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**Polarized parton densities and higher twist
corrections in the light of the
recent CLAS and COMPASS data**

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OUTLINE

- Method of analysis – **higher twist** corrections are taken into account
- Two **new** sets of very precise data are included in the analysis

- **low** Q^2 CLAS data

- COMPASS data mainly at **large** Q^2

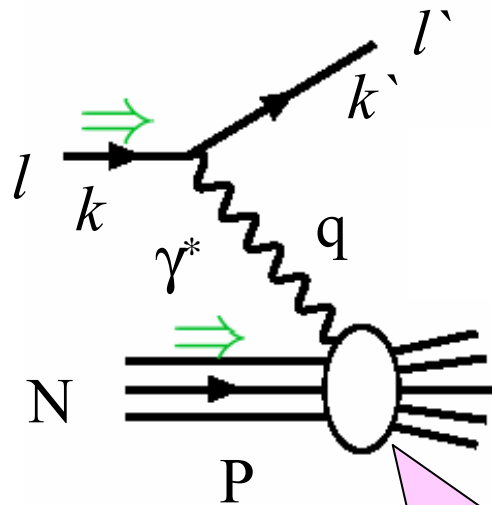


*Very different
kinematic regions*

- Impact of the **new** data on **LSS'05 polarized PD and HT**
- The sign of the gluon polarization
- Summary

Inclusive DIS

one of the best tools to study
the structure of **nucleon**



$$Q^2 = -q^2 = 4EE' \sin^2(\theta/2)$$

$$x = Q^2/(2Mv) \quad v = E - E'$$

DIS regime $\implies Q^2 \gg M^2, v \gg M$

$F_i(x, Q^2)$ $g_i(x, Q^2)$

unpolarized SF

polarized SF

$$Q^2 \approx 1 - 5 \text{ GeV}^2, 4 < W^2 < 10 \text{ GeV}^2$$

preasymptotic region

As in the unpolarized case the main goal is:

- to test **QCD**
- to extract from the DIS data the **polarized PD**

$$\Delta q(x, Q^2) = q_+(x, Q^2) - q_-(x, Q^2)$$

$$\Delta \bar{q}(x, Q^2) = \bar{q}_+(x, Q^2) - \bar{q}_-(x, Q^2)$$

$$\Delta G(x, Q^2) = G_+(x, Q^2) - G_-(x, Q^2)$$

where "+" and "-" denote the helicity of the parton, along or opposite to the helicity of the parent nucleon, respectively.

The knowledge of the polarized PD will help us:

- to make predictions for other processes like polarized **hadron-hadron** reactions, etc.
- more generally, to answer the question how the helicity of the nucleon is divided up among its constituents:

$$S_z = 1/2 = 1/2 \Delta\Sigma(Q^2) + \Delta G(Q^2) + L_z(Q^2)$$

$$\Delta\Sigma = \Delta u + \Delta\bar{u} + \Delta d + \Delta\bar{d} + \Delta s + \Delta\bar{s}$$

the parton polarizations Δq_a and ΔG are the first moments

$$\Delta q_a(Q^2) = \int_0^1 dx \Delta q_a(x, Q^2) \quad \Delta G(Q^2) = \int_0^1 dx \Delta G(x, Q^2)$$

of the helicity densities: $\Delta u(x, Q^2), \Delta\bar{u}(x, Q^2), \dots, \Delta G(x, Q^2)$

DIS Cross Section Asymmetries

Measured quantities

$$A_{\parallel} = \frac{d\sigma^{\downarrow\uparrow} - d\sigma^{\uparrow\uparrow}}{d\sigma^{\downarrow\uparrow} + d\sigma^{\uparrow\uparrow}},$$

$$A_{\perp} = \frac{d\sigma^{\downarrow\Rightarrow} - d\sigma^{\uparrow\Rightarrow}}{d\sigma^{\downarrow\Rightarrow} + d\sigma^{\uparrow\Rightarrow}}$$

$$(A_{\parallel}, A_{\perp}) \Rightarrow (A_1, A_2) \Rightarrow (g_1, g_2)$$

where A_1, A_2 are the virtual photon-nucleon asymmetries.

At present, A_{\parallel} is much better measured than A_{\perp}

If A_{\parallel} and A_{\perp} are measured

$$\Rightarrow g_1 / F_1$$

If only A_{\parallel} is measured

$$\Rightarrow \frac{A_{\parallel}^N}{D} \approx (1 + \gamma^2) \frac{g_1}{F_1}$$

$$\gamma^2 = 4M_N^2 x^2 / Q^2 \quad \text{- kinematic factor}$$

NB. γ cannot be neglected in the **SLAC**,
HERMES and **JLab** kinematic regions

Theory

In QCD

$$g_1(x, Q^2) = g_1(x, Q^2)_{LT} + g_1(x, Q^2)_{HT}$$

$$g_1(x, Q^2)_{LT} = g_1(x, Q^2)_{pQCD} + \frac{M^2}{Q^2} h^{TMC}(x, Q^2) + O\left(\frac{M^4}{Q^4}\right)$$

$$g_1(x, Q^2)_{HT} = h(x, Q^2) / Q^2 + O\left(\frac{\Lambda^4}{Q^4}\right)$$

dynamical HT power corrections ($\tau=3,4$)
=> non-perturbative effects (model dependent)

target mass corrections
which are calculable
A. Piccione, G. Ridolfi

In NLO pQCD

$$g_1(x, Q^2)_{pQCD} = \frac{1}{2} \sum_q^{N_f} e_q^2 [(\Delta q + \Delta \bar{q}) \otimes (1 + \frac{\alpha_s(Q^2)}{2\pi} \delta C_q) + \frac{\alpha_s(Q^2)}{2\pi} \Delta G \otimes \frac{\delta C_G}{N_f}]$$

$\delta C_q, \delta C_G$ – Wilson coefficient functions

polarized PD evolve in Q^2

$N_f (=3)$ - the number of flavors

according to **NLO DGLAP** eqs.

- An important difference between the kinematic regions of the unpolarized and *polarized* data sets
- A lot of the present data are at moderate Q^2 and W^2 :

$$Q^2 \approx 1-5 \text{ GeV}^2, \quad 4 < W^2 < 10 \text{ GeV}^2$$

*preasymptotic
region*

While in the determination of the PD in the unpolarized case we can cut the low Q^2 and W^2 data in order to eliminate the less known non-perturbative HT effects, it is **impossible** to perform such a procedure for the present data on the spin-dependent structure functions without losing too much information.

$$\alpha(\Lambda^2/Q^2)$$

➔ HT corrections have to be **accounted for** in **polarized** DIS !

Method of analysis

$$\left[\frac{g_1(x, Q^2)}{F_1(x, Q^2)} \right]_{\text{exp}} \xleftrightarrow{\chi^2} \frac{g_1(x, Q^2)_{\text{LT}} + h^{g_1}(x)/Q^2}{F_1(x, Q^2)_{\text{exp}}}$$

$F_2^{\text{NMC}}, R_{1998}(\text{SLAC})$

in model independent way

Input PD

$$\Delta f_i(x, Q_0^2) = A_i x^{\alpha_i} f_i^{\text{MRST}}(x, Q_0^2) \quad Q_0^2 = 1 \text{ GeV}^2, A_i, \alpha_i - \text{free par.}$$

$h^p(x_i), h^n(x_i) - 10$ parameters ($i = 1, 2, \dots, 5$) to be determined from a fit to the data

8-2(SR) = 6 par. associated with PD; positivity bounds imposed by MRST'02 unpol. PD

SUM

$$a_3 = g_A = (\Delta u + \Delta \bar{u})(Q^2) - (\Delta d + \Delta \bar{d})(Q^2) = F - D = 1.2670 \pm 0.0035$$

RULES

$$a_8 = (\Delta u + \Delta \bar{u})(Q^2) + (\Delta d + \Delta \bar{d})(Q^2) - 2(\Delta s + \Delta \bar{s})(Q^2) = 3F - D = 0.585 \pm 0.025$$

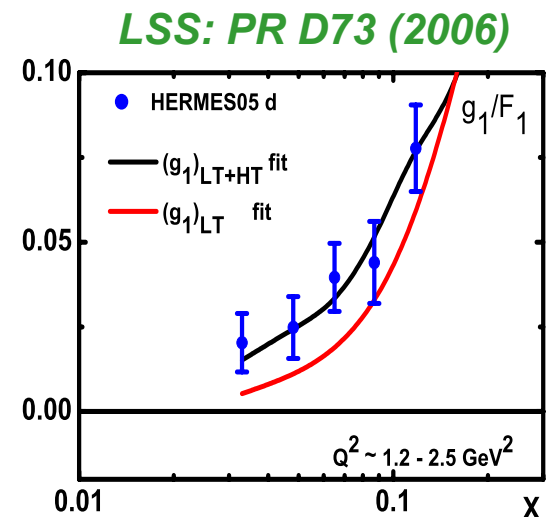
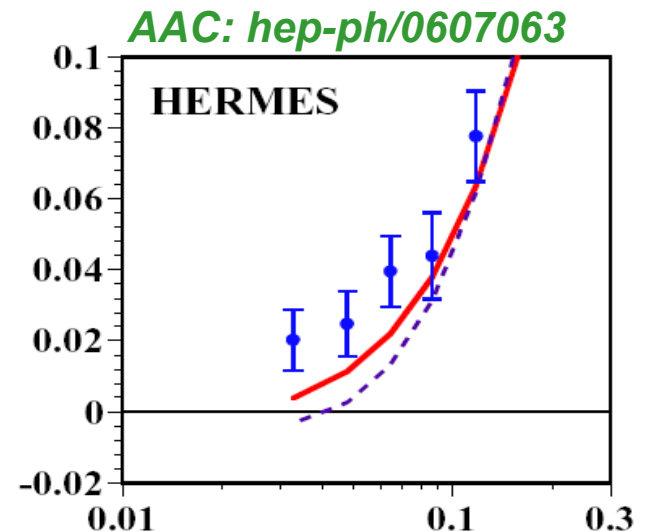
Flavor symmetric sea convention: $\Delta u_{\text{sea}} = \Delta \bar{u} = \Delta d_{\text{sea}} = \Delta \bar{d} = \Delta s = \Delta \bar{s}$

