solutions ($U(x) = 0$)

Let's determine ω_j^2 , ($j = 1, \dots, n$)-eigen values of matrix $\varkappa V''_{\tau\tau}(0)$

Theorem: Let the following conditions to be satisfied: $\varkappa V''_{\tau\tau}(0) > 0$, $V(\tau) \in C^3(\mathbb{R}^n)$ and $\omega_j^2 \neq \omega_k^2$, $j \neq k$, then semiclassical solyton type solutions has the form:

$$\Psi(x, t, h) = N_\nu \exp\left[\frac{i}{h}(S_\varkappa(t, h) + (P(t, h), x - X(t, h)))\right] \exp\left[\frac{i}{2h}(x - X(t, h), \dot{C}C^{-1}(t)(x - X(t, h)))\right] \frac{1}{\sqrt{C(t)}}, \quad (1)$$

$$S_\varkappa(t, h) = \int_0^t \frac{\dot{X}^2}{2}(\tau, h) d\tau - \varkappa V(0)t - \frac{\varkappa}{2} \int_0^t \text{Sp}(V''_{\tau\tau}(0)\sigma_{xx}(\tau, h)) d\tau.$$

$$X(t, h) = x_0 + p_0 t + \frac{(a_0 + h a_1) t^2}{2} +$$

$$+ h \sum_{\pm} \sum_{l, m=1, l \neq m}^n (d_{l, m}^{\pm} \cos((\omega_l \pm \omega_m) t) + k_{l, m}^{\pm} \sin((\omega_l \pm \omega_m) t)) +$$

$$+ h \sum_{l=1}^n (e_l \cos(2\omega_l t) + s_l \sin(2\omega_l t)),$$

$$P(t, h) = \dot{X}(t, h), \quad a_0 = -\varkappa \nabla_{\tau} V(0),$$

$$\sigma_{xx}(t, h) = h \Gamma_0 + \sum_{\pm} \sum_{l, m=1, l \neq m}^n (\cos((\omega_l \pm \omega_m) t) A_{l, m}^{\pm} -$$

$$- \sin((\omega_l \pm \omega_m) t) B_{l, m}^{\pm}) + \sum_{l=1}^n A_{l, l}^+ \cos(2\omega_l t) - B_{l, l}^+ \sin(2\omega_l t),$$

$$\|\sigma_{xx}(t)\| < \text{Const}, \quad \forall t \in \mathbb{R}$$

where Γ_0 , $A_{l, m}^{\pm}$, $B_{l, m}^{\pm}$ $-n \times n$ is the constant real matrix, a_1 , $d_{l, m}^{\pm}$, $k_{l, m}$, s_l , e_l real vectors from \mathbb{R}^n , $l, m = 1, \dots, n$ explicit form is easily calculated with formula $\Delta(t, h) = A(t) \Delta|_{t=0} A^+(t)$ with the Caushi matrix $A =$

$$\begin{pmatrix} \cos[t\sqrt{R}] & R^{-1/2} \sin[t\sqrt{R}] \\ -\sqrt{R} \sin[t\sqrt{R}] & \cos[t\sqrt{R}] \end{pmatrix}, \quad R = \varkappa V''_{\tau\tau}(0).$$

$(U(x) = \langle E(t), x \rangle)$:

Let be $V(\tau)$ –smooth function, $\tau \in \mathbb{R}^3$, grow with all it's derivates at $|\tau| \rightarrow \infty$ not faster, than $Ce^{|\tau|}$, $C = Const > 0$, and $\nabla_{\tau} V(0) = 0$, $\kappa V''_{\tau\tau}(0) > 0$. Then semiclassical solyton type solutions mod $h^{3/2}$ of the equation (1) in non-stationary homogeneous electric field ($U(x) = \langle E(t), x \rangle$, $E(t) \in \mathbb{R}^3$) have the form

$$\Psi(x, t, h) = \exp\left[-\frac{i}{h} \lambda_v t\right] \exp\left[\frac{i}{h} \int_0^t L(X(\tau), \dot{X}(\tau), \tau) d\tau\right] \exp\left[\frac{i}{h} \langle \dot{X}(t) m, x - X(t) \rangle\right] \Psi_{\lambda_v}(x - X(t))$$

where $L(X(t), \dot{X}(t), t) = \frac{\langle m\dot{X}(t), \dot{X}(t) \rangle}{2} - \langle E(t), X(t) \rangle$,

$X(t)$ -centre of gravity of the semiclassical solyton (-the any solution of the equation of Newton)– $m\ddot{X}(t) = -E(t)$,