Higher twist effects

(CLAS'06 and COMPASS'06 not included)

$$g_1 = (g_1)_{LT} + h^{g_1}(x)/Q^2$$

- The low x and low Q^2 ($1.2 \sim 2.5 \text{ GeV}^2$) HERMES/d data can **not** be described by the **LT** (logarithmic in Q^2) term in $g_1 \Rightarrow$ **red curves**
- Excellent agreement with the data if the **HT corrections** to g_1 are taken into account in the analysis



DATA
(old set)

CERN **EMC** - A_1^p **SMC** - A_1^p, A_1^d **COMPASS'05** - A_1^d

DESY **HERMES** - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

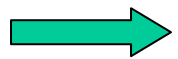
SLAC **E142, E154** - A_1^n **E143, E155** - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

JLab **Hall A** - $\frac{g_1^n}{F_1^n}$

$$A_1^N \approx (1 + \gamma^2) \frac{g_1^N}{F_1^N}$$

$$\gamma^2 = 4M^2x^2/Q^2 - \text{kinematic factor}$$

Number of exp. points: **190**



LSS'05 polarized PD and HT (PR D73, 2006)

DATA

CERN **EMC** - A_1^p **SMC** - A_1^p, A_1^d **COMPASS'05** - A_1^d

DESY **HERMES** - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

SLAC **E142, E154** - A_1^n **E143, E155** - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

JLab **Hall A** - $\frac{g_1^n}{F_1^n}$ **CLAS'06** - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

$$A_1^N \approx (1 + \gamma^2) \frac{g_1^N}{F_1^N}$$

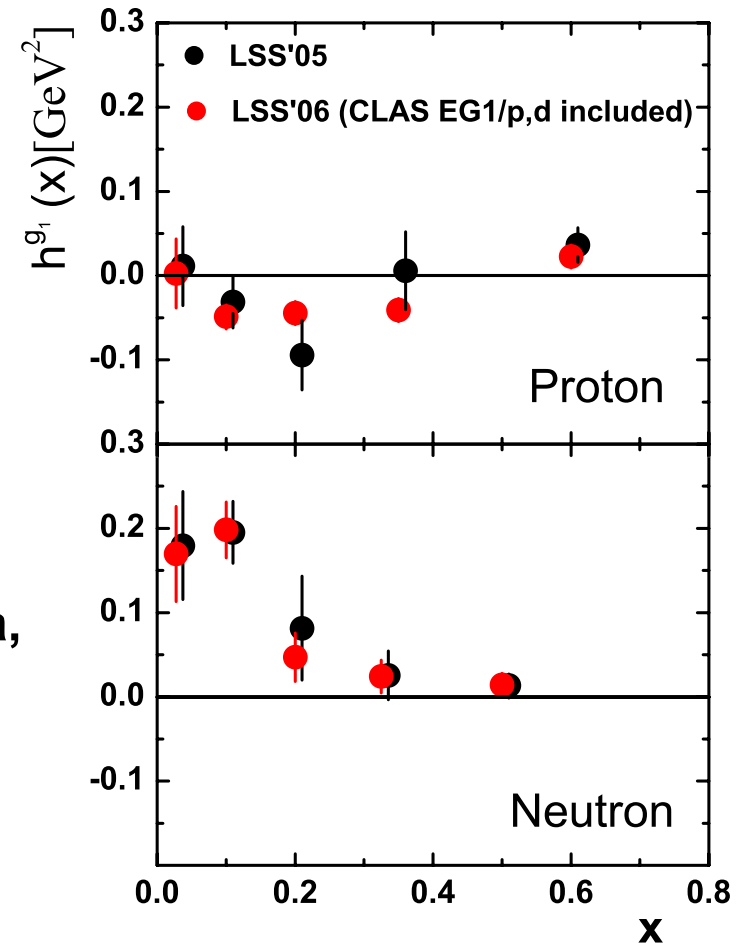
$$\gamma^2 = 4M^2x^2/Q^2 - \text{kinematic factor}$$

Number of exp. points: **190** \longrightarrow **823**

Effect of CLAS'06 p and d data (*PL B641, 11, 2006*) on polarized PD and HT

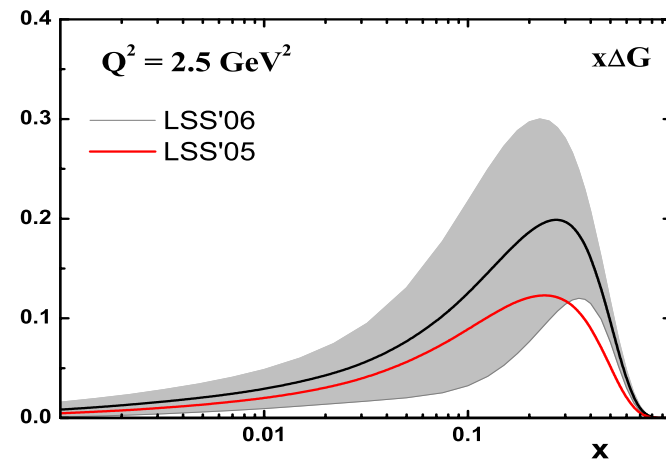
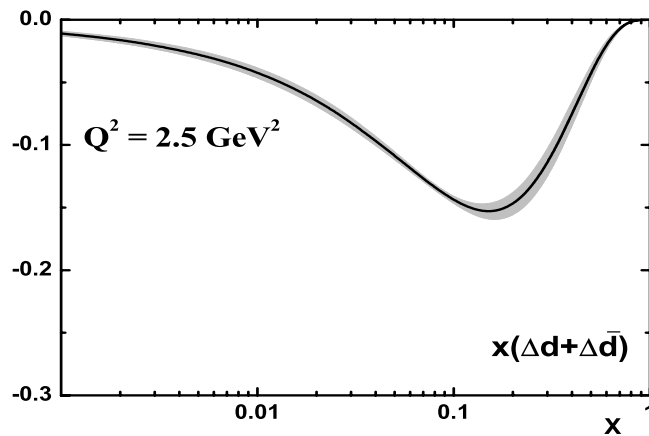
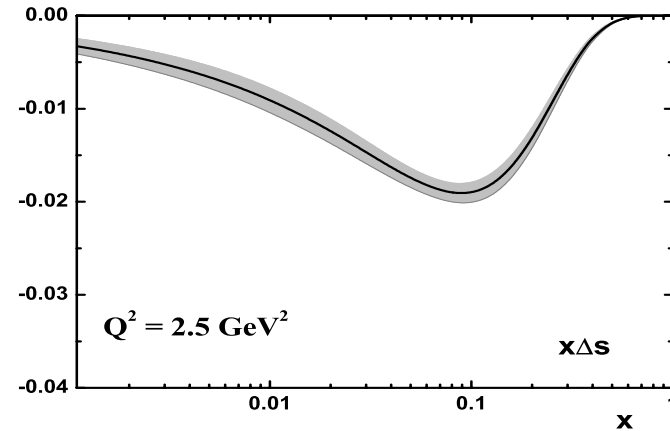
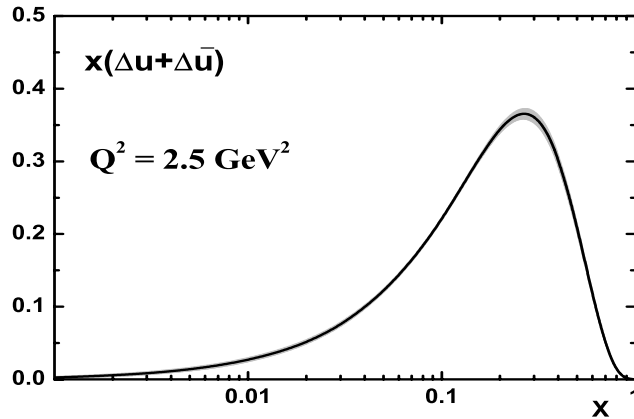
- Very accurate data on g_1^p and g_1^d at **low Q^2 : 1~4 GeV²** for **$x \sim 0.1 - 0.6$**
- The determination of HT/p and HT/n is **significantly improved** in the *CLAS* x region compared to HT(LSS'05)
- As expected, the central values of PPD are practically **not** affected by *CLAS* data, but the accuracy of its determination is **essentially improved** (a **consequence** of much better determination of HT corrections to g_1)

LSS'05: PR D73 (2006)

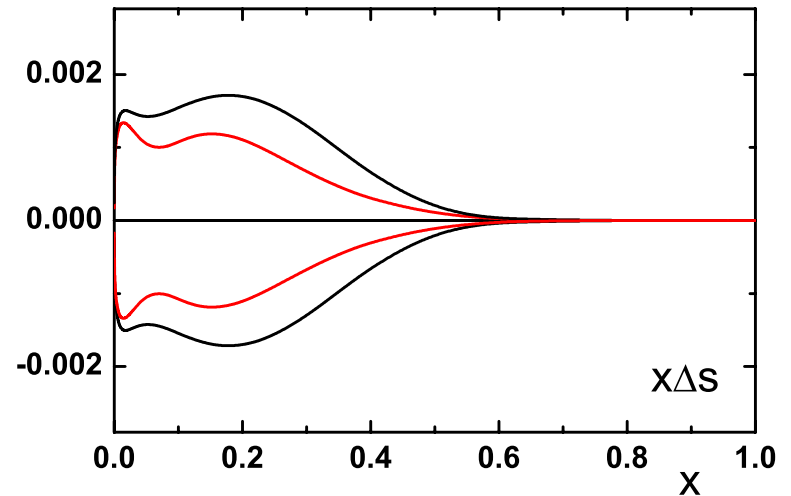
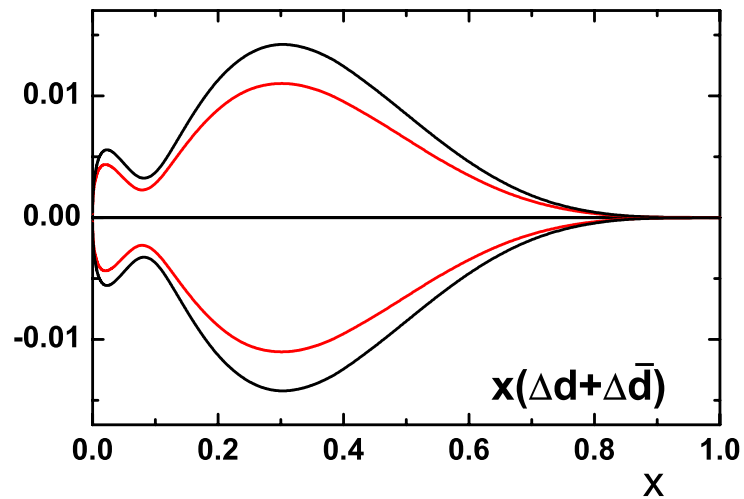
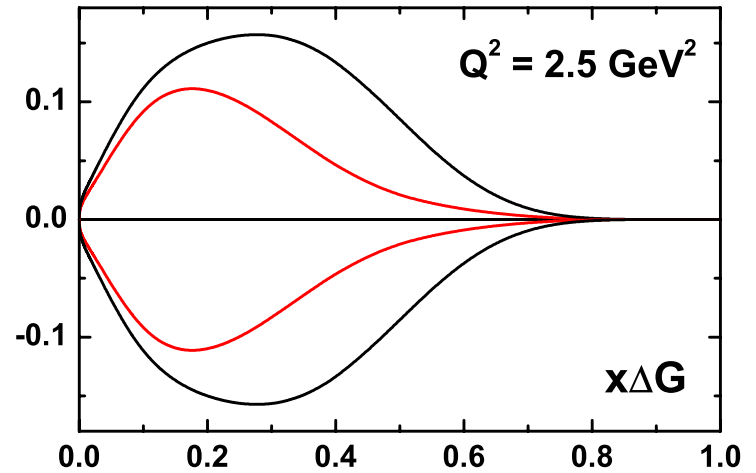
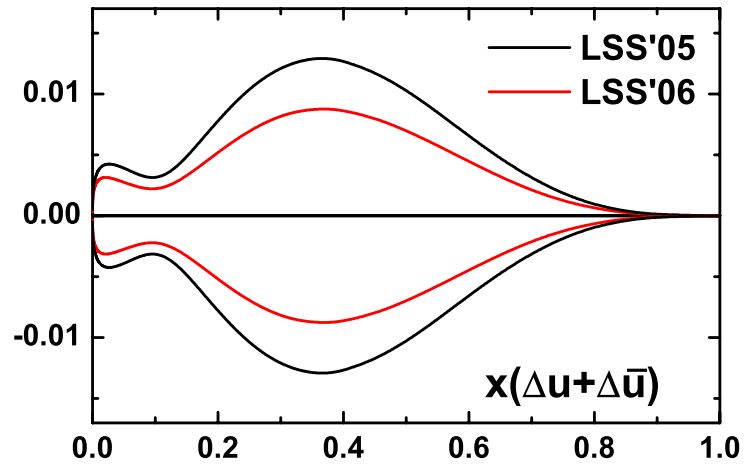


LSS'06 NLO($\overline{\text{MS}}$) polarized PDFs

The **quark** densities (central values) are identical with those of **LSS'05**.



Impact of CLAS'06 data on the uncertainties for NLO polarized PD



The first moments of higher twist

- Thanks to the **very precise CLAS** data the **first** moments of HT corrections are now **much better** determined.

$$\bar{h}^N = \int_{0.0045}^{0.75} dx h^N(x), \quad N = p, n$$

$$\bar{h}^p = (-0.014 \pm 0.005) \text{ GeV}^2$$

$$\bar{h}^n = (0.037 \pm 0.008) \text{ GeV}^2$$

$$\bar{h}^p - \bar{h}^n = (-0.051 \pm 0.009) \text{ GeV}^2$$

$$\bar{h}^p + \bar{h}^n = (0.023 \pm 0.009) \text{ GeV}^2$$

- $\bar{h}^p - \bar{h}^n < 0$ ← In agreement with the **instanton model** predictions and **sum rules** in QCD

- $\bar{h}^p + \bar{h}^n < |\bar{h}^p - \bar{h}^n|$ ← In agreement with **1/N_c** expansion in QCD (*Balla et al., NP B510, 327, 1998*)

The main message from this analysis

→ It is **impossible** to describe the very precise CLAS data if the HT corrections are **NOT taken into account**

NOTE: If the **low Q^2** data are **not too accurate**, it would be possible to describe them using only the leading twist term (logarithmic in Q^2) of g_1 , *i.e.* to **mimic** the power in Q^2 dependence of g_1 with a logarithmic one (using different forms for the input PDFs and/or more free parameters associated with them) which was done in the analyses of another groups before the CLAS data have appeared.

DATA

CERN EMC - A_1^p SMC - A_1^p, A_1^d COMPASS'06 - A_1^d

DESY HERMES - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

SLAC E142, E154 - A_1^n E143, E155 - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

JLab Hall A - $\frac{g_1^n}{F_1^n}$ CLAS'06 - $\frac{g_1^p}{F_1^p}, \frac{g_1^d}{F_1^d}$

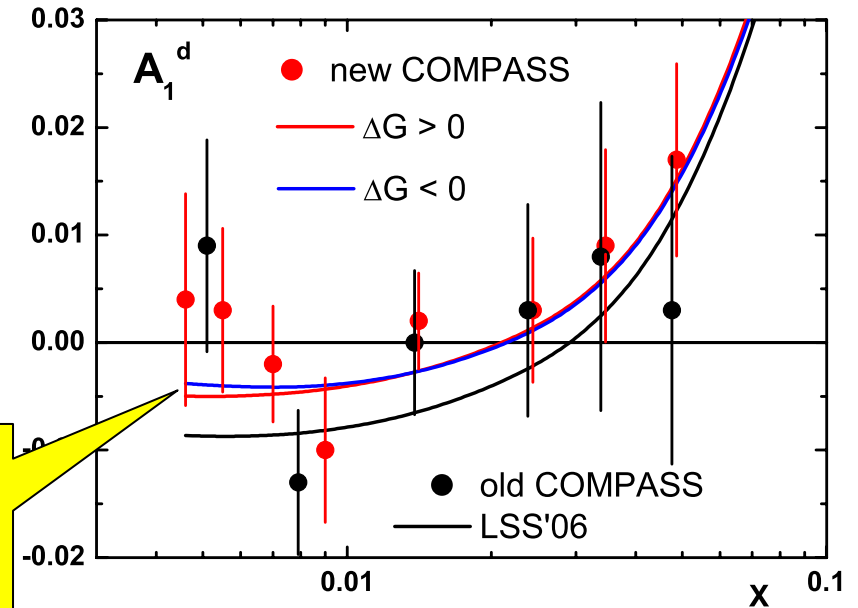
$$A_1^N \approx (1 + \gamma^2) \frac{g_1^N}{F_1^N}$$

$$\gamma^2 = 4M^2x^2/Q^2 - \text{kinematic factor}$$

Number of exp. points: 823 \longrightarrow 826

Effect of COMPASS'06 A_1^d data (*hep-ex/0609038*) on polarized PD and HT

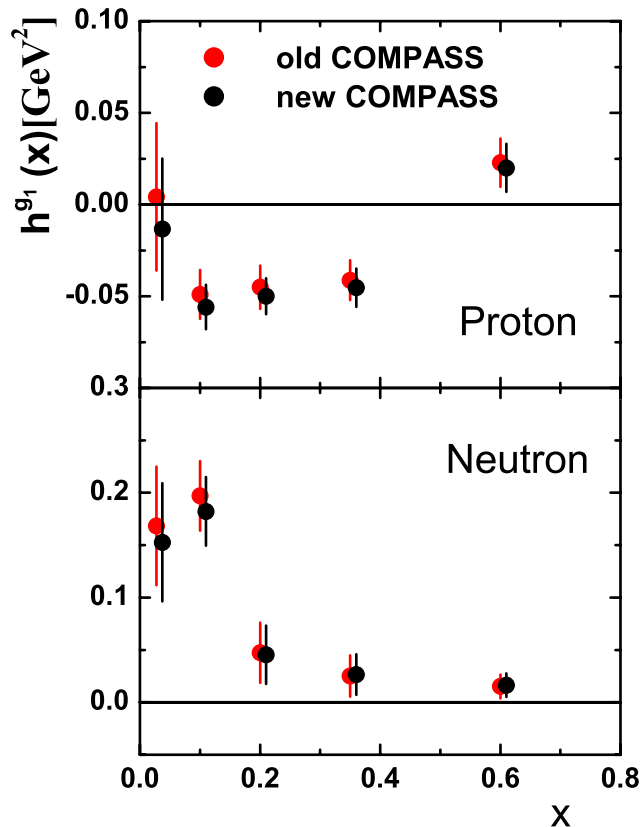
In contrast to the *CLAS* data, the *COMPASS* data are mainly at large Q^2 and the only precise data at small x : $0.004 < x < 0.02$. The new data are based on 2.5 times larger statistics than those of *COMPASS'05*