$(\lambda_\nu, \Psi_{\lambda_\nu}(x)), x \in R^3$ - E. values and E.- function operator $\frac{(\hat{p})^2}{2m} + \kappa \int_{R^3} V(x-y) |\Psi|^2 dy$.

$$\lambda_\nu, \nu = (\nu_1, \nu_2, \nu_3), \nu_j \in Z: \quad \lambda_\nu = \kappa V(0) + h \sum_{k=1}^3 \Omega_k \left(\nu_k + \frac{1}{2} \right) + \frac{1}{2} Sp[V_{\tau\tau}(0) \sigma_{xx}],$$

$$\sigma_{xx} = \frac{h}{4} (C D_\nu C^+ + C^* D_\nu C^T)$$

where 3×3 matrix C consist of the E. vectors $f_j, \langle f_j, f_{j'} \rangle = \delta_{jj'}, j, j' = 1, 2, 3$, which associate E. values $\Omega_j^2, j = 1, 2, 3$ of the matrix $\kappa V''_{\tau\tau}(0)$. $D_\nu = diag(2\nu_k + 1), k = 1, 2, 3$.

$$\Psi_{\lambda_\nu}(x, h) = (\pi h)^{-3/4} (\det C)^{-1/2} \exp\left[\frac{i}{h} \frac{1}{2} \langle \Delta x, B C^{-1} \Delta x \rangle\right] He_\nu \left(-i (\tilde{C}^*)^{-1} \frac{\Delta x \sqrt{2}}{\sqrt{h}} \right)$$

$He_\nu(\xi), \xi \in R^3$ - multidimensional Hermit polynomial, which associate matrix $W = -\tilde{C}^+ (\tilde{C}^{-1})^T$,

$$(He_\nu(\xi) = \frac{(-1)^{|\nu|}}{\sqrt{n!}} \left(\frac{\partial}{\partial \xi} - W \xi \right)^\nu \cdot 1), \Delta x = x - x_0, x_0 \in R^3$$

Definition. Class of trajectory-concentrated functions

$$Z(t, h) = (P(t, h), X(t, h)) \in \mathbb{R}^{2n}, S(t, h) \in \mathbb{R}, S(0, h) = 0$$

$$\forall t K = \left\{ \Phi, \Phi = \phi\left(\frac{x - X(t, \hbar)}{\sqrt{\hbar}}, t\right) \exp\left[\frac{i}{\hbar}(S(t, \hbar) + (P(t, \hbar), x - X(t, \hbar)))\right] \right\},$$

where $\phi(\xi, t) \in S(\mathbb{R}^n)$

Assumption: Let $\Psi(x, t, h)$ exist and satisfy eq. Hartree (mod $h^{3/2}$, $\Psi \in K$)

Example

$$V(\tau) = \exp[-\sum_{i=0}^3 \gamma_i(\tau_i, \tau_i)], \tau_i \in \mathbb{R}^1, \gamma_i < 0.$$

Soliton type solution has the form

$$\Psi_\nu(x, t, h) = N_\nu \exp\left[\frac{i}{h}(S_\kappa(t, h) + (P(t, h), x - X(t, h)))\right] \\ \exp\left[\frac{i}{2h}(x - X(t, h), \dot{C}C^{-1}(t)(x - X(t, h)))\right] \frac{1}{\sqrt{C(t)}} \hat{A}^{+\nu}(t) 1,$$

$$((C(t)))_{ij} = \left(\left(\frac{b_{ij}}{\Omega_i} \sin(t\Omega_i) + \delta_{ij} \cos(t\Omega_i)\right)\right)_{ij}$$

$$S_\kappa(t, h) = \int_0^t \frac{\dot{X}^2}{2}(\tau, h) d\tau - \kappa \int_0^t \sum_{i=0}^3 \gamma_i \sigma_{x_i x_i}(\tau, h) d\tau,$$

$$\sigma_{x_i x_i} = \frac{h(1 + 2\mu_i)}{8\text{Im}B\Omega_i^2} (|b|^2 + \Omega_i^2 + \cos(2t\Omega_i)(-|b|^2 + \Omega_i^2) + \sin(2t\Omega_i)2(\text{Re}b)^2\Omega_i)$$

$$\hat{A}^{+\nu} = [\hat{A}_1^+]^{\nu_1} \dots [\hat{A}_n^+]^{\nu_n},$$

$$\hat{A}_j^+(t) = \frac{1}{\sqrt{h}} ((z_j^*(t), \hat{p}) - (\dot{z}_j^*(t) - \dot{C}C^{-1}(t)z_j^*(t), (x - X(t, h))))),$$

$z_j(t)$ – columns of matrix $C(t)$