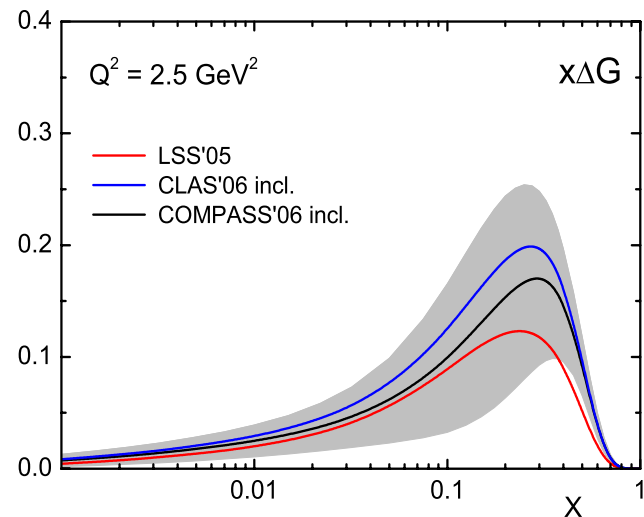
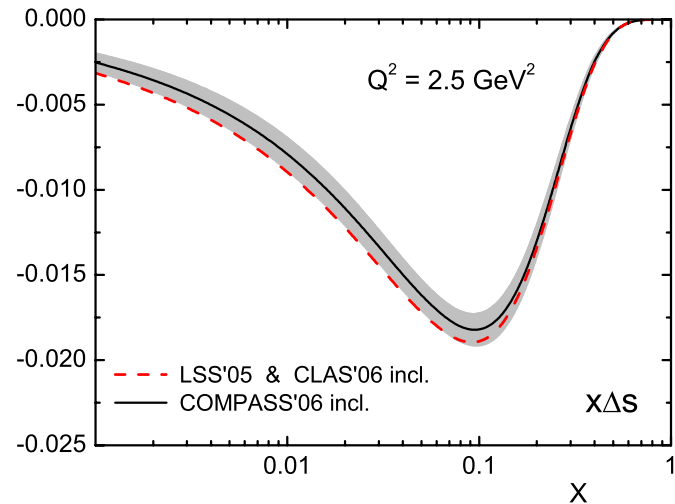
The new QCD curves corresponding to the best fits lie above the old one at $x < 0.1$

- $(\Delta u + \Delta \bar{u}), (\Delta d + \Delta \bar{d})$ do **NOT** change
- $x|\Delta s(x)|$ and $x\Delta G(x)$ and their first moments Δs and ΔG slightly **decrease**

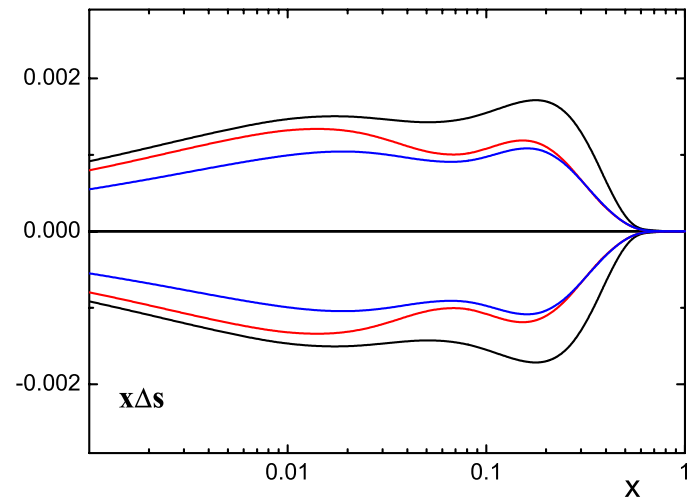
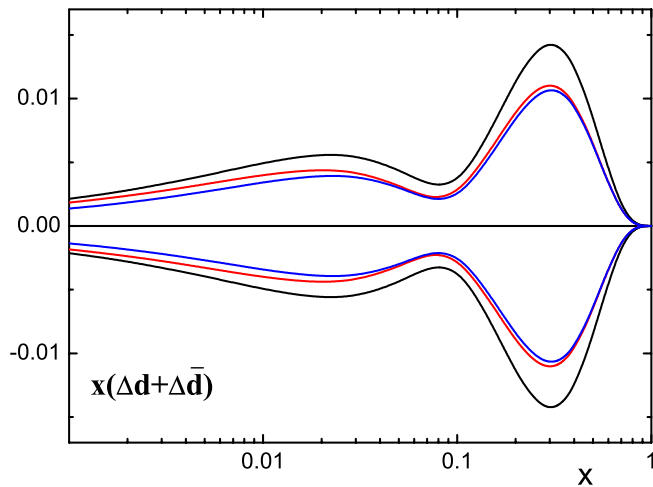
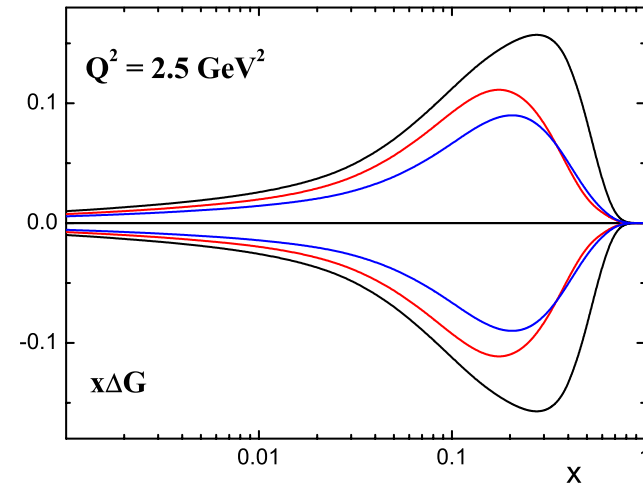
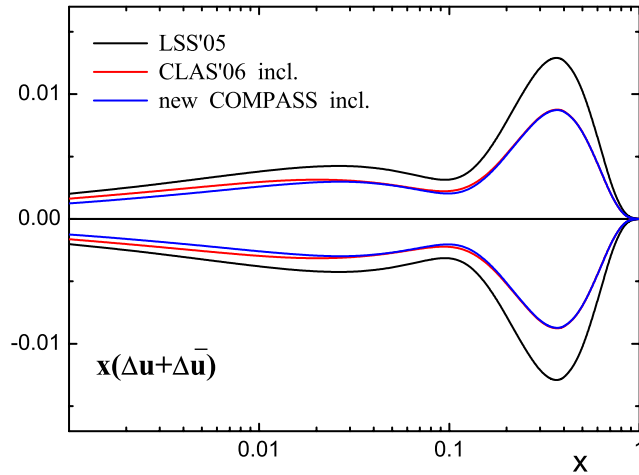
5 x-bins for HT



The values of HT are practically **NOT** affected by *COMPASS* data excepting the **small x** where **Q^2** are also **small**



Impact of COMPASS'06 data on the uncertainties for NLO polarized PD



$$Q^2 = 1 \text{ GeV}^2$$

| <i>COMPASS</i> | Δs | ΔG | $a_0 = \Delta \Sigma_{\text{MS}}$ |
|----------------|--------------------|-------------------|-----------------------------------|
| old | -0.070 ± 0.007 | 0.296 ± 0.197 | 0.164 ± 0.048 |
| new | -0.063 ± 0.005 | 0.237 ± 0.153 | 0.207 ± 0.039 |

Spin of the proton

$$\begin{aligned} S_z = 1/2 &= 1/2 \Delta \Sigma(Q^2) + \Delta G(Q^2) + L_q(Q^2) + L_g(Q^2) \\ &= \mathbf{0.34 \pm 0.16} + L_q(Q^2) + L_g(Q^2) \end{aligned}$$

The **big** uncertainty is coming from gluons

To be determined from forward extrapolations of **generalized** PD

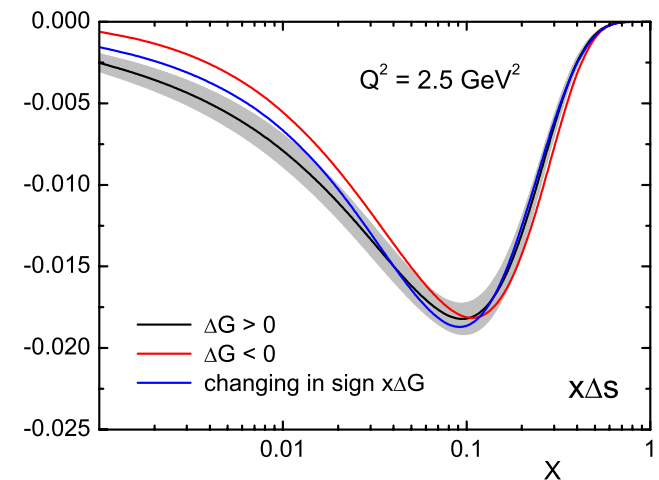
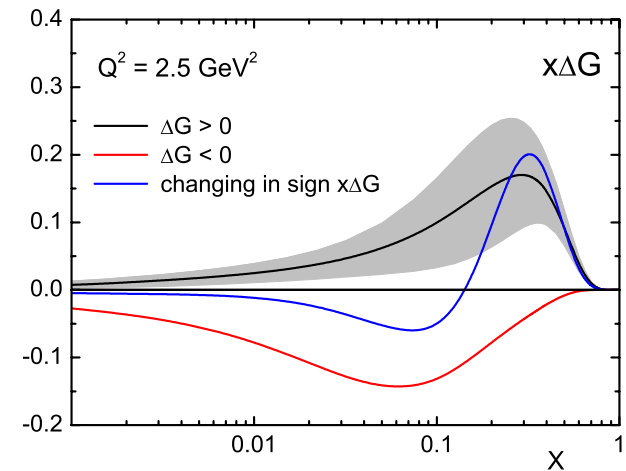
The sign of gluon polarization

- The present **inclusive** DIS data **cannot rule out** the solutions with negative and changing in sign gluon polarizations

$$\chi_{DF}^2(\Delta G > 0) = 0.892$$

$$\chi_{DF}^2(\Delta G < 0) = 0.895, \chi_{DF}^2(x\Delta G / \text{chsign}) = 0.888$$

- The shape of the negative gluon density **differs** from that of positive one
- In all the cases the magnitude of ΔG is small: $|\Delta G| \leq 0.4$ at $Q^2 = 1 \text{ GeV}^2$
- The corresponding polarized quark densities are **very close** to each other



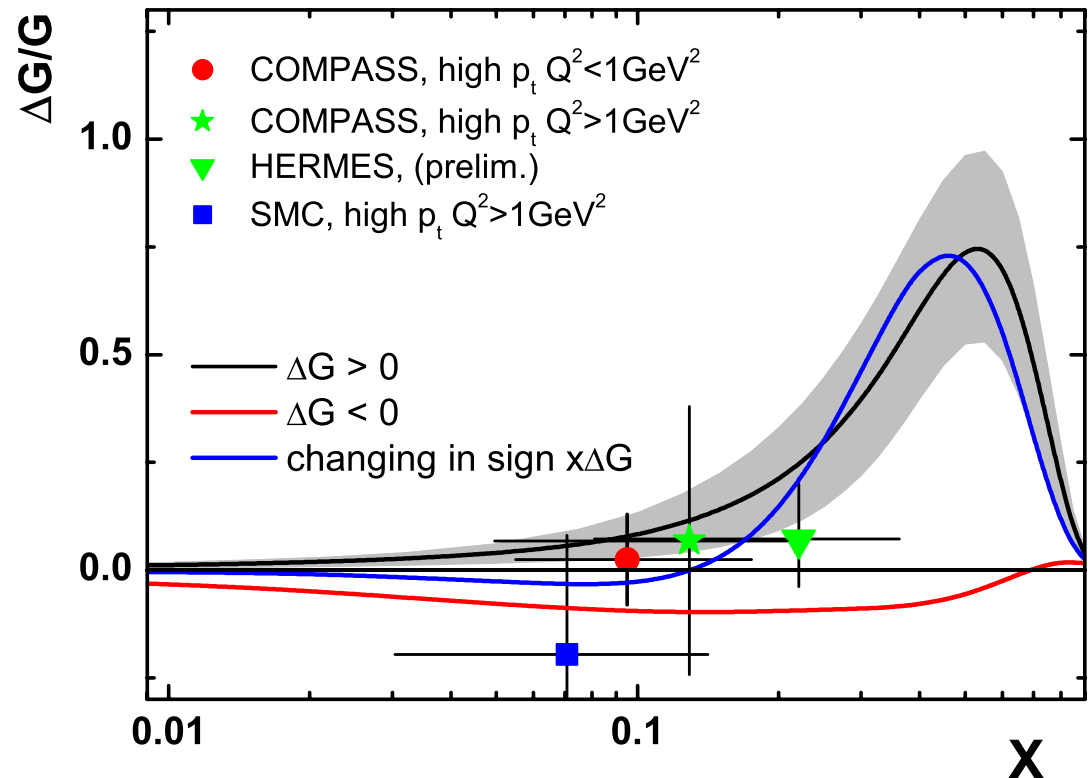
Comparison with directly measured $\Delta G/G$ at $Q^2 = 3 \text{ GeV}^2$

MRST'02 unpolarized gluon density is used for $G(x)$

The error band corresponds to statistic and systematic errors of ΔG

The error bars of the experimental points represent the **total errors**

The most precise value of $\Delta G/G$, the **COMPASS** one, is **well consistent** with any of the polarized gluon densities determined in our analysis



SUMMARY

- The **low Q^2** *CLAS* data improve **essentially** our knowledge of **higher twist** corrections to g_1 structure function
- The central values of polarized PD are **NOT affected**, but the accuracy of its determination is **essentially improved**
- The *COMPASS* data (mainly at **large Q^2**) influence $|\Delta s|$ and ΔG which slightly **decrease**, but practically do **NOT** change HT



Strong support of the QCD framework

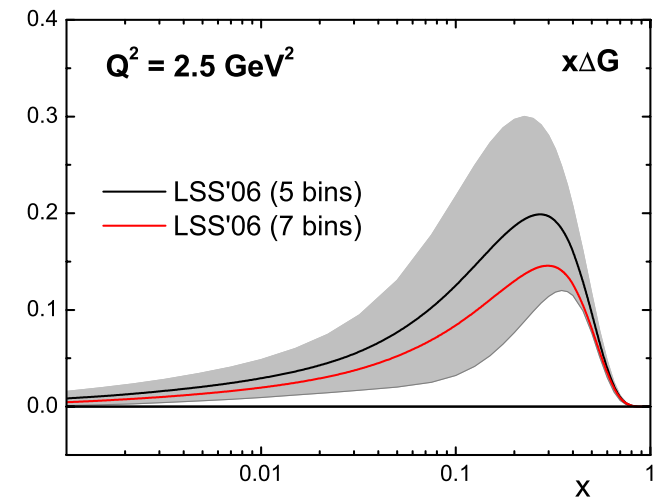
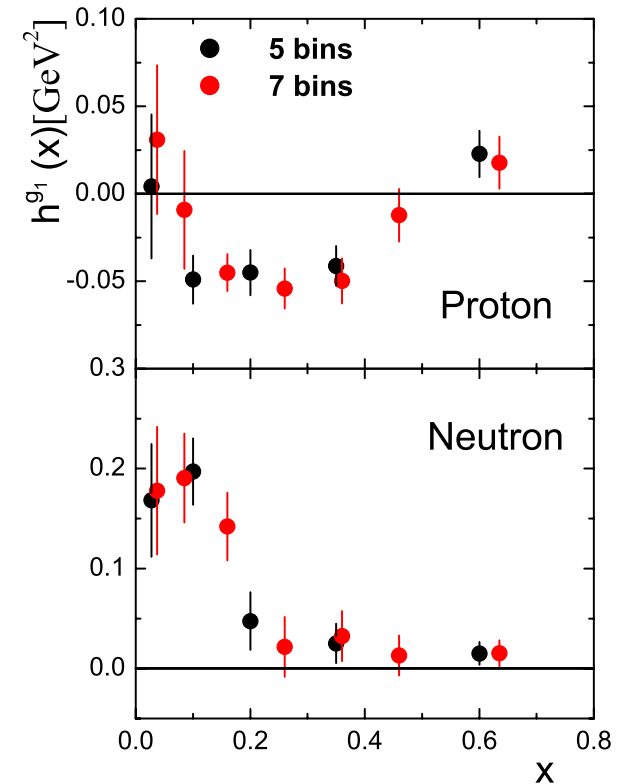
- **Large (40%)** contribution of HT to $(g_1)^d$ at small x (**low Q^2**)
- The present **inclusive** DIS data **cannot rule out** the negative and changing in sign gluon densities
- **Good agreement** with the directly measured $\Delta G/G$

OPEN QUESTIONS

- To constrain better ΔG \longrightarrow directly from *COMPASS, RHIC*;
more precise experiments on g_1^d - *JLab Hall C*
- $\Delta\bar{u}, \Delta\bar{d}$ \longrightarrow from *SIDIS (COMPASS, JLab)* and $A_L(W^{+(-)})$ at *RHIC*
- L_q (from generalized PD - *HERMES, COMPASS, JLab*) and L_g ?
- $a_8 \neq 3F - D = 0.585$? (how much $SU(3)_f$ is broken) \rightarrow *NA48*
at *CERN*
- HT corrections in *SIDIS*, $O(\Lambda^4/Q^4)$ term in HT expansion in Bjorken x-space
...etc.

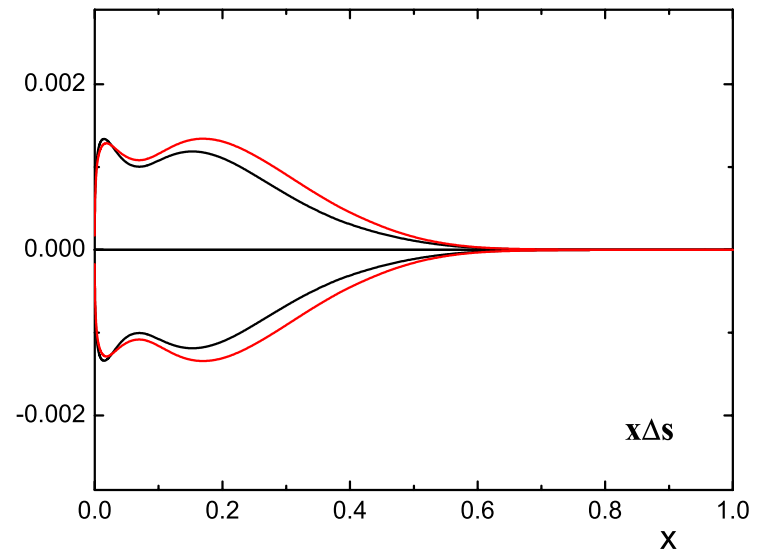
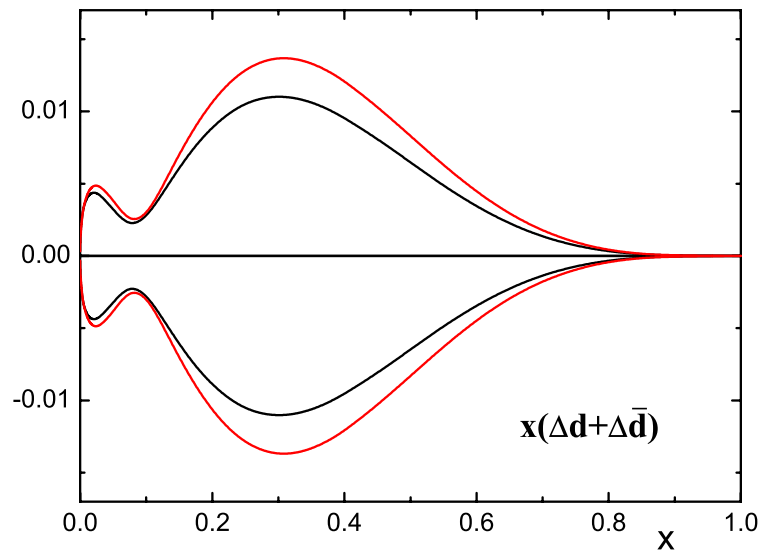
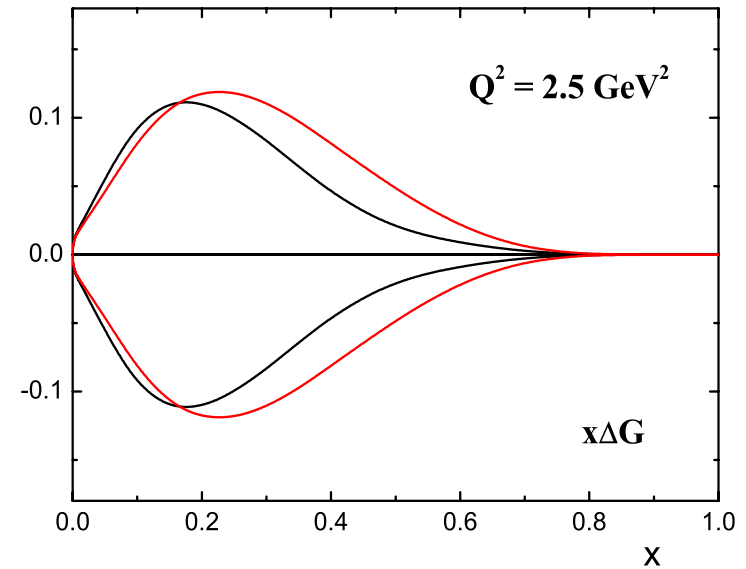
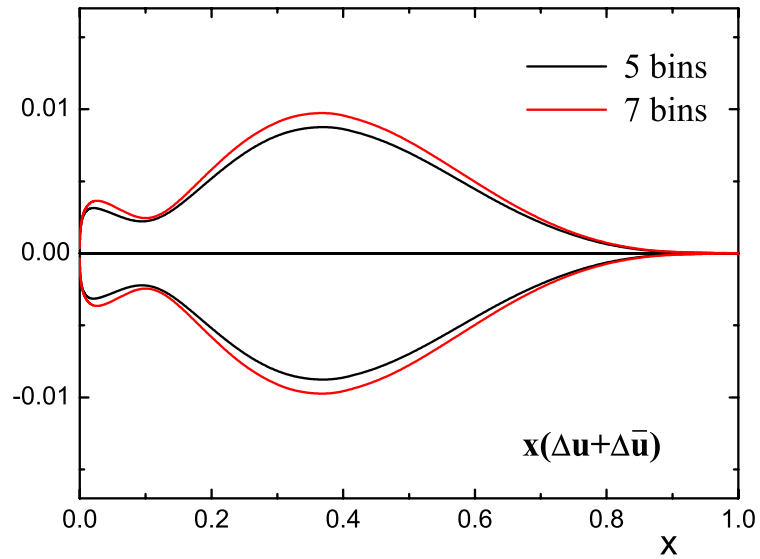
Additional slides

- Due to the good accuracy of the *CLAS* data, one can split the measured x region of the *world+CLAS* data set into **7 bins** instead of 5, and to determine **more precisely** the x -dependence of HT
- The corresponding PPD are practically **identical** with those of LSS'06 (**5 bins**)
- The only exception is $x\Delta G$, but it lies **within** the error band of $x\Delta G$ (**5 bins**) \rightarrow small correlation between gluons and HT



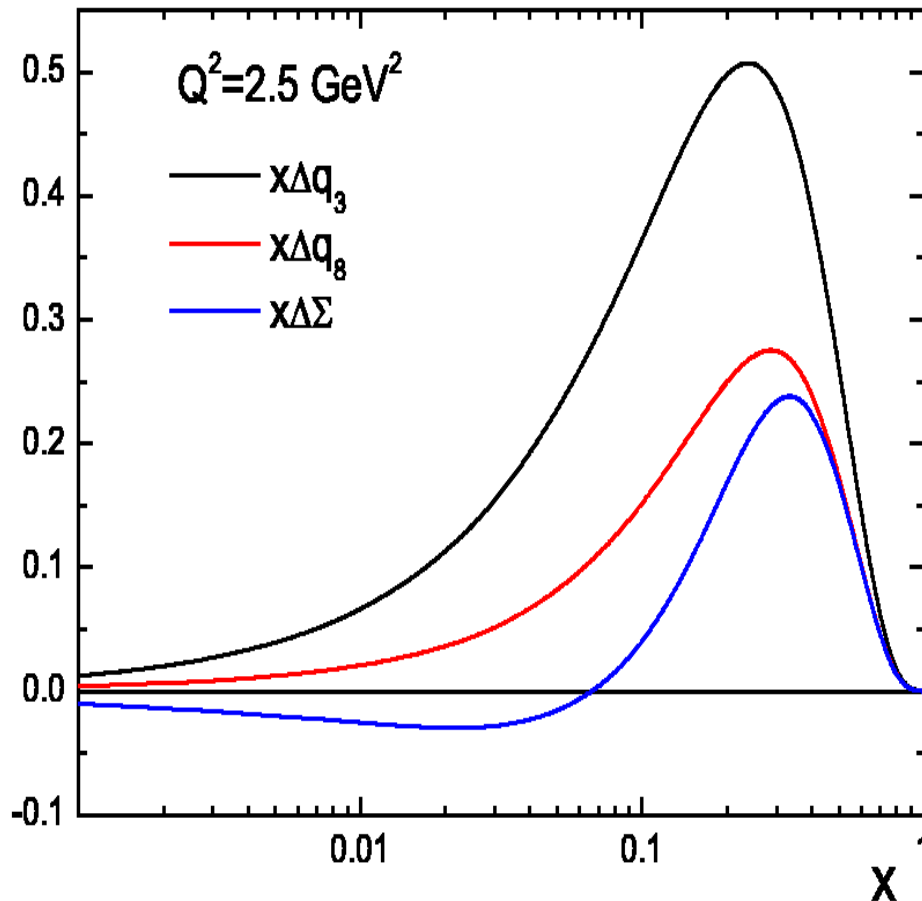
5 \rightarrow 7 x-bins

Impact on the uncertainties for NLO polarized PDFs



Why deuteron best for $\Delta G(x)$?

$$g_1^{p(n)}(x, Q^2) = \frac{1}{9} \left[\left(\pm \frac{3}{4} \Delta q_3 + \frac{1}{4} \Delta q_8 + \Delta \Sigma \right) \otimes \left(1 + \frac{\alpha_s(Q^2)}{2\pi} \delta C_q \right) + \frac{\alpha_s(Q^2)}{2\pi} \Delta G \otimes \delta C_G \right]$$



- The Δq_3 terms from p and n about twice size of Δq_8 and $\Delta \Sigma$ terms, **cancel** in deuteron.

- **Relative** gluon contributions largest in deuteron: relevant because experimental errors dominated by systematic scale factors.

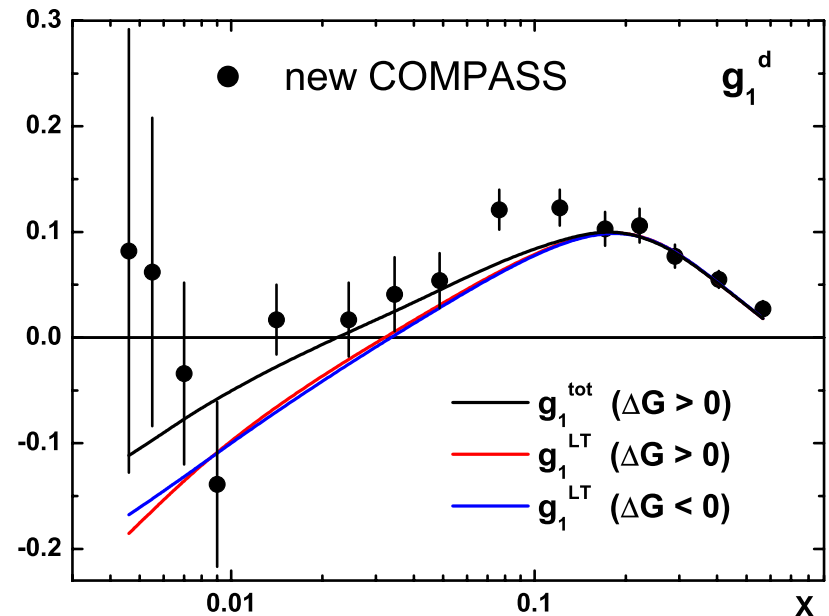
LSS'06 vs COMPASS'06

- At small x : $0.004 - 0.02$ ($Q^2 \sim 1-3 \text{ GeV}^2$) our results differ from those of *COMPASS*
- *COMPASS* → significant difference between $(g_1)_{\text{th}}$ corresponding to the best fits for $\Delta G > 0$ and $\Delta G < 0$
- *LSS'06* → the theoretical curves for both cases are very close to each other
- The reason → HT effects (40% at small x) which are NOT taken into account by *COMPASS*

$$(g_1)_{\text{exp}} \leftrightarrow$$

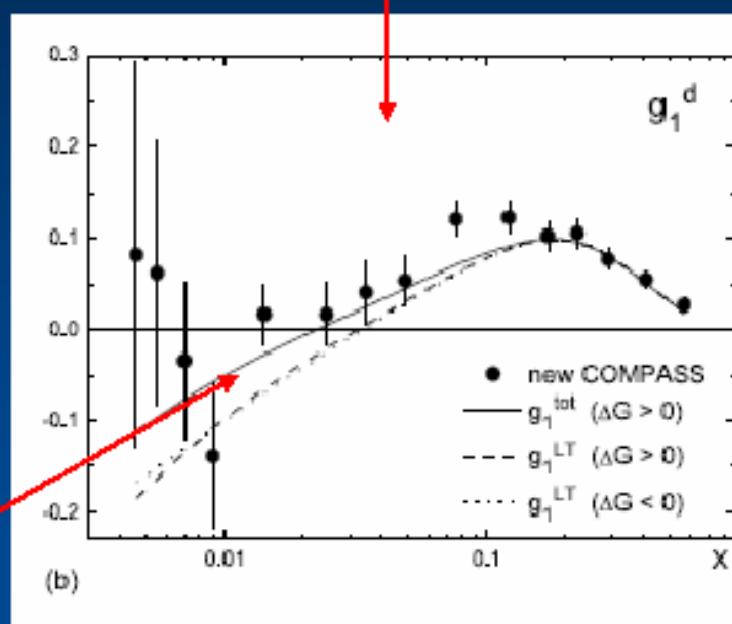
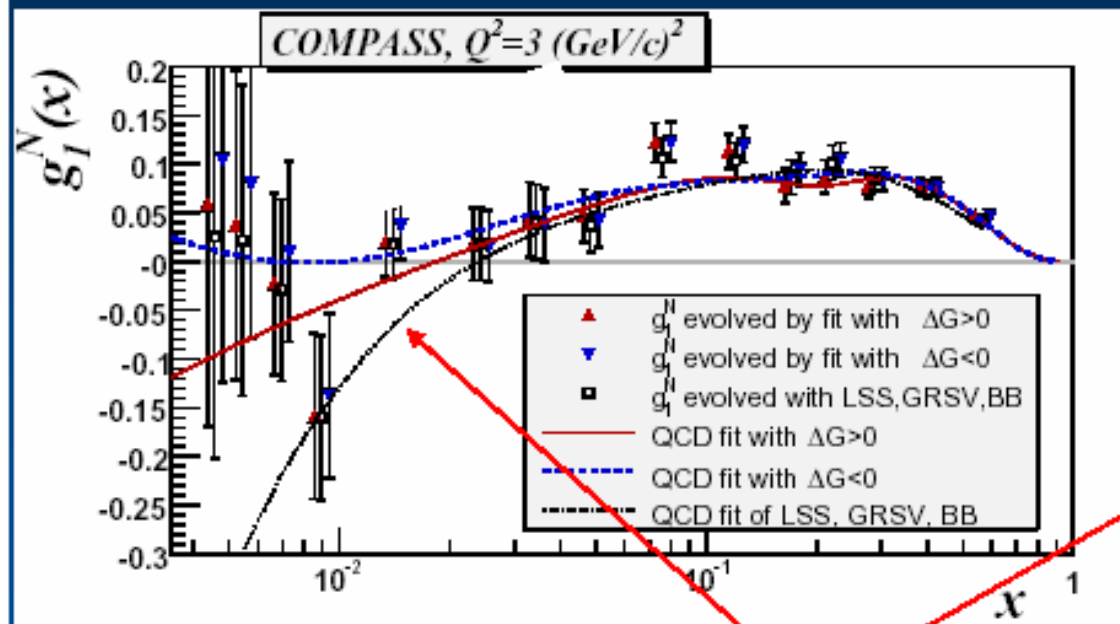
$$(g_1)_{\text{LT}}(\text{COMPASS}) \approx$$

$$(g_1)_{\text{LT}}(\text{LSS}) + h^d(x)/Q^2$$



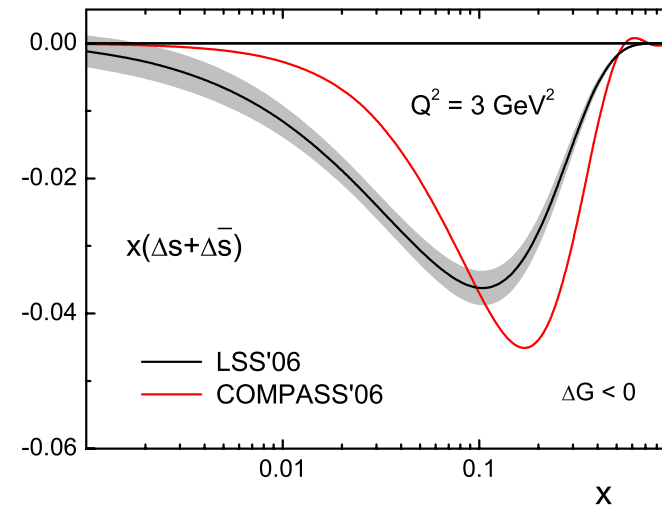
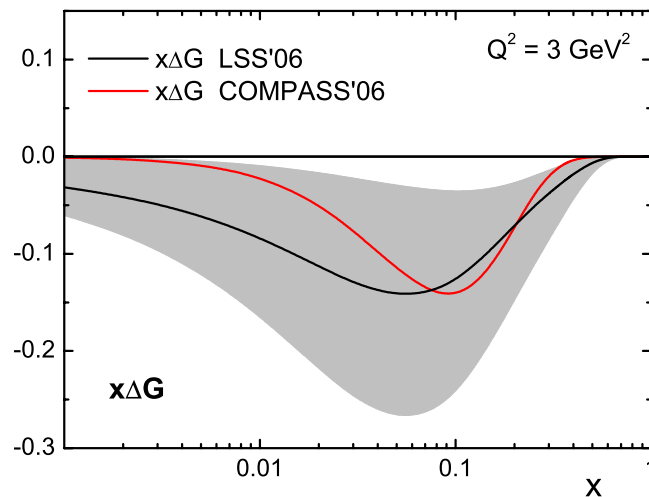
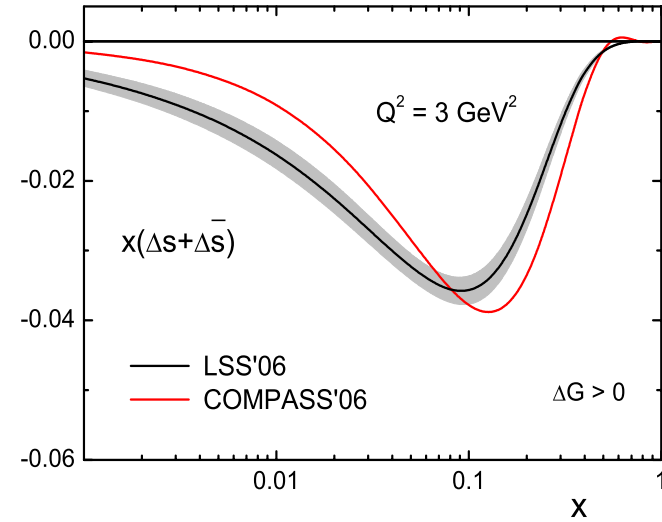
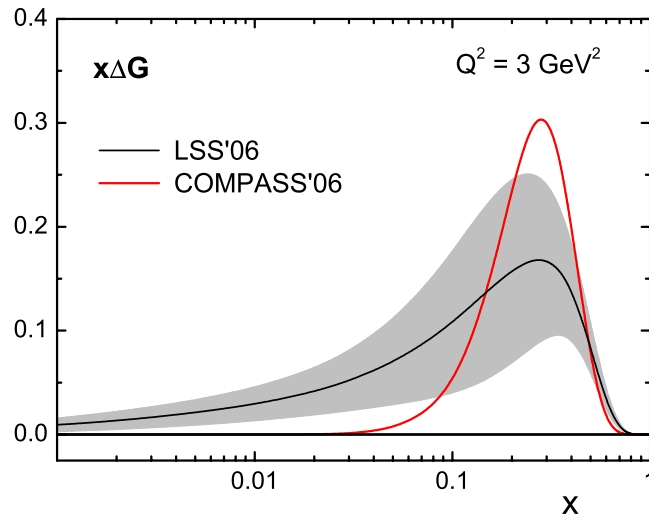
QCD analysis of the world data on structure function g_1

Comparison of data and fits - LSS06 (hep-ph/0612360)



LSS05 vs LSS06

- $x\Delta_S$ are different, especially in the case of $\Delta G < 0$
- $x\Delta_G$ positive obtained by COMPASS is more peaked than our



Constraint on ΔG from π^0 production at RHIC (*AAC, hep-ph/0612037*)

$$\vec{p} + \vec{p} \rightarrow \pi^0 + X$$

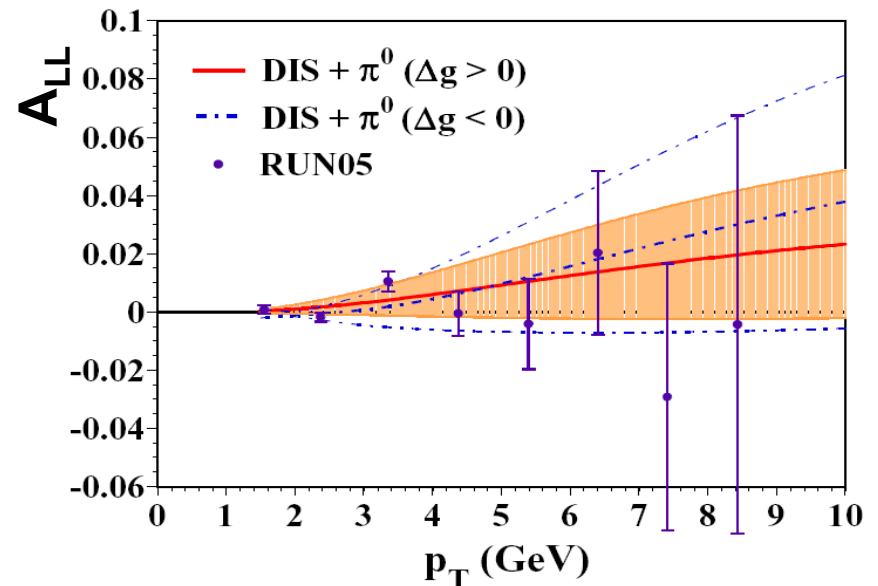
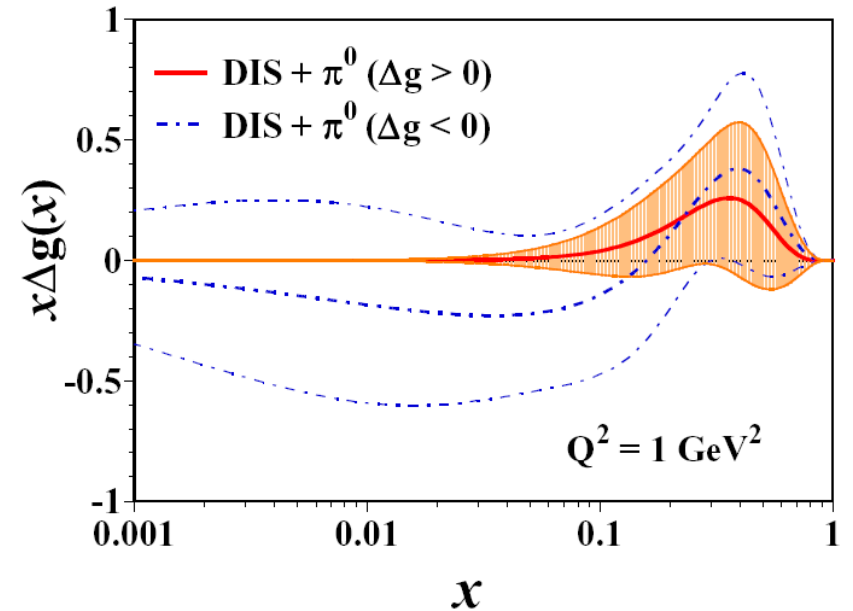
From DIS + π^0 analysis:

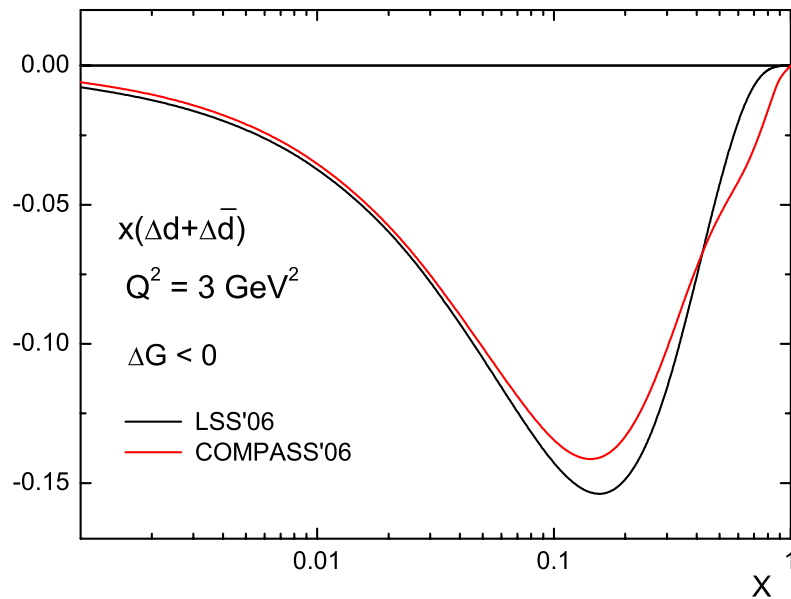
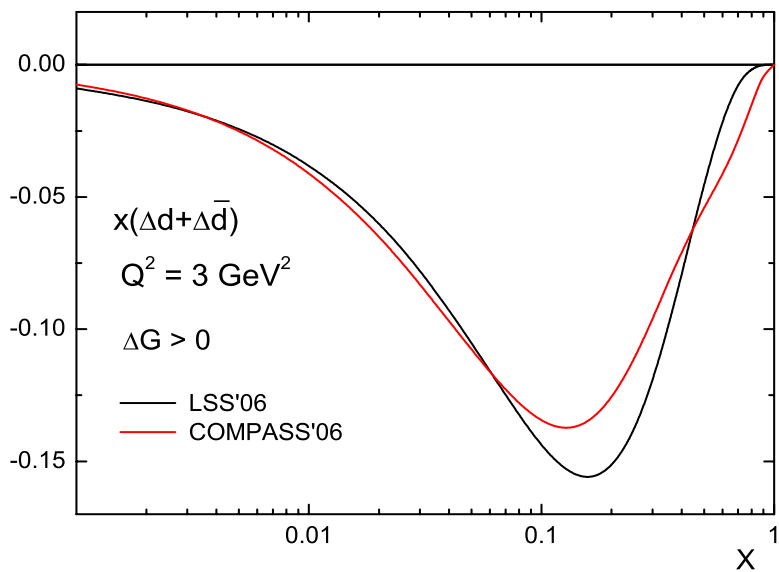
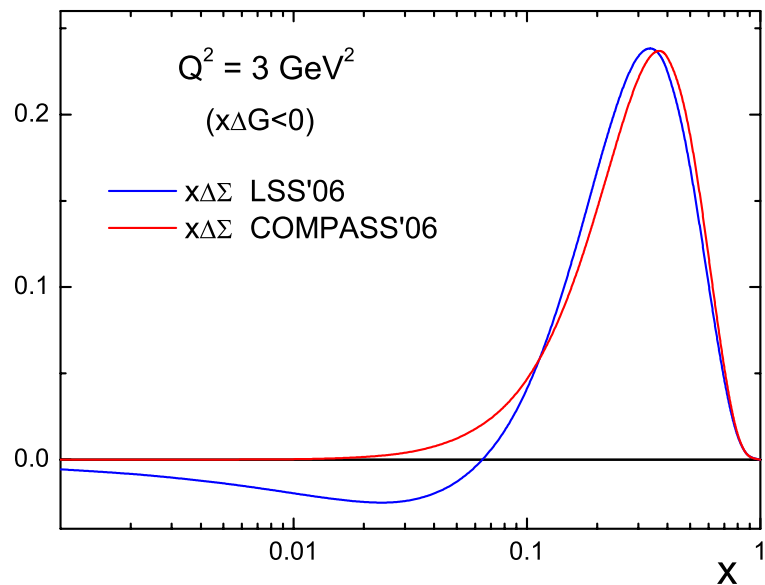
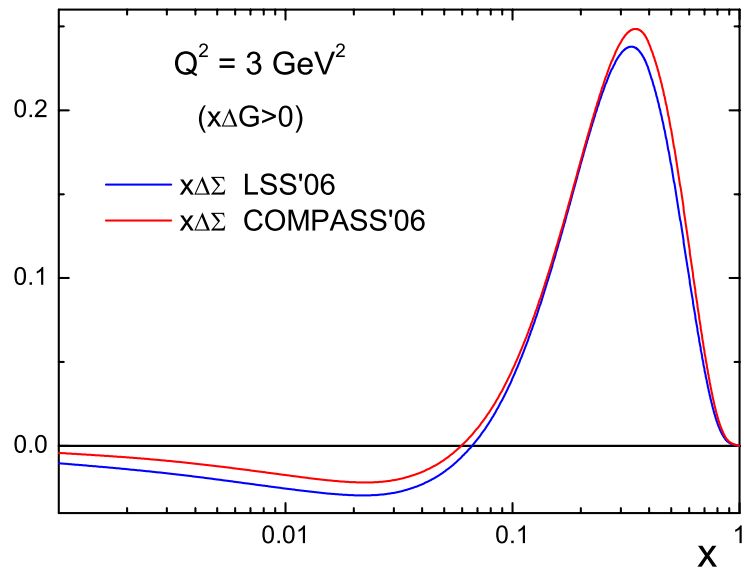
$$\Delta G = 0.31 \pm 0.32$$

$$\Delta G = -0.56 \pm 2.16$$

$$(Q^2 = 1 \text{ GeV}^2)$$

Note: In contrast to LSS changing in sign $x\Delta G$, which for $Q^2 > 6 \text{ GeV}^2$ is **positive** for any x , $x\Delta G_{\text{AAC}}$ becomes **negative** for large x too with increasing of Q^2 .





The expected uncertainties for NLO($\overline{\text{MS}}$) polarized PDFs including the CLAS12 “data” set

